



Investigation of the decays $B^0_{d,s} \rightarrow J/\psi \eta$ in the ATLAS experiment at the LHC

by
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Physics Motivation (1)

- Probing the mixing phase: $\phi_s = -2\lambda\eta = 2\lambda\sin\gamma |v_{ub}/v_{cb}|$
 $\phi_s^{SM} \sim 0.039 \rightarrow \gamma$ is very hard to extract!

$B_{s,d}^0 \rightarrow J/\psi\eta$ modes are U-spin related (d \leftrightarrow s) can be used to extract the γ -angle of the Unitarity triangle ([hep-ph9903455](#)):

- General parametrization of the $B_{s,d}^0 \rightarrow J/\psi\eta$ decay amplitudes:

$$A_{s,d}(t) \propto 1 - b_{s,d} e^{i\rho_{s,d}} e^{i\gamma} \quad \rho_{s,d} \text{ the strong phases.}$$

$$b_s = \frac{\lambda^2}{1 - \lambda^2} a_s \quad b_d = a_d$$

b_s is strongly Cabibbo suppressed \rightarrow Small CP violation effects

\rightarrow use $B_d^0 \rightarrow J/\psi\eta$ channel to calculate CP asymmetry





Physics Motivation (2)

- Calculate the time-dependent asymmetries: A_{cp}^{dir} , A_{cp}^{mix} , $A_{\Delta\Gamma}$

However due to: $(A_{cp}^{dir})^2 + (A_{cp}^{mix})^2 + (A_{\Delta\Gamma})^2 = 1$

→ need new observable :

$\langle\Gamma_s\rangle$ and $\langle\Gamma_d\rangle$ → build new observable $H=f(b,\rho,\gamma)$

3 unknowns (b,ρ,γ) , 3 independent observables :

U-spin symmetry: $a_s = a_d$ $\rho_s - \rho_d = \pi$

$$\tan \gamma \approx \frac{\sin \phi - \eta A_{cp}^{mix}}{(1-H)\cos \phi_M} \quad \phi \text{ is the } B_d\text{-}B_d \text{ mixing phase.}$$





Physics Motivation (3)

- Branching ratio of $B_{s,d}^0 \rightarrow J/\psi\eta$:

$$\text{LEP/L3} \quad \text{Br}(B_s^0 \rightarrow J/\psi\eta) < 3.8 \cdot 10^{-3}$$

By combining these modes and using the U-spin symmetry one was able to predict the branching ratios for both channels depending on the η - η' mixing angle (-20 to -10):

$$\text{Br}(B_s^0 \rightarrow J/\psi\eta) = 9.5 - 8.3 \cdot 10^{-4}$$

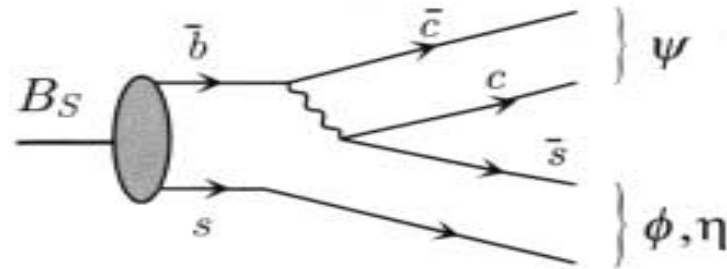
$$\text{Br}(B_d^0 \rightarrow J/\psi\eta) = 1.6 - 4.1 \cdot 10^{-6}$$





Physics Motivation (5)

- Decays $B_s^0 \rightarrow J/\psi\phi$, $B_{s,d}^0 \rightarrow J/\psi\eta$, $B_{s,d}^0 \rightarrow D_s^+ D_s^-$ are analogous, but up to now, only the first one has been studied.
- Advantages of $B_s^0 \rightarrow J/\psi\eta$: no angular analysis needed since the final state is a CP eigenstate (V-PS).



- Cross-check of the decay mode $B_s^0 \rightarrow J/\psi\phi$. New input to B-Physics measurements: $B_d^0 \rightarrow J/\psi\eta \equiv B_d^0 \rightarrow D_d^+ D_d^- \Rightarrow \sin 2\beta$





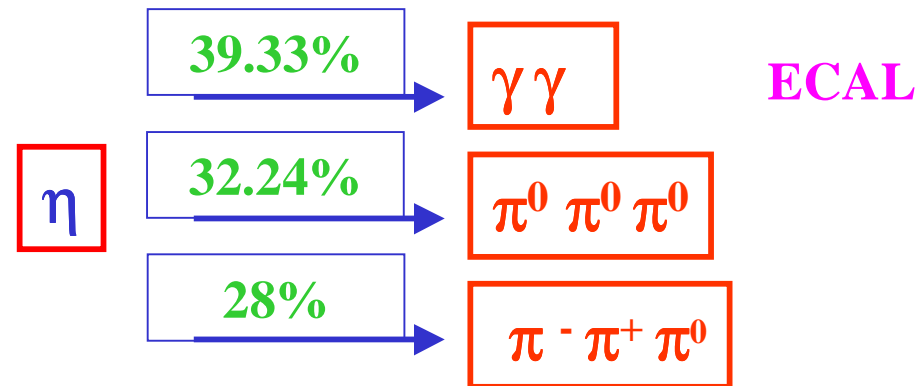
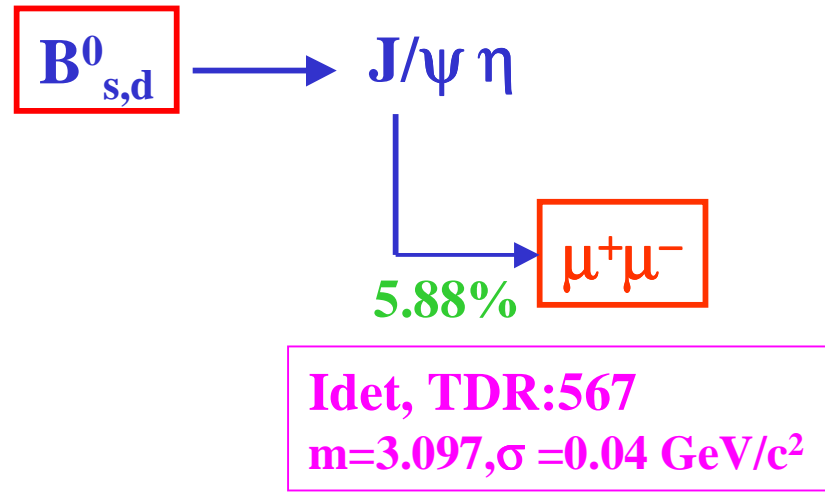
Simulation Tools

- **Generation: ATGEN-B (PYTHIA 5.7)**
 - b-quarks produced via: lowest order process, gluon splitting, flavour excitation.
 - Exclusive modes $B_{s,d} \rightarrow J/\psi\eta$ ($\mu\bar{\mu}\mu^3$, $|\eta| < 2.5$) $\eta(\gamma\gamma)$ for **signal**.
 - Inclusive mode $B \rightarrow J/\psi(\mu\bar{\mu}\mu^3$, $|\eta| < 2.5$) X for **physics back**.
 - p-p $\rightarrow J/\psi X$ is not used, harmless for $B_s \rightarrow J/\psi\phi$ (TDR).
- **Simulation: Dice 98_2, full detector**
 - srt release 1.3.0, OLD layout
 - 17 000** $B_s \rightarrow J/\psi\eta$ events, **15 000** $B_d \rightarrow J/\psi\eta$ events
 - 23 000** $B \rightarrow J/\psi X$ events
- **Reconstruction: ATRECON**
 - srt release 1.3.0, using xKalman





Physics analysis





J/psi Reconstruction

- Pairs of muons with $p_T(\mu_1) > 6 \text{ GeV}$, $p_T(\mu_2) > 3 \text{ GeV}$ have been selected: efficiency: 85% for μ_1 (incl. level-1 and reconstruction in muon spectrometer), 78% for μ_2 (incl. Reconstruction in muon spectrometer and tilecal). Fake rate is very small.
 - Cuts:
 - Fit Quality: $\chi^2/\text{ndf} < 6$
 - Transverse Decay Length: $L_{xy} > 250 \text{ } \mu\text{m}$
- J/ψ mass resolution ~ 39 MeV**
→ J/ψ reconstruction efficiency: 79%





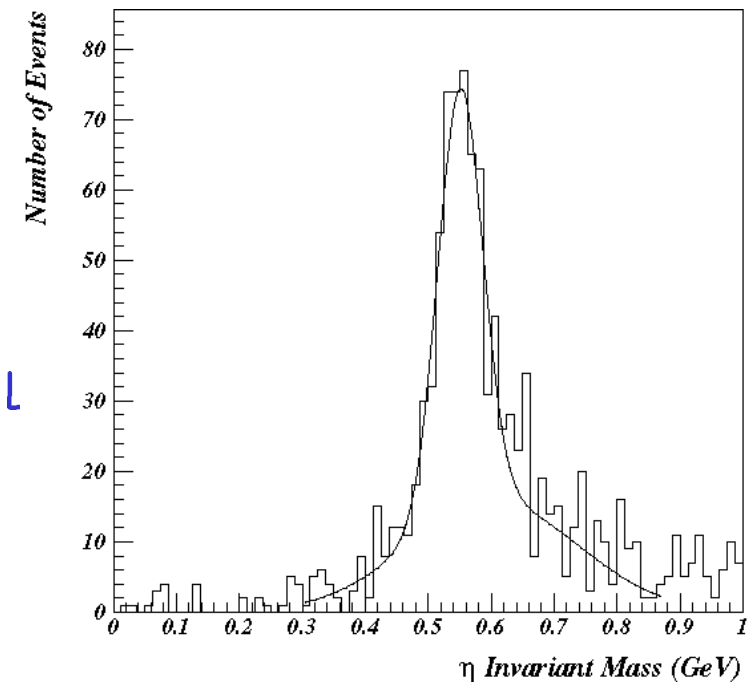
η Reconstruction

- Reconstructed in the ATLAS LAr Calorimeter:
 - Cell size: $\Delta\eta \times \Delta\phi = 0.025 \times 0.025$
 - Energy resolution: $\sigma(E)/E = 10\%/\sqrt{E} \oplus 1\%$
 - Cluster size 3x3 cells with $E_{\text{thr}} > 1 \text{ GeV}$:
 - $\gamma\gamma$ opening angle $> 5^\circ$

→ 20% of $\eta \rightarrow \gamma\gamma$ retained.

- Further selection cuts:
 - Energy sharing in the 3 ECAL/lay and HCAL
 - Cluster width
 - Energy containment in the 3x3 window

→ η mass resolution $\sim 35 \text{ MeV}$
→ Reconstruction efficiency $\sim 2.3\%$

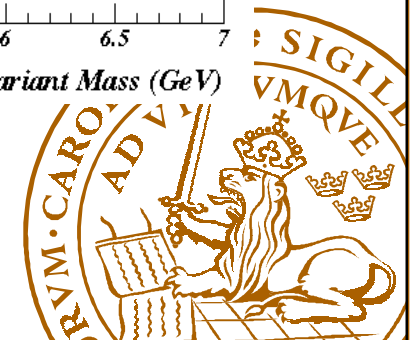
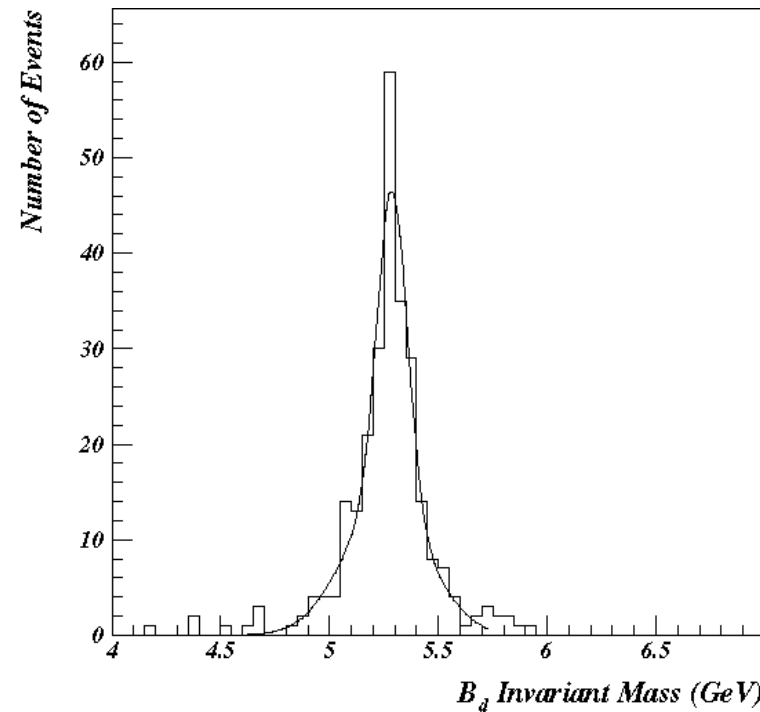
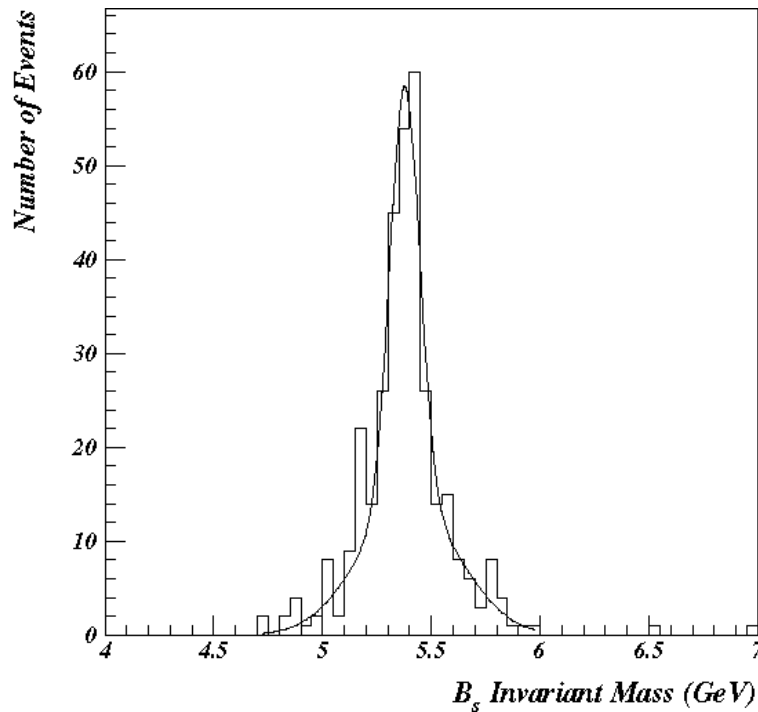




$B^0_{s,d}$ Reconstruction

Combining reconstructed J/ψ and η candidates, and choosing events with $P_{\tau}(B) > 5 \text{ GeV}$.

→ mass resolution : $\sigma(B^0_{s,d}) = 67 \text{ MeV}$.





Results for 30 fb⁻¹ (???)

$$N_{\text{signal}}^{\text{prod}} = 2\sigma(\text{pp} \rightarrow \text{b}\bar{\text{b}}) \text{Br}(\text{b} \rightarrow \text{B}_{\text{s,d}}^0) \text{Br}(\text{B}_{\text{s,d}}^0 \rightarrow \text{J}/\psi\eta) \text{Br}(\text{J}/\psi \rightarrow \mu^+\mu^-) \cdot A(\mu 6) A(\mu 3) \text{Br}(\eta \rightarrow \gamma\gamma) \int \mathcal{L} dt,$$

$$N_{\text{back}}^{\text{prod}} = \sigma(\text{pp} \rightarrow \text{b}\bar{\text{b}} \rightarrow \text{J}/\psi\text{X}) A(\mu 6) A(\mu 3) \int \mathcal{L} dt.$$

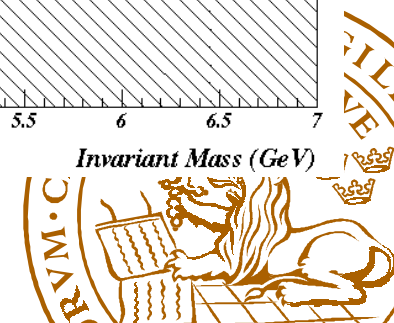
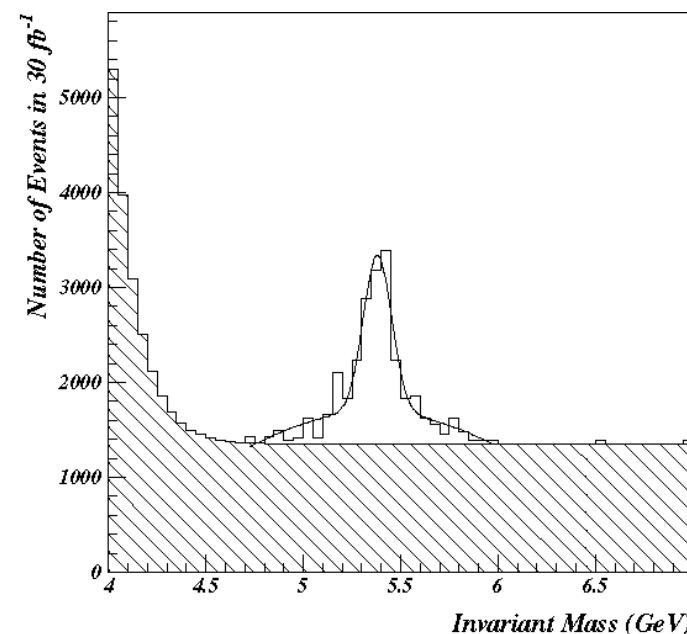
$$N_{\text{Bs}}^{\text{prod}} = 9.2 \cdot 10^8 \cdot \text{Br}(\text{B}_s^0 \rightarrow \text{J}/\psi\eta),$$

$$N_{\text{Bd}}^{\text{prod}} = 4.9 \cdot 10^9 \cdot \text{Br}(\text{B}_d^0 \rightarrow \text{J}/\psi\eta).$$

$$N_{\text{back}}^{\text{prod}} = 1.2 \cdot 10^8.$$

	$\theta_P = -10^\circ$	$\theta_P = -20^\circ$
$\text{Br}(\text{B}_s^0 \rightarrow \text{J}/\psi\eta)$	$8.3 \cdot 10^{-4}$	$9.5 \cdot 10^{-4}$
$N_{\text{Bs}}^{\text{obs}}$	8 500	9 700
$\text{Br}(\text{B}_d^0 \rightarrow \text{J}/\psi\eta)$	$4.1 \cdot 10^{-6}$	$1.6 \cdot 10^{-6}$
$N_{\text{Bd}}^{\text{obs}}$	200	80
$N_{\text{back}}^{\text{obs}}$	10 800	10 800

- Clear B_s signal S/B~0.8
- B_d CANNOT be observed





CP-asymmetry measurement (1)

Reminder:

Decay amplitude parametrization :

$$A_s(t) \propto 1 + a_s \frac{\lambda^2}{1 - \lambda^2} e^{i\theta_s} e^{i\gamma} \quad A_d(t) \propto 1 - a_d e^{i\theta_d} e^{i\gamma}$$

where $\theta_{s,d}$ the strong phases, $a_{s,d}$ their respective amplitude, λ the wolfenstein parameter.

CP asymmetry is doubly cabbibo suppressed in the $B_s \rightarrow J/\psi$.

Since the $B_d \rightarrow J/\psi\eta$ cannot be observed, one thought of calculating the CP asymmetry in the $B_s \rightarrow J/\psi\eta$ channel.





CP-asymmetry measurement (2)

- Events need to be tagged as B or anti-B:
 - **Lepton tagging:** Not possible due to low statistics.
 - **Jet Charge Tagging (SST):** exploits the correlation between the charge of a jet and the charge of the quark producing the jet:

$$Q_{jet} = \frac{\sum_i q_i P_i^k}{\sum_i P_i^k}$$

q_i is the charge of the i^{th} particle in the jet, P_i its momentum.

Charged particles with: $P_T > 0.5 \text{ GeV}$ and $|\eta| < 2.5$, $\Delta R < 0.8$, $d_0 < 1 \text{ cm}$, $|z| < 5 \text{ cm}$ were considered. $\mu^+ \mu^-$ were excluded.





CP-asymmetry measurement (3)

The reconstructed B-meson was defined as B_s -meson (anti B_s -meson) if the $Q_{jet} > +c$ ($Q_{jet} < -c$). c is a tunable cut.

The exponent (K) was optimized to maximize the the tagging quality factor $Q = \epsilon_{tag} \cdot D_{tag}^2$, where D_{tag} is the dilution factor and ϵ_{tag} is the tagging efficiency

Optimal parameters were found to be: $K=0.2$, $c=0.2$

$$\rightarrow Q = 3.85\% \quad D_{tag} = 0.26 \quad \epsilon_{tag} = 0.57$$





CP-asymmetry measurement (4)

The observable asymmetry is:

$$a_{obs}(t) = D \cdot a_{cp}(t) = D \cdot \sin\phi_s \cdot \sin\Delta m_s t$$

where $D = D_{tag} \cdot D_{back}$, $D_{back} = N_s / (N_s + N_{back})$.

Assuming the decay time resolution to be $\Delta\tau = 0.073\text{ps}$, the B_s lifetime $\tau_{B_s} = 1.464\text{ps}$:

Fit:

$$\begin{aligned} \rightarrow \delta(\sin\phi_s) &= 0.27 \quad \text{for } x_s = \Delta m_s / \Gamma_s = 19 \\ \rightarrow \delta(\sin\phi_s) &= 0.31 \quad \text{for } x_s = \Delta m_s / \Gamma_s = 30 \end{aligned}$$



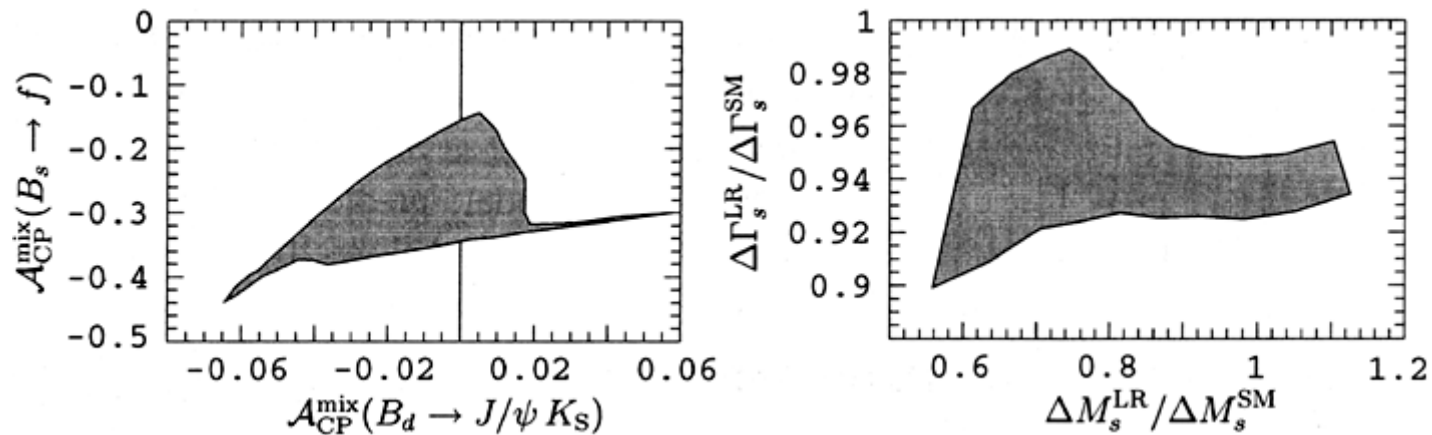


CP-asymmetry measurement (5)

In case of New Physics, ex: SB-LR model (phys. lett. B 475 (2000) 111)

$$A_{CP}^{\text{mix}}(B_s \rightarrow f) = \sin \phi_s, \quad \phi_s = \phi_s^{\text{SM}} + \phi_s^{\text{NP}} = -2\lambda^2\eta + \phi_s^{\text{NP}}$$

$$\Delta\Gamma_s = \Delta\Gamma_s^{\text{SM}} \cos(\phi_s)$$



- CP asymmetry as large as $\vartheta(40\%)$ may arise in the B_s channel while for $B_d \rightarrow J/\psi K_s$, A_{CP} as large as $\vartheta(10\%)$.
- Affects also the B^0 - B_s^0 mass and width differences, where Δm_s may be reduced significantly





Lund Layout v.s. Old layout

Geometry has been changed

- Simulation: 6000 B_s --> $J/\psi\eta$ events (srt release 1.3.7).
- Reconstruction: xKalman++.

Calo : No damage caused ($\eta \rightarrow \gamma\gamma$).

INDet : Mass resolution of ($J/\psi \rightarrow \mu^+\mu^-$) ~ 50 MeV.





Comparison and Conclusions

Simple estimates:

- B-factories PEPII, KEKB cannot produce B_s
- Tevatron CDF 2 fb^{-1} : 190 $B_s \rightarrow J/\psi \eta$
- LHC ATLAS 30 fb^{-1} : 9700 $B_s \rightarrow J/\psi \eta$
- LHC LHCb 10 fb^{-1} : 260 000 $B_s \rightarrow J/\psi \eta$, 2150 $B_d \rightarrow J/\psi \eta$

ATLAS can see a clear signal of $B_s \rightarrow J/\psi \eta$, measure BR and other parameters measurement (mixing parameters), cross-check of other experiments. New sizeable effects beyond the SM can be observed.

Improve statistics by

- investigating $\eta \rightarrow \pi^+ \pi^- \pi^0$ (BR=23%)
- improve η efficiency by better use of the fine granularity of the first ECAL sampling (cluster fine structure)

