

B Physics in LHC Era

The 17'th Nordic Particle Physics Meeting

4-10 January 2002

Spåtind, Norway



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CERN

and

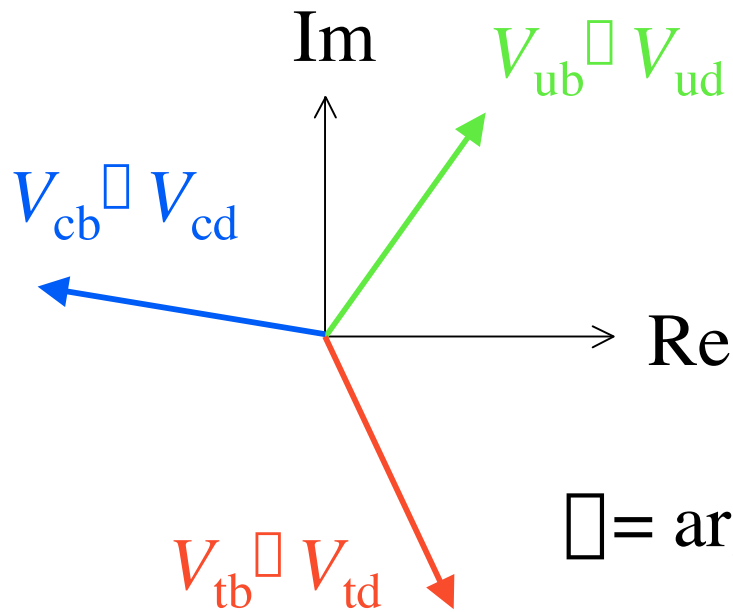
University of Lausanne

1) CKM picture

$$V = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

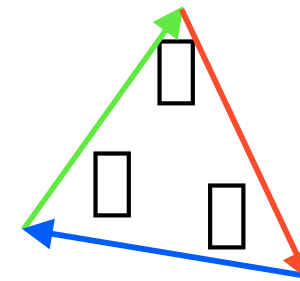
Unitarity: $V^\dagger V = \mathbf{1}$

1) $V_{ub}^* V_{ud} + V_{cb}^* V_{cd} + V_{tb}^* V_{td} = 0$



$$\alpha = \arg \frac{V_{tb}^* V_{td}}{V_{ub}^* V_{ud}}$$

Unitarity means \Rightarrow



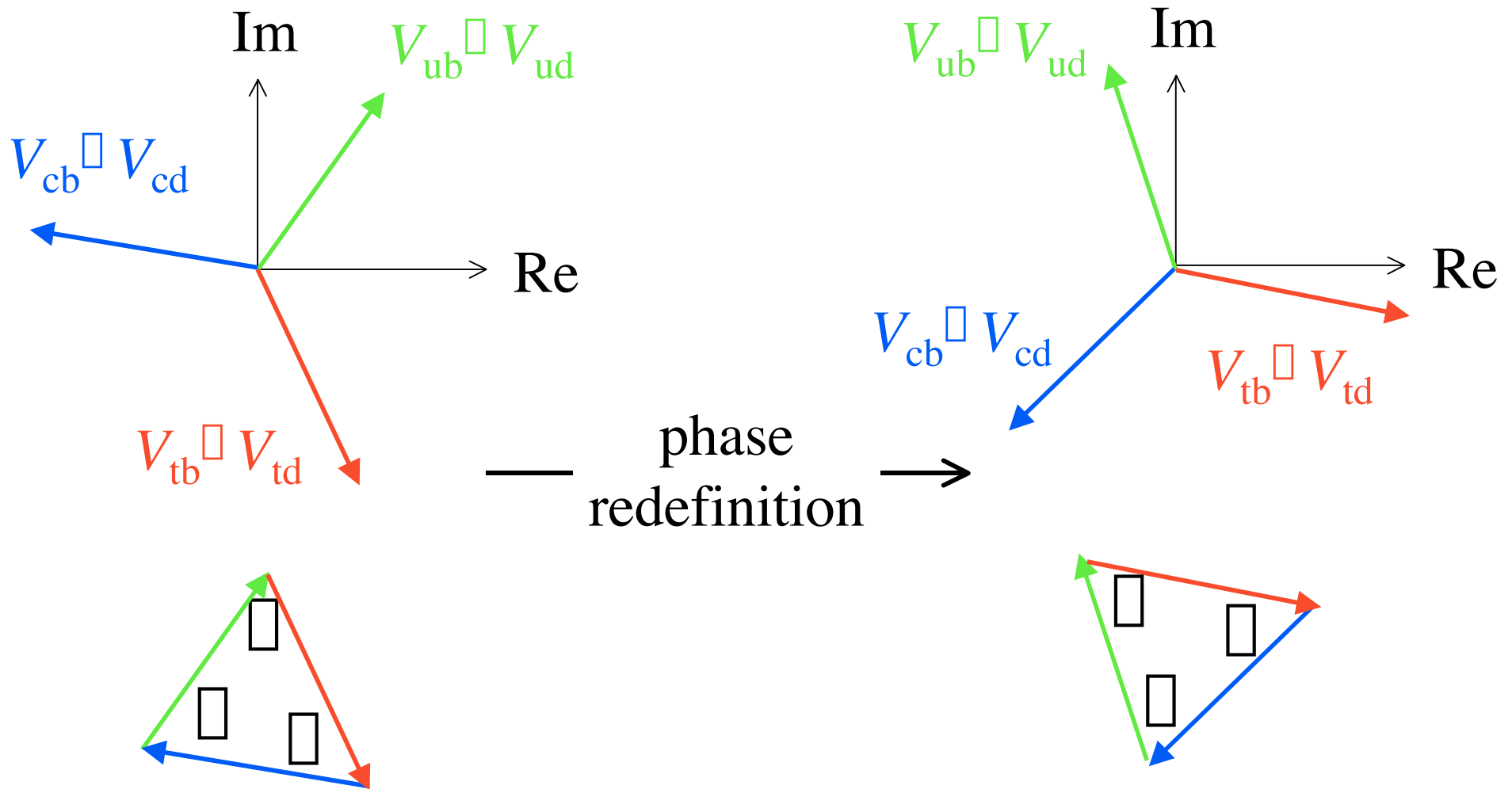
$$\beta = \arg \frac{V_{ub}^* V_{ud}}{V_{cb}^* V_{cd}}$$

$$\gamma = \arg \frac{V_{tb}^* V_{td}}{V_{cb}^* V_{cd}}$$

$$\alpha + \beta + \gamma = \pi$$

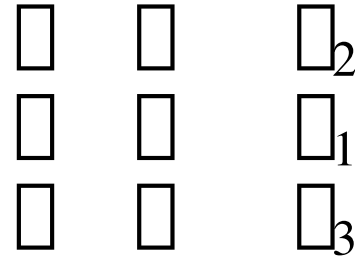
$$\beta = \arg \frac{V_{tb} V_{td}}{V_{ub} V_{ud}} \quad \beta = \beta \quad \beta = \arg \frac{V_{tb} V_{td}}{V_{cb} V_{cd}} \quad \beta = \arg \frac{V_{ub} V_{ud}}{V_{cb} V_{cd}}$$

- invariant under quark phase redefinition -



英和辞典

West Cost

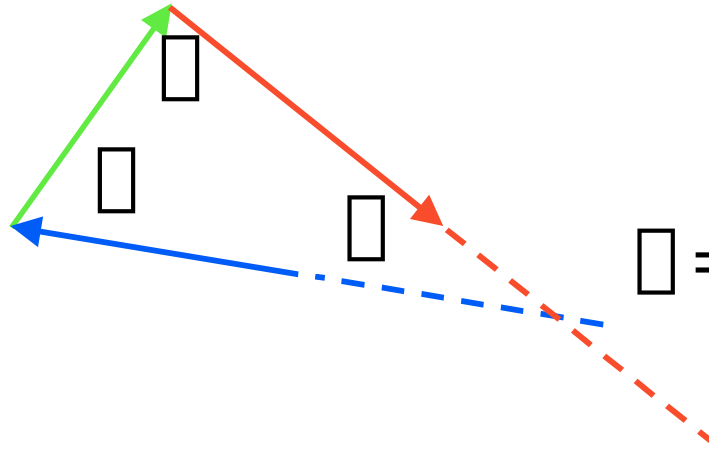


Far East

If CKM is not unitary...

$$\alpha = \arg \frac{V_{tb} V_{td}}{V_{ub} V_{ud}}$$

$$\beta = \arg \frac{V_{ub} V_{ud}}{V_{cb} V_{cd}}$$



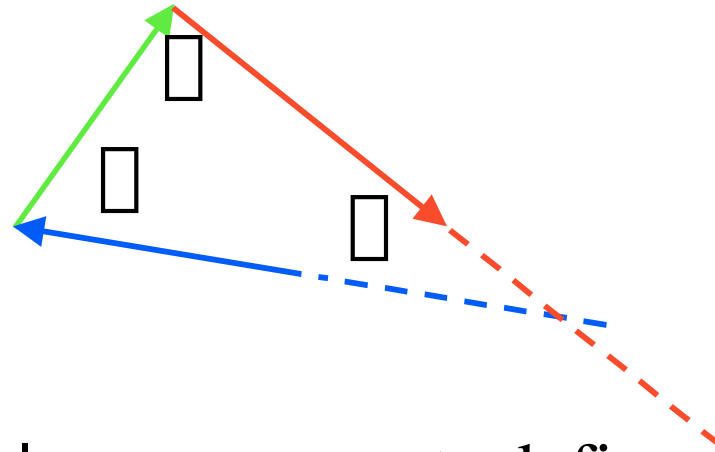
$$\gamma = \arg \frac{V_{tb} V_{td}}{V_{cb} V_{cd}}$$

$\alpha + \beta + \gamma = \pi$ is still valid!!

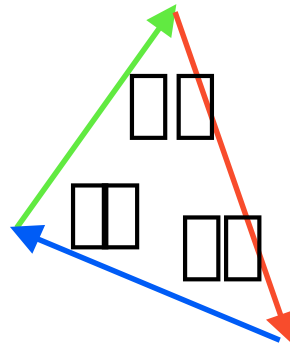
This is not a test of unitarity.

CP violation measurements

measure the angles (α and β):

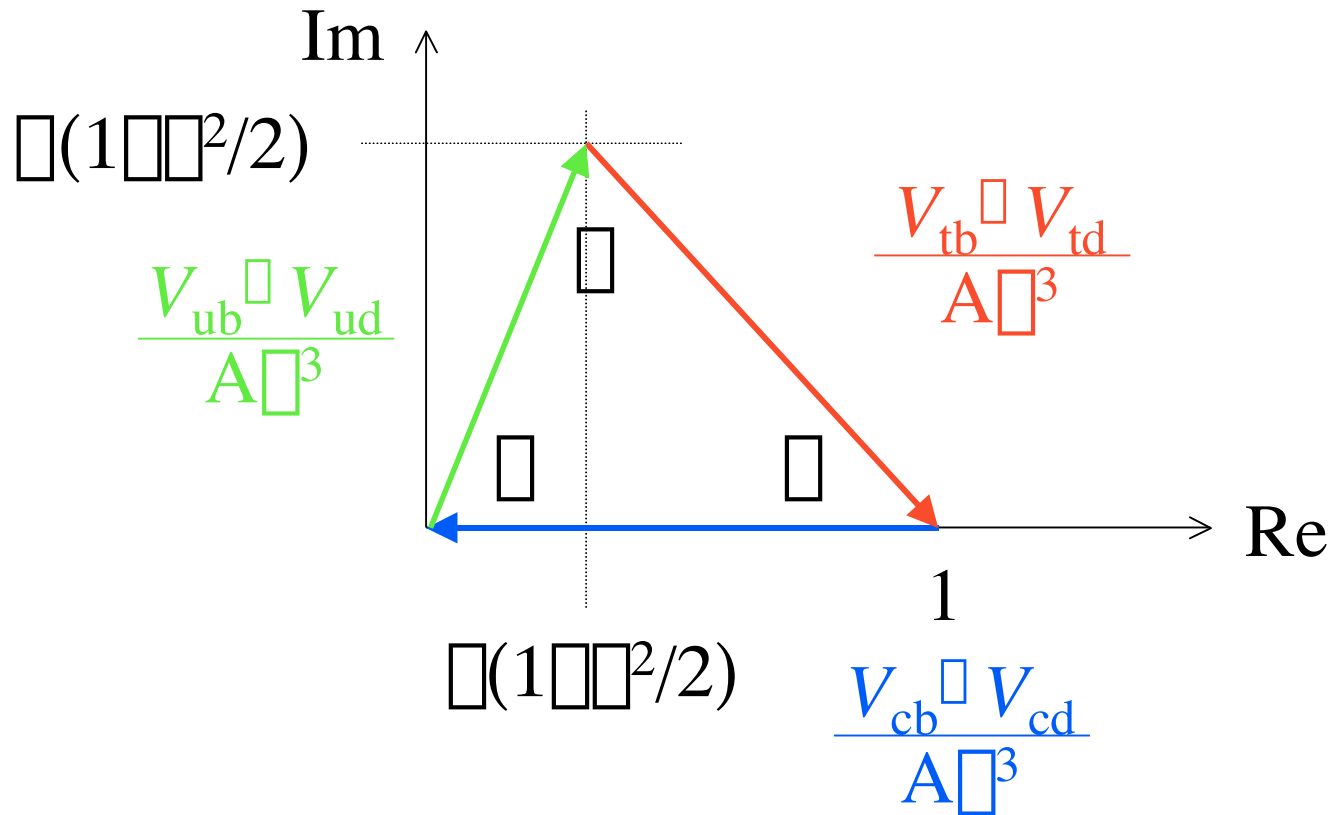


$|V_{cb}|$, $|V_{ub}|$ and $|V_{td}|$ measurements define a triangle:

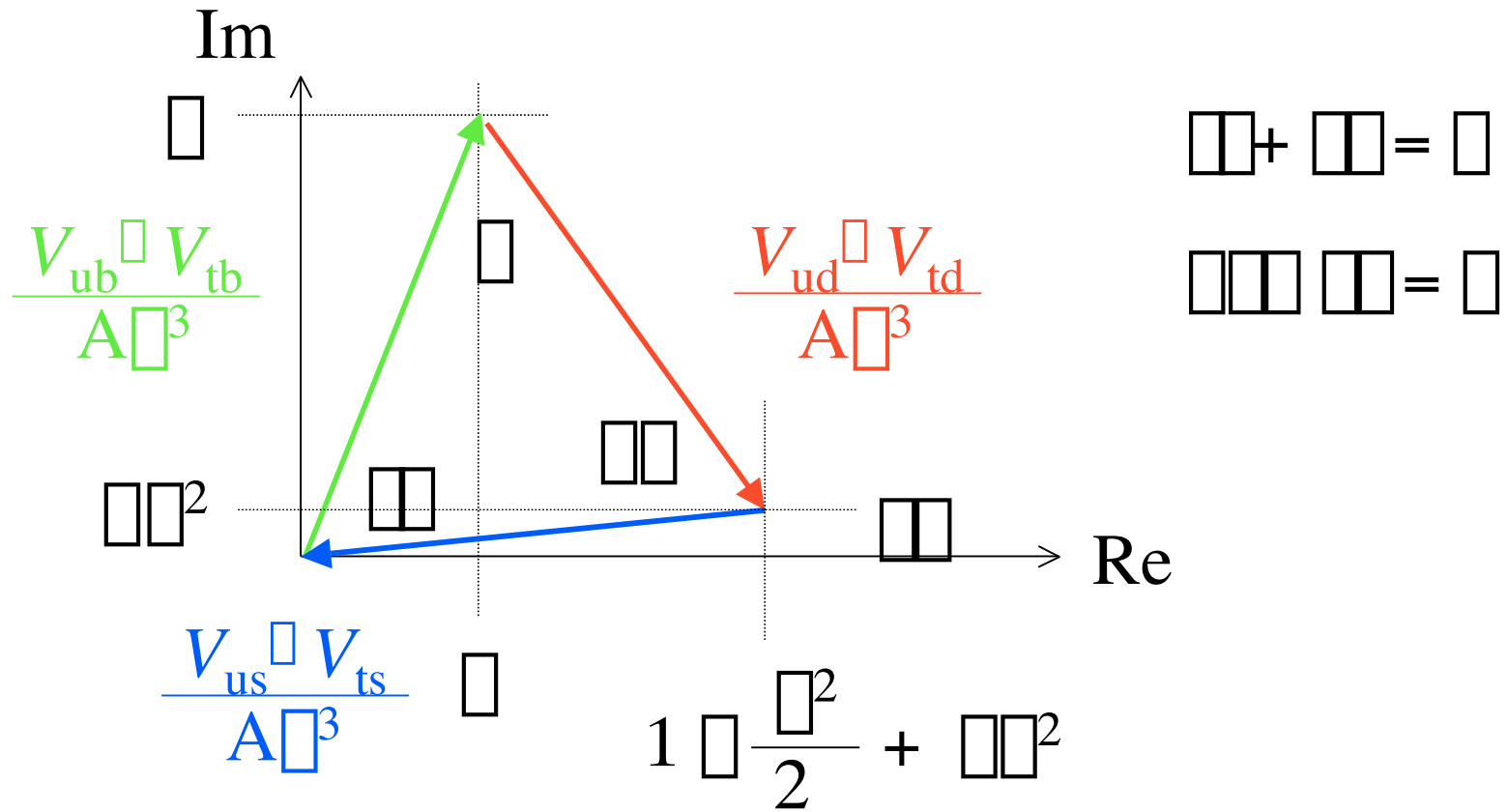


$$\alpha \neq \beta, \alpha \neq \gamma$$

Using the Wolfenstein's parametrization (ρ, A, ϕ, θ)



$$2) V_{ud} \square V_{td} + V_{us} \square V_{ts} + V_{ub} \square V_{tb} = 0$$



Important conclusions:

$$\begin{aligned} \arg V_{td} &= \alpha \\ \arg V_{ub} &= \alpha \\ \arg V_{ts} &= \alpha + \beta \end{aligned}$$

$$\beta = \tan^{-1} \frac{\lambda}{1 - \lambda} \left(1 - \frac{\lambda^2}{2(1 - \lambda)} \right)$$

$$\beta = \tan^{-1} \frac{\lambda}{\lambda}$$

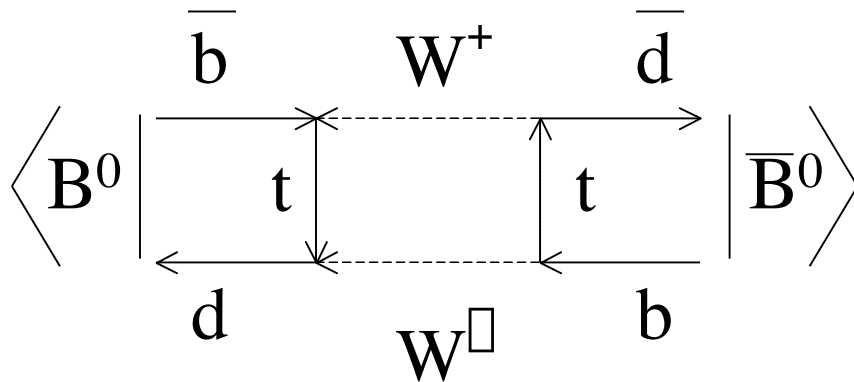
$$\lambda = \lambda^2$$

If we ignore the λ^2 correction, 1) and 2) are degenerate.

A must for the next generation experiments!

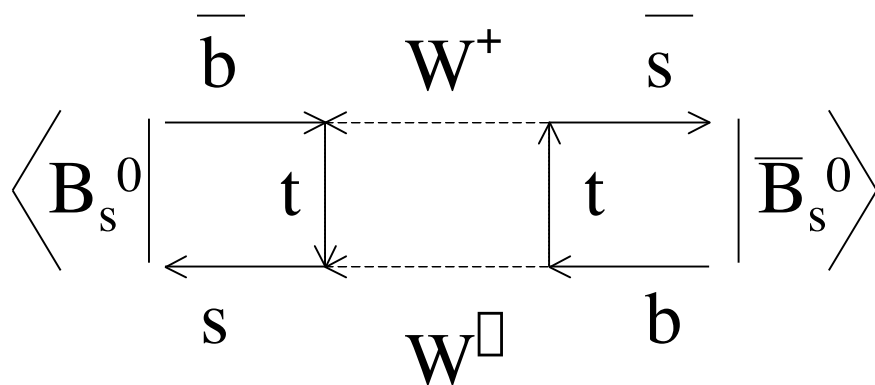
2) \bar{b} , d , and \bar{d} measurements

B- \bar{B} oscillation dispersive part: M_{12}



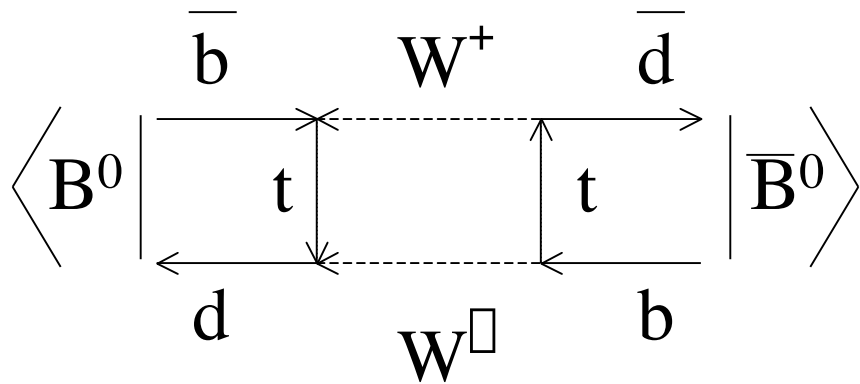
$$\Delta m = 2|M_{12}| \mu_B f_d^2 |V_{td}|^2 |V_{tb}|^2$$

$$\begin{aligned} \arg M_{12} &= \arg (V_{td}^* V_{tb})^2 + \arg \mu_B \\ &= 2\arg(V_{td}^* V_{tb}) + \arg \mu_B \end{aligned}$$

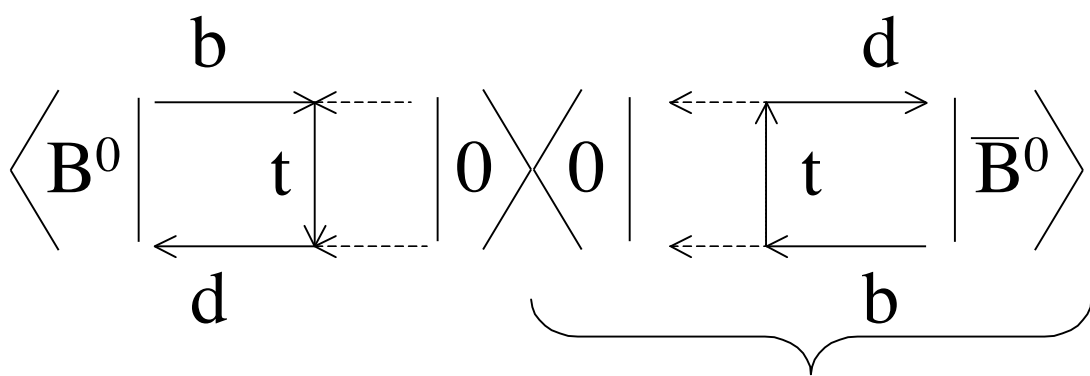


$$\Delta m_s = 2|M_{12}| \mu_{B_s} f_s^2 |V_{ts}|^2 |V_{tb}|^2$$

$$\begin{aligned} \arg M_{12} &= \arg (V_{ts}^* V_{tb})^2 + \arg \mu_{B_s} \\ &= 2\arg(V_{ts}^* V_{tb}) + \arg \mu_{B_s} \end{aligned}$$

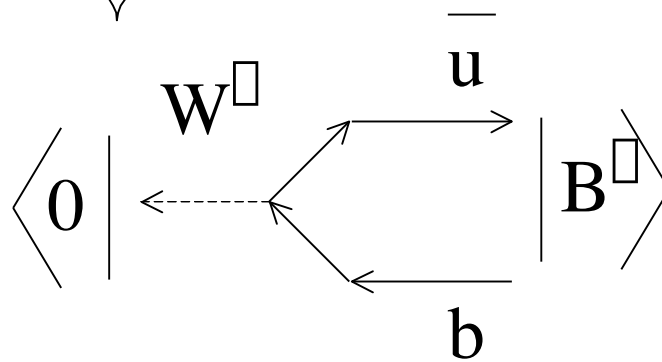


$= B$



vacuum saturation approximation

leptonic decay



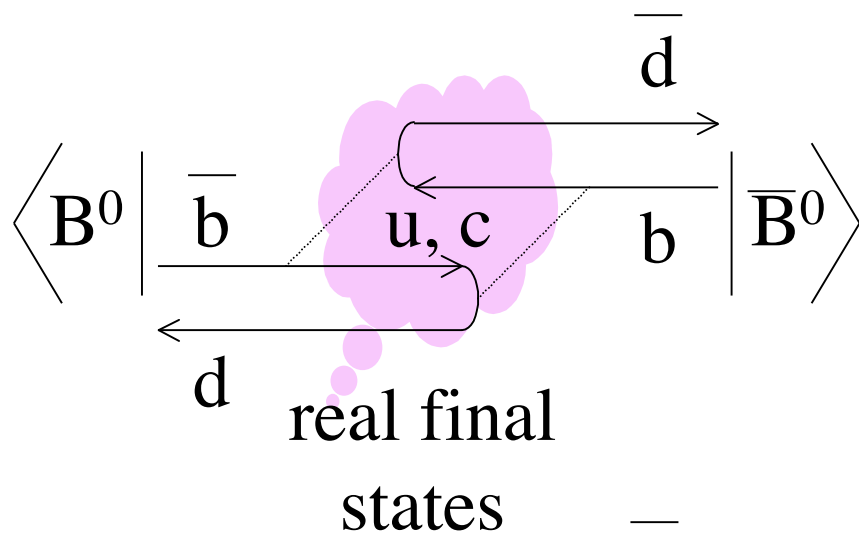
$\mu f_B |V_{ub}|$

decay constant

$$\Gamma_{B \rightarrow \ell \bar{\nu}} \approx f_B^2 |V_{ub}|^2$$

B- \bar{B} oscillation absorptive part: Γ_{12}

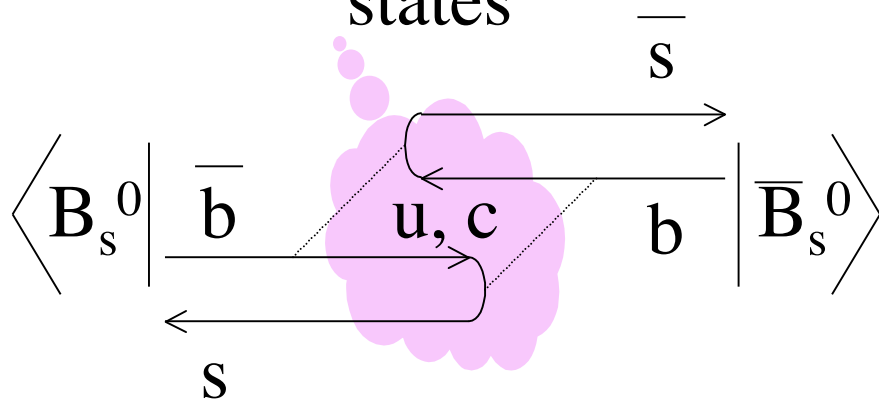
$$\Gamma\Gamma = 2 |\Gamma_{12}|$$



$$\frac{\Gamma\Gamma}{\Gamma m} = \frac{3\Gamma m_b^2}{2m_W^2 S(x_t)} \approx 5 \times 10^3$$

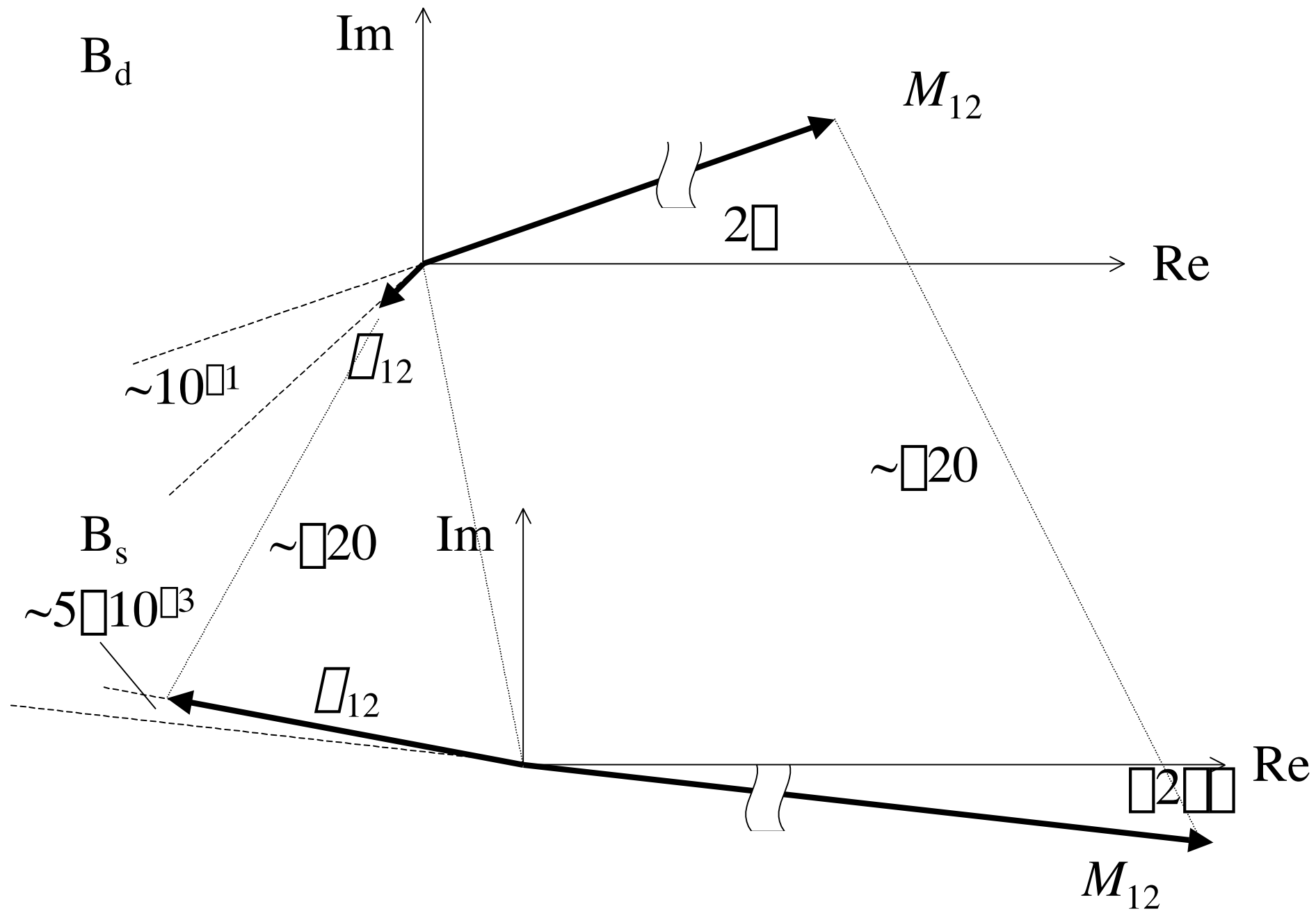
$$x_d = \Gamma m / \Gamma \approx 0.73 \text{ (measurements)}$$

$$x_s = 20 \text{ (SM)} > 16 \text{ (measurements)}$$



$$\frac{\Gamma\Gamma}{\Gamma} = \begin{cases} \sim 10^3: B_d \\ \sim 10^1: B_s \end{cases}$$

$$\arg \frac{M_{12}}{\Gamma_{12}} = \Gamma + \frac{8m_c^2}{3m_b^2} \Gamma\Gamma \begin{cases} \frac{\Gamma}{(1\Gamma\Gamma)^2 + \Gamma^2}: B_d \approx \Gamma\Gamma 10^1 \\ \Gamma^2: B_s \approx \Gamma + 5\Gamma 10^3 \end{cases}$$



To be more precise...

The mass and decay width differences:

$$\begin{aligned}
 m_{\text{heavy}} - m_{\text{light}} &\equiv \Delta m = 2|M_{12}| \cos(\arg M_{12}/\Gamma_{12}) \\
 \Gamma_{\text{heavy}} - \Gamma_{\text{light}} &\equiv \Delta\Gamma = 2|\Gamma_{12}| \cos(\arg M_{12}/\Gamma_{12})
 \end{aligned}$$

~ 1

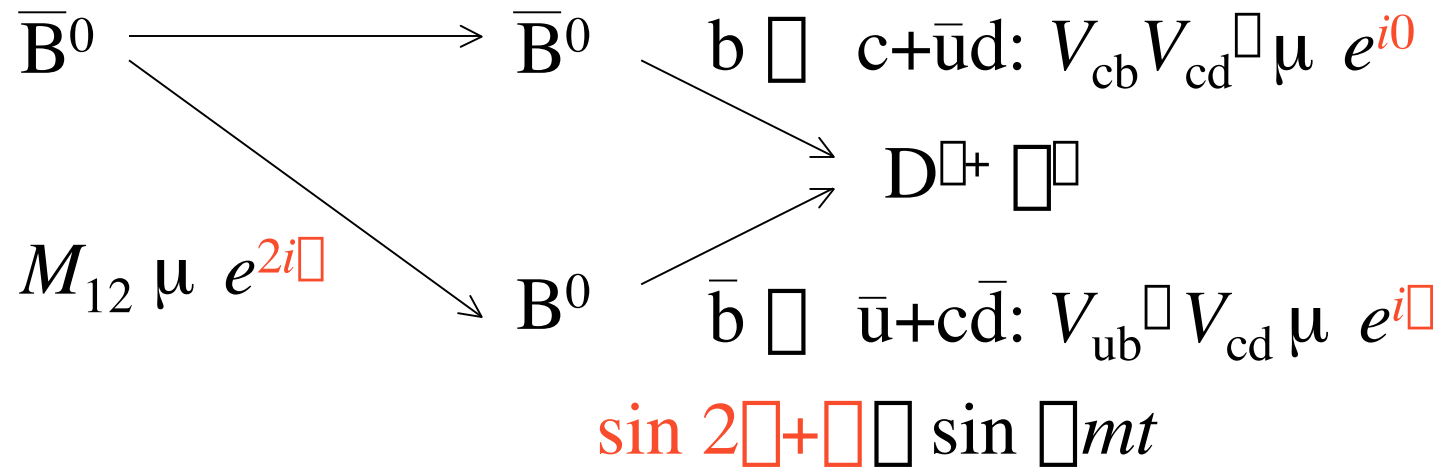
CP violation in the oscillation:

$$d \equiv \text{Im} \frac{\Gamma_{12}}{M_{12}} = \frac{|\Gamma_{12}|}{\Delta m} \sin(\arg M_{12}/\Gamma_{12})$$

5×10^{-3}

$\left\{ \begin{array}{l} \sim 10^{-1} \quad (\text{B}_d) \\ \sim 5 \times 10^{-3} \quad (\text{B}_s) \end{array} \right.$

$$\frac{\overline{\text{B}}_{t=0} - \text{B}_t - \text{B}_{t=0} - \overline{\text{B}}_t}{\overline{\text{B}}_{t=0} + \text{B}_t + \text{B}_{t=0} + \overline{\text{B}}_t} = 2d < 10^{-3}$$



No other diagram contribute to the process.

NB: strong phase

$b \rightarrow c + \bar{u}d$ and $b \rightarrow u + \bar{c}d$

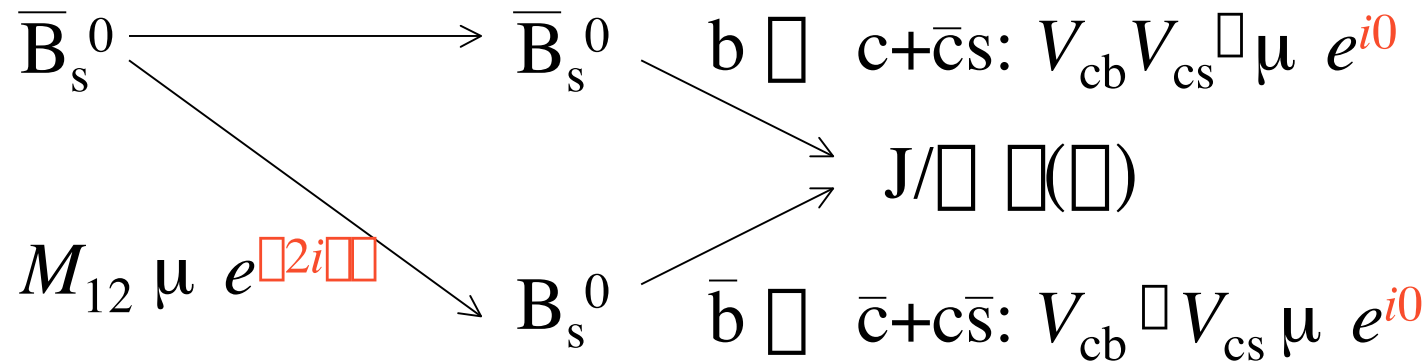
may have a strong phase difference θ !

$\bar{B}^0 \rightarrow D^+ K^-$ and $\bar{B}^0 \rightarrow D^0 K^+$

$2\theta + \theta + \theta$ and $2\theta + \theta$

Strong phase difference can be measured.

Theoretically VERY clean measurement.



$$\sin 2\theta \sin \theta mt$$

$$\text{CP}(J/\psi) = +1, \quad \text{CP}(\psi) = +1$$

$$\text{CP}(J/\psi \psi) = (\pm 1)^{L_{J/\psi-\psi}}$$

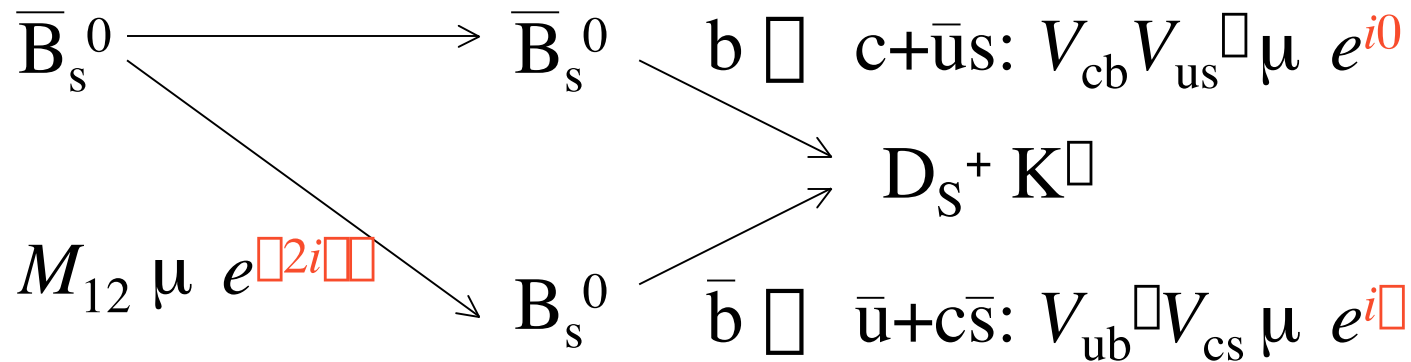
$$S = S_{J/\psi} + S_{\psi}: 0, 1, 2, \quad J = S + L = 0$$

$$L_{J/\psi-\psi} = 0, 1 \text{ or } 2$$

ψ fraction of $L_{J/\psi-\psi} = 1$ needed: angular analysis

$$\text{CP}(J/\psi) = +1, \quad \text{CP}(\psi) = \pm 1$$

$$\text{CP}(J/\psi \psi) = \pm (\pm 1)^{L_{J/\psi-\psi}0} = +1$$



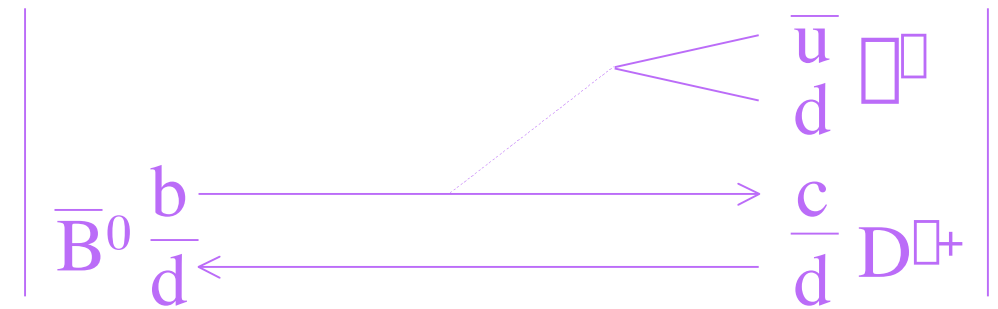
$$\sin 2\phi \sin \phi mt$$

Strong phase: same as the $B_d \rightarrow D^0$ case,
i.e. measurable!



$\phi \approx 0.1 \Rightarrow \cos 2\phi \sinh \phi t$ measurable.

Theoretically VERY clean measurement.

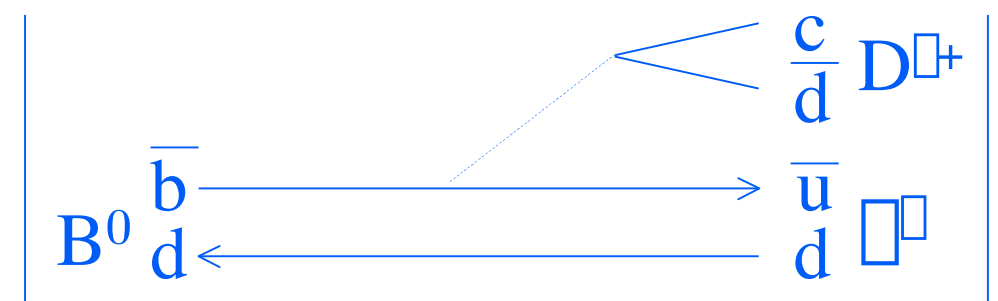


$B_d \rightarrow D^+ D_s^-$

$$\left| \frac{A^2}{A^4 \sqrt{1^2 + 1^2}} \right| \approx 1/0.02$$

small interference term ☹️

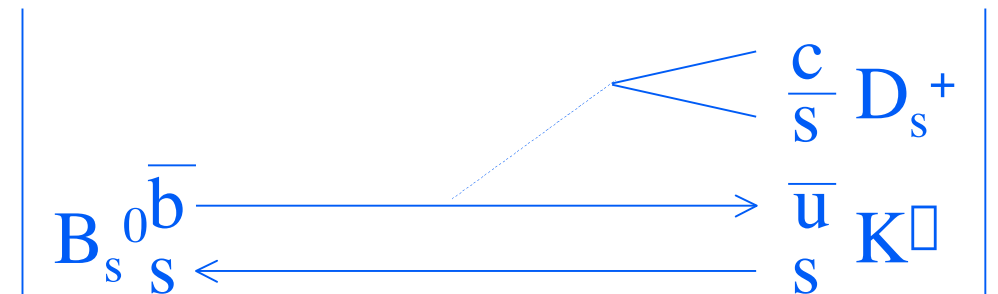
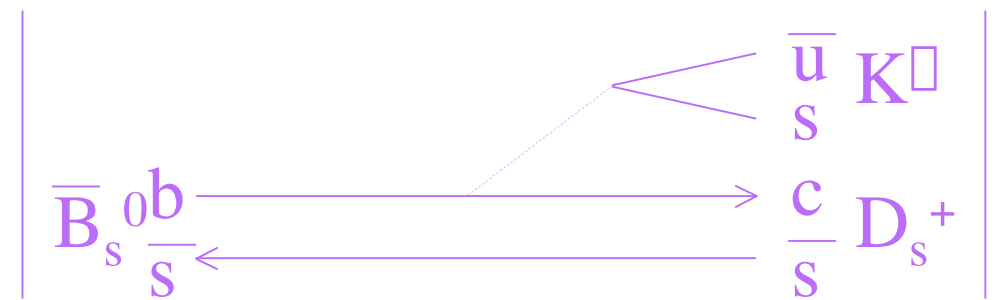
$$\frac{\Gamma_m}{\Gamma} \approx 1 \text{ 😊}$$



$B_s \rightarrow D_s^+ K^-$

$$\left| \frac{A^3}{A^3 \sqrt{1^2 + 1^2}} \right| \approx 0.4$$

large interference term 😊



$$\frac{\Gamma_m}{\Gamma} \approx 20 \text{ ☹️}$$

3) CP landscape

CKM constraints now

+ limit on

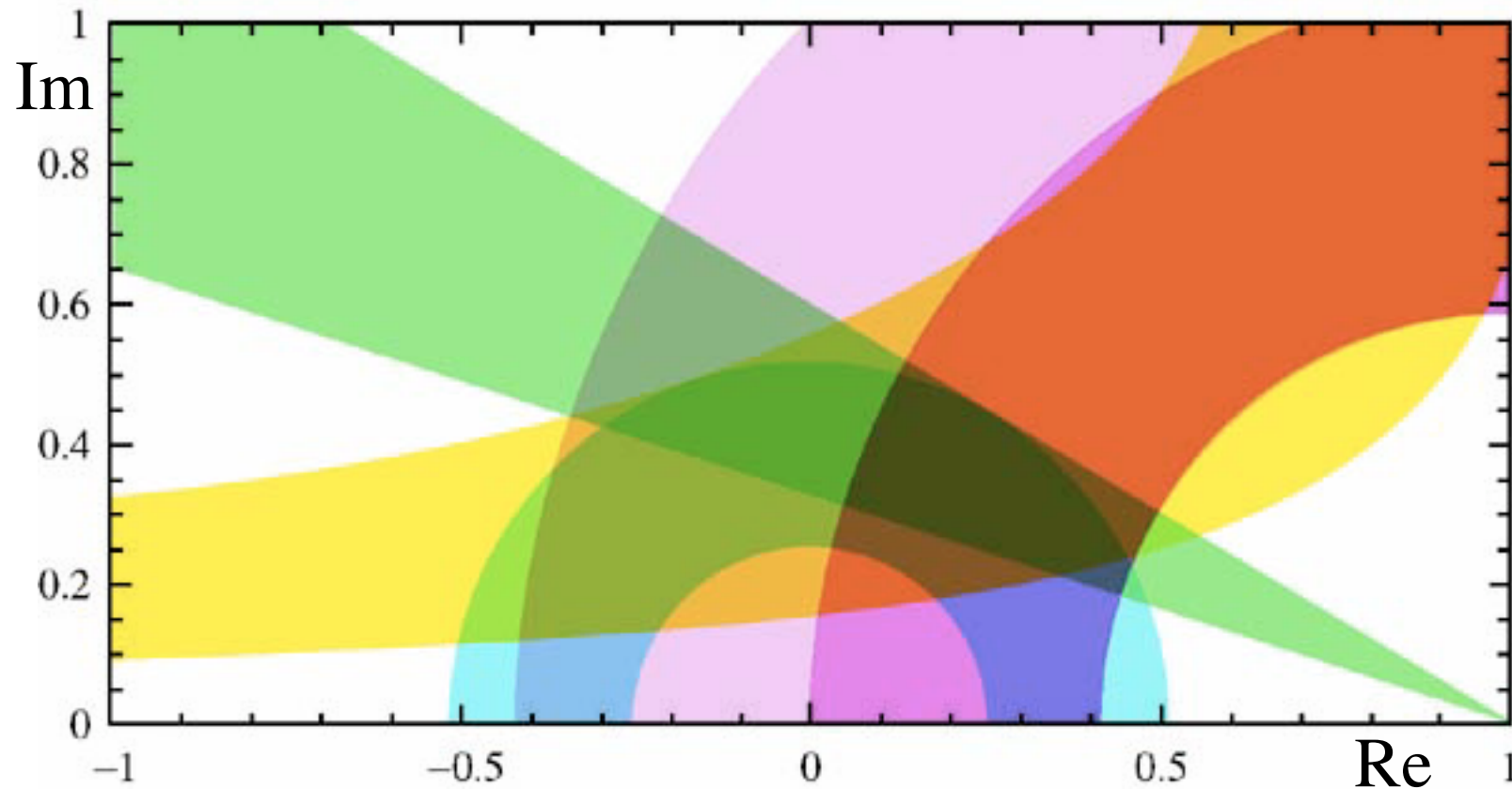
$|V_{ub}/V_{cb}|$

$\Delta m(B_d)$

$\Delta m(B_s)$

$\Delta(K)$

$\sin 2\beta$



Perfectly consistent within the statistics!

CKM fit without $\sin 2\beta$ measurements

$\beta = 0.23 \pm 0.?$, $\beta = 0.33 \pm 0.?$ “?” is somewhat “theological”

$\sin 2\beta = 0.50 \text{ } \square \text{ } 0.86$ CL > 32% A.Höcker et al.
 0.698 ± 0.066 A.Stocchi et al.

$\beta = 23.2^\circ$ without β^2 correction, 22.5° with β^2 correction

$\sin 2\beta$ differs by 0.02

Kaon system:

$$|\epsilon| = (2.271 \pm 0.017)10^{-3}$$

known to better than 1%

theoretical uncertainties $\sim 10\%$

$$\text{Re}(\epsilon/\epsilon') = (17.3 \pm 2.3) \times 10^{-4}$$

known to 13%

theoretical uncertainties $>100\%$

□ No precision test is possible

B-meson system:

OPAL (1998)

$$\sin 2\beta = 3.2^{+1.9}_{-2.1}$$

CDF (2000)

$$= 0.79^{+0.41}_{-0.44}$$

ALEPH (2000)

$$= 0.84^{+0.84}_{-1.05}$$

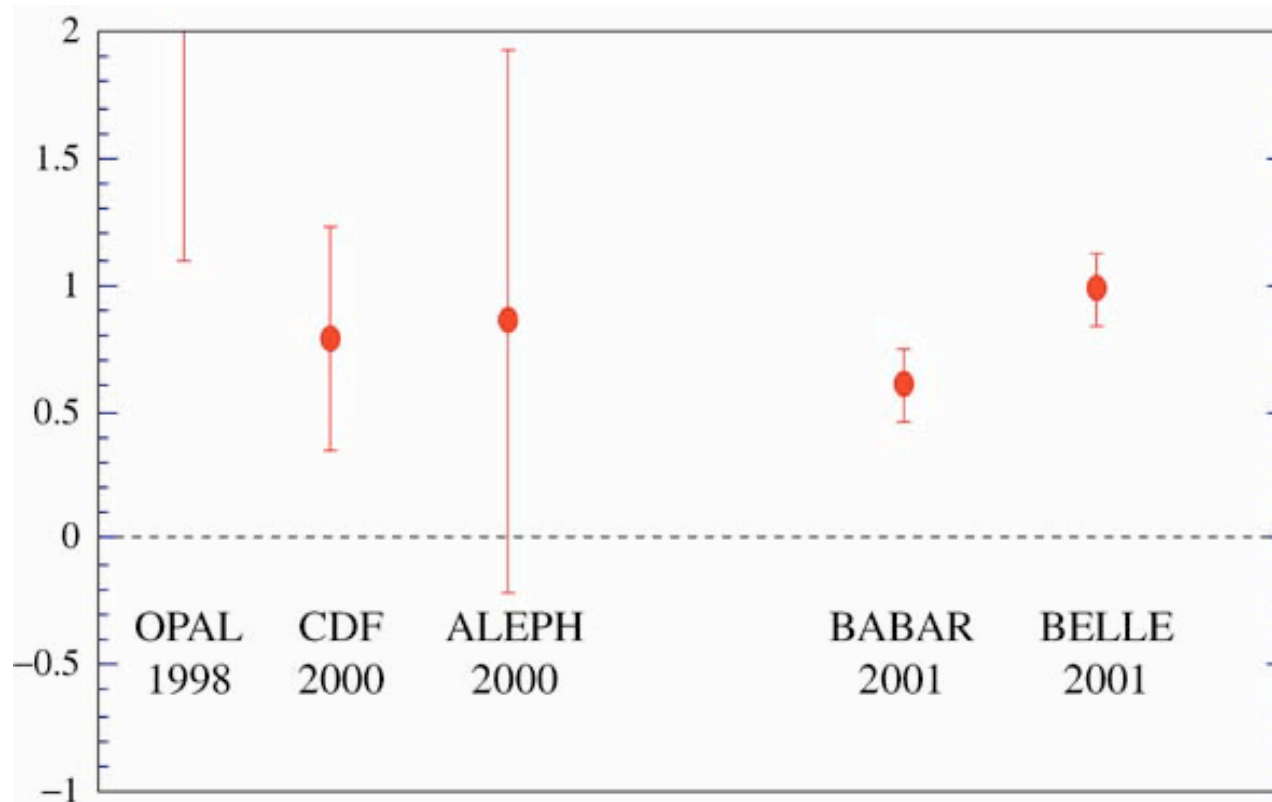
BABAR (2001)

$$= 0.59 \pm 0.15$$

BELLE (2001)

$$= 0.99 \pm 0.15$$

} $0.79 \pm 0.17: \chi^2 = 3.56$
(error scaled)



Major improvements expected...

Theory:

A better understanding of $Bf_B^2 \propto |V_{td}/V_{cb}|^2$ from $\Delta m(B_d)$
 $\sqrt{(1 - \lambda)^2 + \lambda^2}$

Theory and more data

A better understanding of $b \rightarrow u, c \rightarrow \ell$ decays: $|V_{ub}/V_{cb}|$
 $\sqrt{\lambda^2 + \lambda^2}$

More data

An order of magnitude smaller error on $\sin 2\beta$:
before hitting the penguin pollution

$\sqrt{(1 - \lambda)^2 + \lambda^2}$ from the $\Delta m(B_d)/\Delta m(B_s)$ measurement

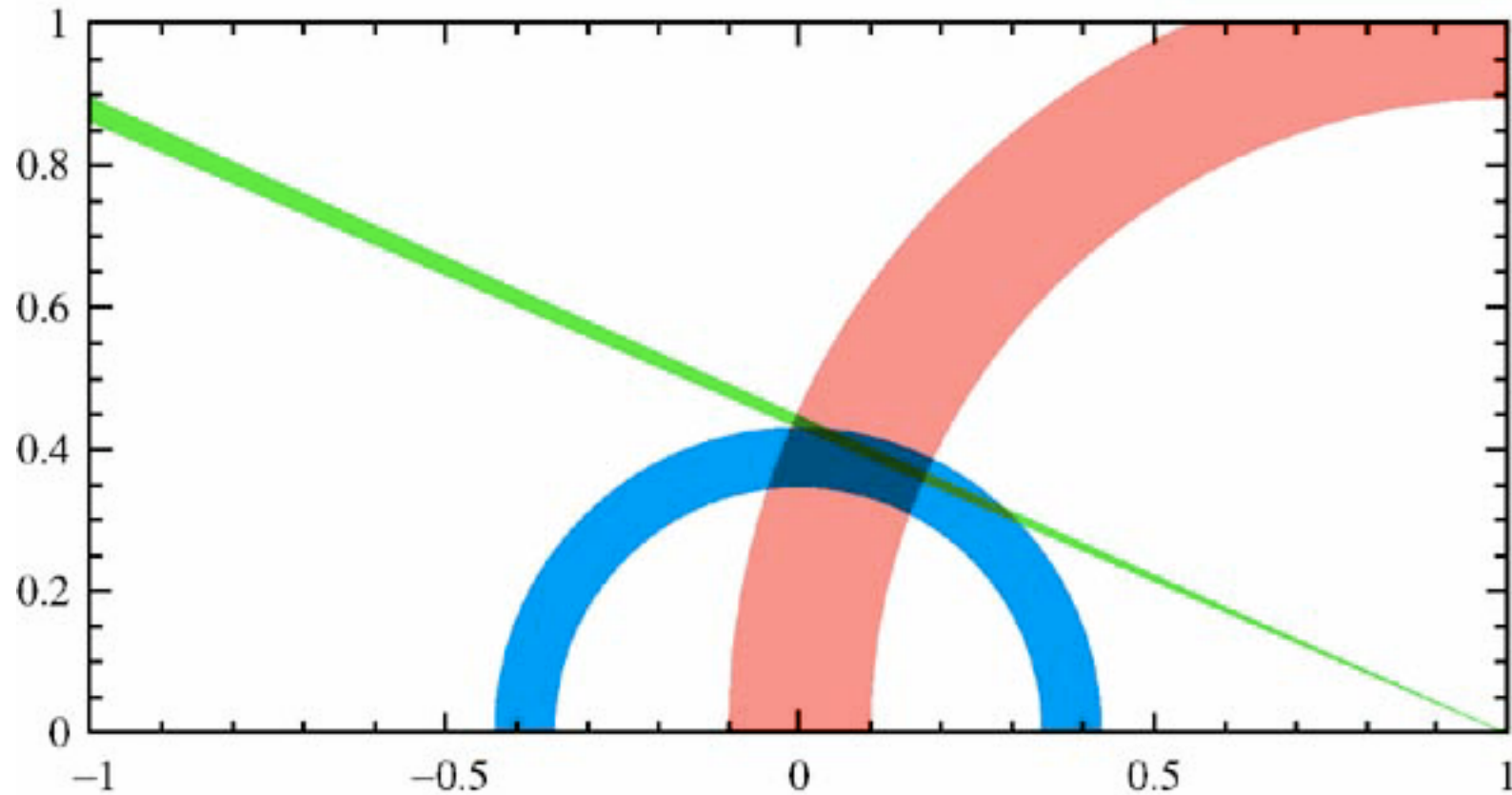
2001

$\sqrt{x^2 + y^2}$: improved by a factor of 3

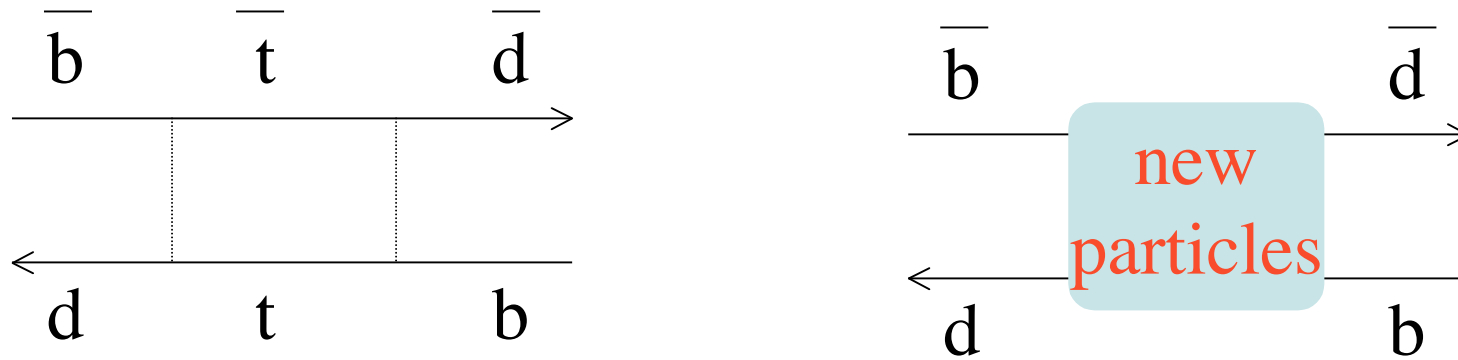
~2007

$\sqrt{(1 - |x|)^2 + y^2}$: improved by a factor of 4

$\sin 2\theta$: ± 0.01



4) Possible effect of new physics



$$M_{12} = M_{12}^{\text{SM}} + M_{12}^{\text{NP}}$$

$$\varphi_m \not\leftrightarrow |V_{td}|^2$$

$$\arg M_{12} = 2\varphi + \varphi_{bd}^{\text{NP}}$$

$$\text{CP asymmetry in } B_d \not\leftrightarrow J/\psi K_S \not\leftrightarrow \sin 2\varphi$$

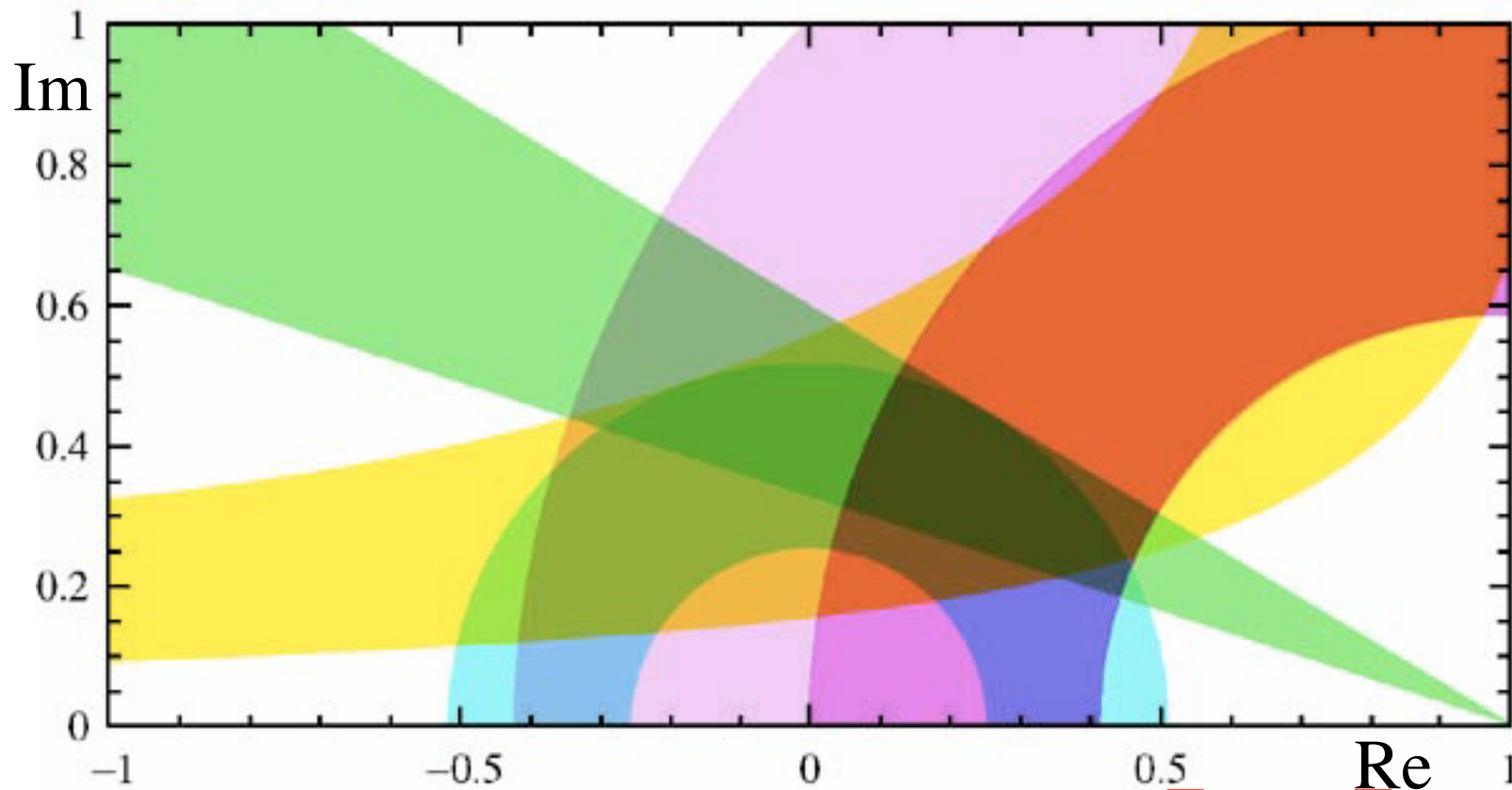
CP asymmetry in $B_d \rightarrow J/\psi K_S \leftrightarrow \sin 2\beta$

New physics in $K-\bar{K}$ oscillation

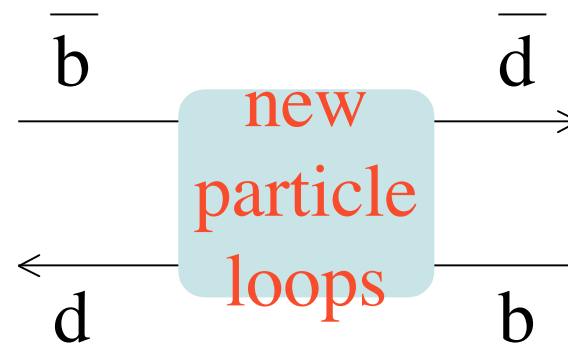
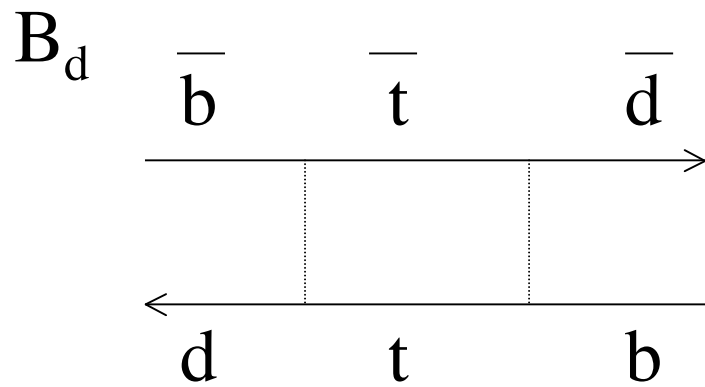
$\Delta m \leftrightarrow |V_{td}|^2$

+ limit on

$|V_{ub}/V_{cb}|$ $\Delta m(B_d)$ $\Delta m(B_s)$ $\Delta(K)$ $\sin 2\beta$



Not enough constraint to determine β and α .



$$\arg M_{12} = 2\phi + \phi_{bd}^{NP}$$

$$\text{CP asymmetry in } B_d \rightarrow J/\psi K_S \iff \sin(2\phi + \phi_{bd}^{NP})$$

$$\text{CP asymmetry in } B_d \rightarrow D^+ D^- \iff \sin(2\phi + \theta + \phi_{bd}^{NP})$$

B_s

$$\arg M_{12} = 2\phi + \phi_{bs}^{NP}$$

$$\text{CP asymmetry in } B_s \rightarrow J/\psi \phi \iff \sin(2\phi + \phi_{bs}^{NP})$$

$$\text{CP asymmetry in } B_s \rightarrow D_s^+ K^- \iff \sin(2\phi + \theta + \phi_{bs}^{NP})$$

θ can be determined!

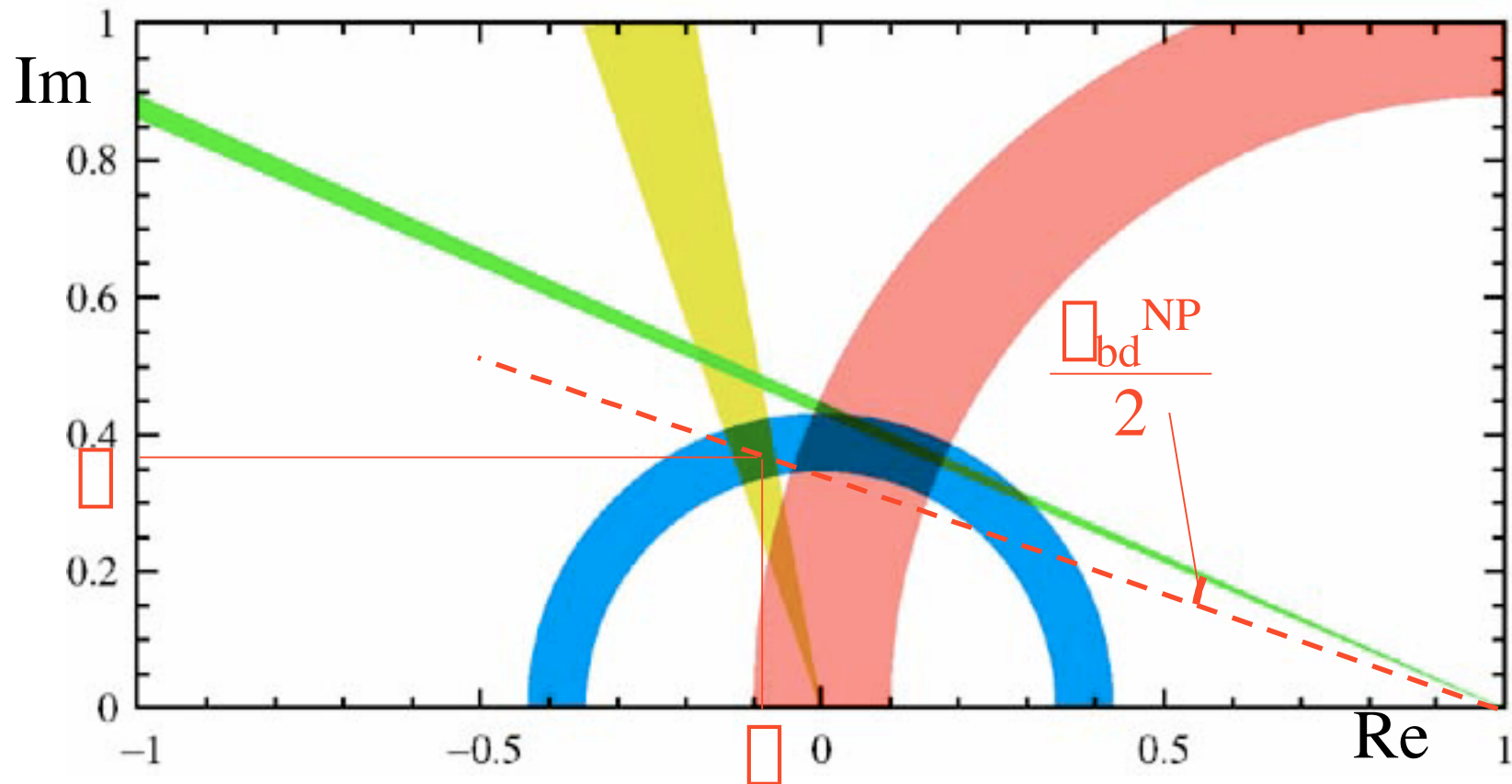
~2007

clean measurements

BABAR, BELLE: $|V_{cb}|$, $|V_{ub}|$, φ_m^d , $\sin 2\varphi$,

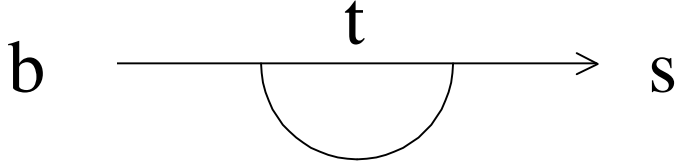
CDF (D0): φ_m^s , $\sin 2\varphi$,

Clean measurements of $\varphi \pm 5^\circ$

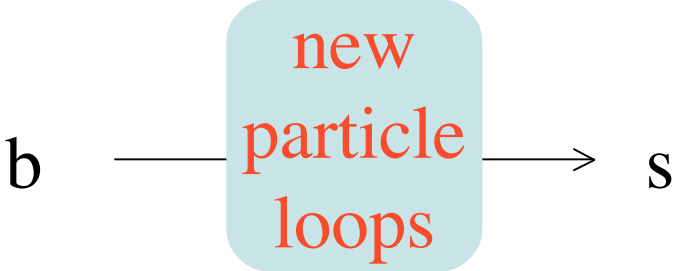


Dedicated B experiments at LHC (and Tevatron)!!!

Penguin could be affected by new physics?

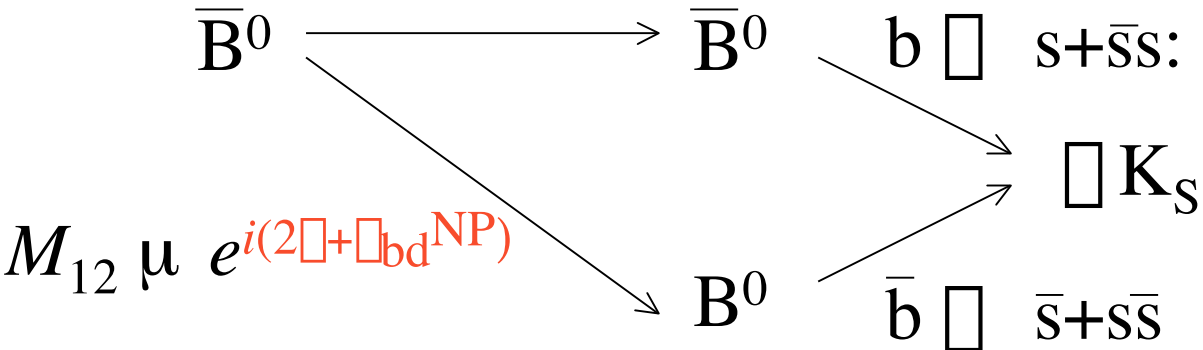


$$V_{bt} V_{ts}^* \mu e^{i\phi}$$



$$A_{bs}^{NP} \mu e^{i\phi_{pbs}^{NP}}$$

$B_d \rightarrow K_s$



$B_s \rightarrow K_s$
as well

$$\sin \phi \approx \sin \phi_{mt}$$

Future experiments need to have:

High statistics B_d and B_s .

Trigger sensitive to the final state with leptons
and with **only hadrons**.

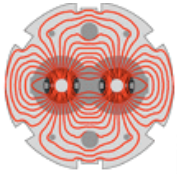
Good proper time resolution for measuring the CP violating
oscillation amplitudes of the B_s meson.

Good $\pi/K/\pi/e$ separation to reduce the background from
both combinatorics and **other B meson decays**.
-kaon identification also useful for the flavour tag-

Good momentum and vertex resolutions
to reduce background.

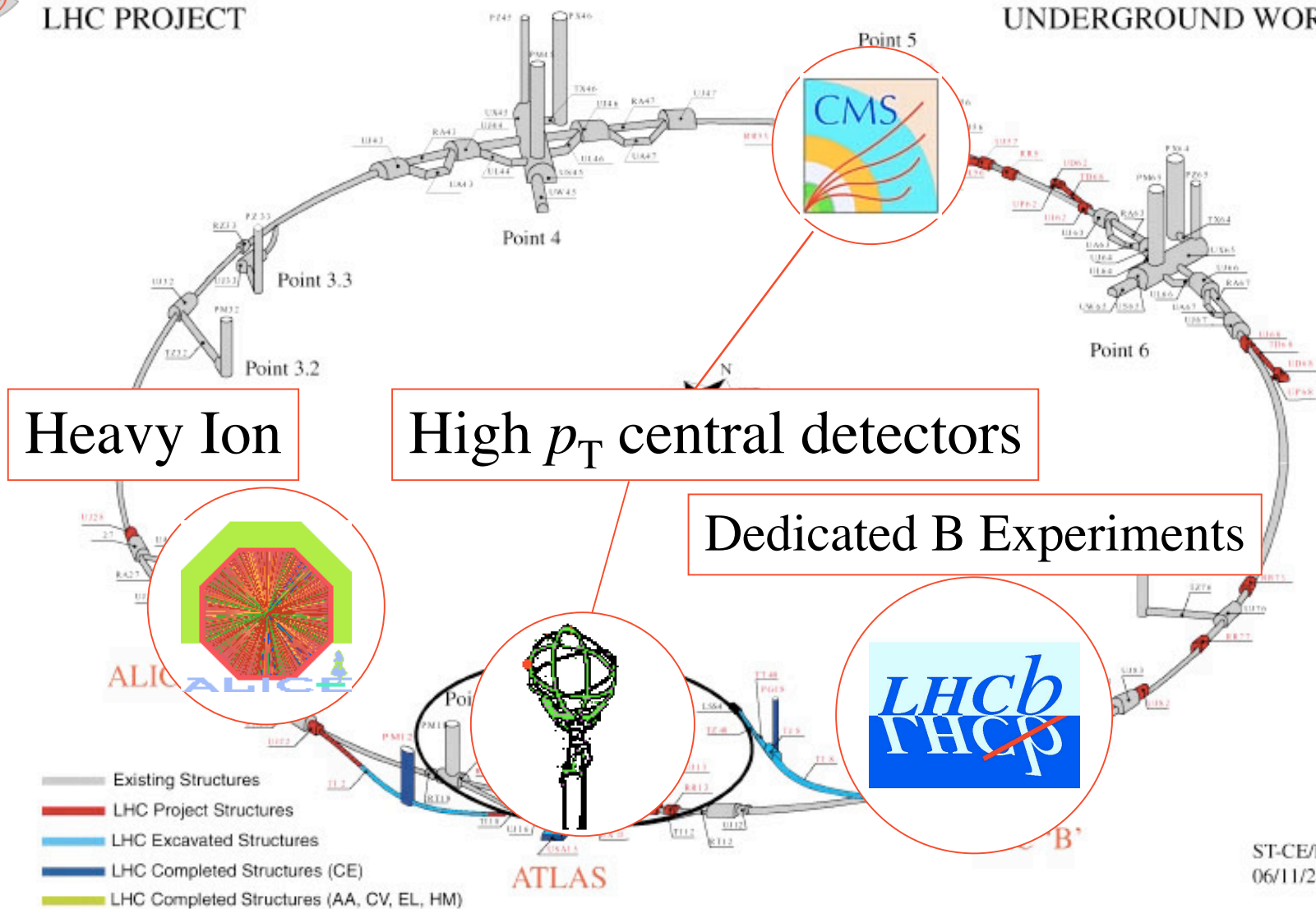
5) LHC experimental conditions





LHC PROJECT

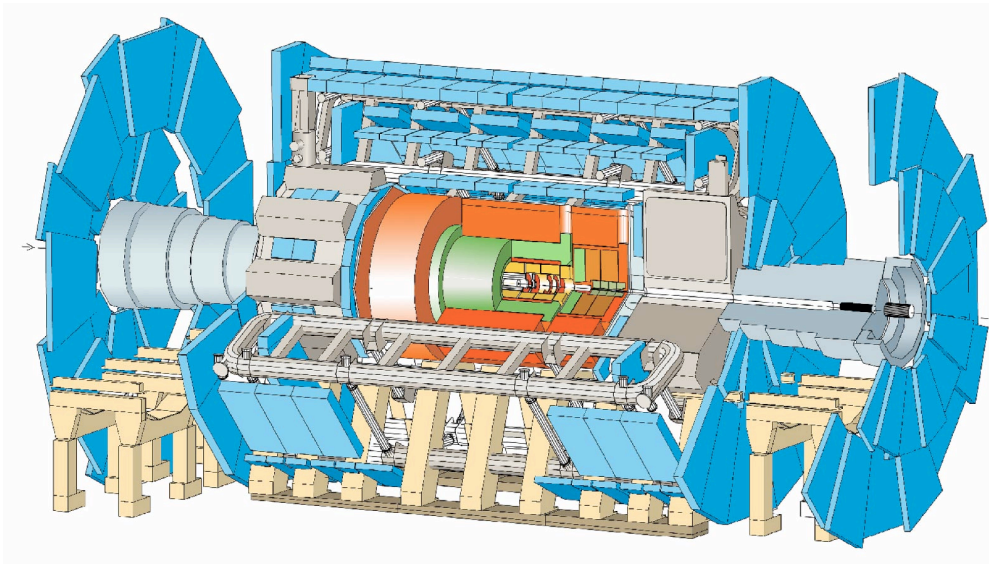
UNDERGROUND WORKS



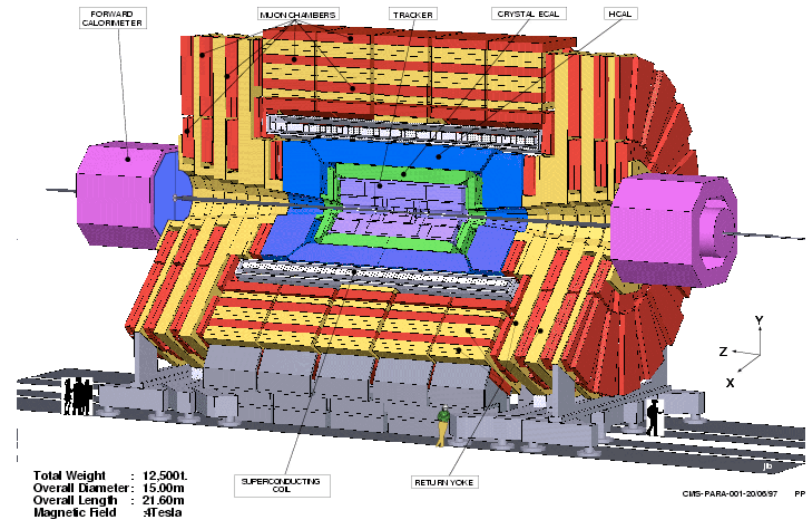
Heavy Ion

High p_T central detectors

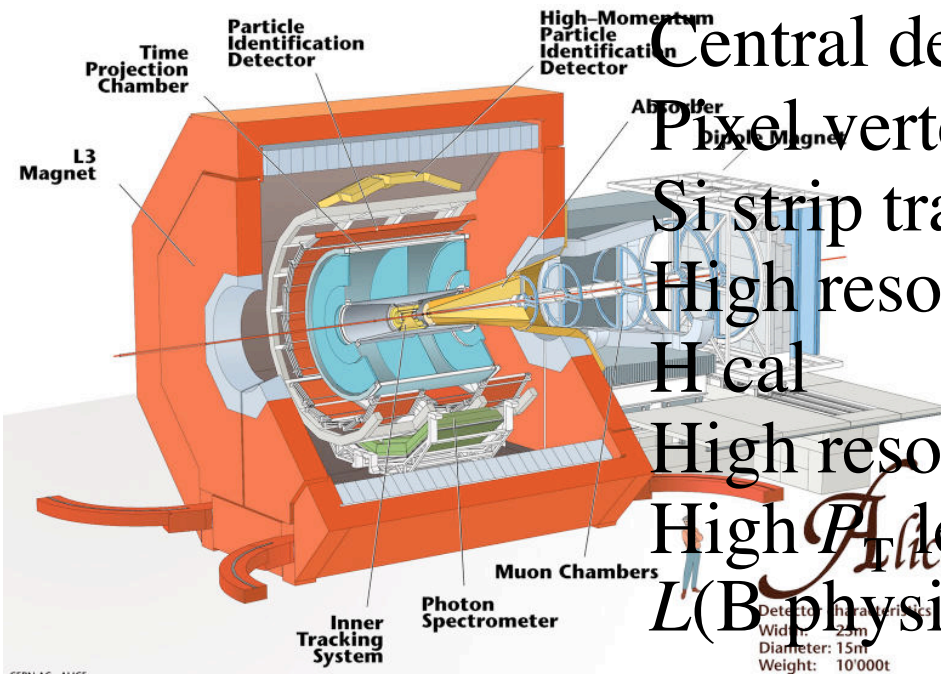
Dedicated B Experiments



ATLAS



CMS



Central detector: $|\eta| < 2.5$

Pixel vertex detector

Si strip tracker

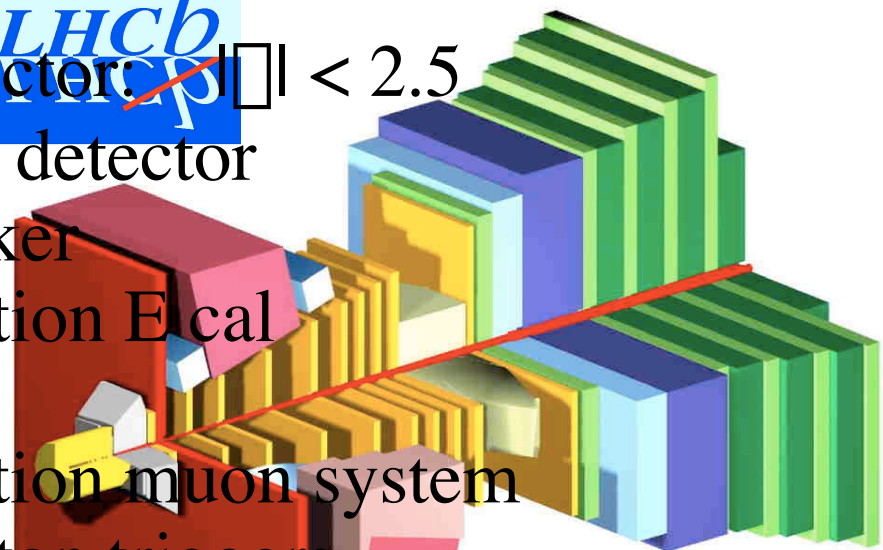
High resolution Ecal

H cal

High resolution muon system

High P_T lepton triggers

$L(B \text{ physics}) 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$



LHC will become operational in ~2006

with a starting luminosity of 10^{32} to 10^{33} $\text{cm}^{-2}\text{s}^{-1}$
(design luminosity $> 10^{34}$ $\text{cm}^{-2}\text{s}^{-1}$)

$\sigma_{b\bar{b}}$ expected in pp collisions at $\sqrt{s} = 14$ TeV: $500 \mu\text{b}$
 5×10^{11} to 5×10^{12} $b\bar{b}$ pairs in 1 year (10^7 s)
 \square powerful source of b quarks!

cf.

\circ (4S) B factories: 10^7 $B\text{-}\bar{B}$ /year @ $L = 10^{33}$ $\text{cm}^{-2}\text{s}^{-1}$

LHC could have a big potential for B physics.

However not too easy experimental environment...

$$\frac{\sigma_{b\bar{b}}}{\sigma_{\text{inelastic}}} = 5 \times 10^{-3}$$

cf

° (4S) B factories: 0.2

HERA-B 10^{-6} ☹️

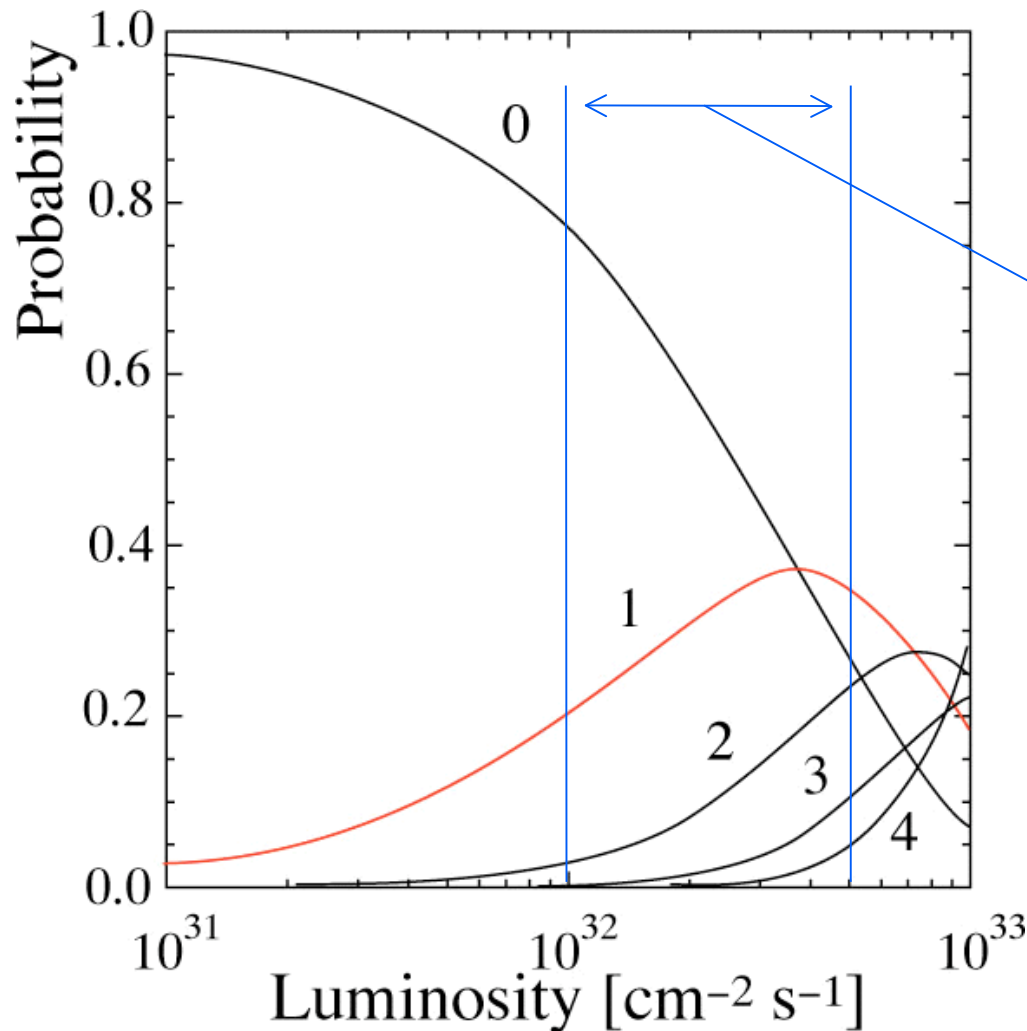
Fixed target charm experiments

$$\frac{\sigma_{c\bar{c}}}{\sigma_{\text{inelastic}}} \approx 10^{-3}$$



Bunch crossing frequency: $f_{pp} = 40$ MHz, i.e. every 25 nsec

Number of $p\bar{p}$ inelastic interactions in one bunch crossing ($\sigma_{\text{inelastic}} = 80$ mb), 0, 1, 2 ...

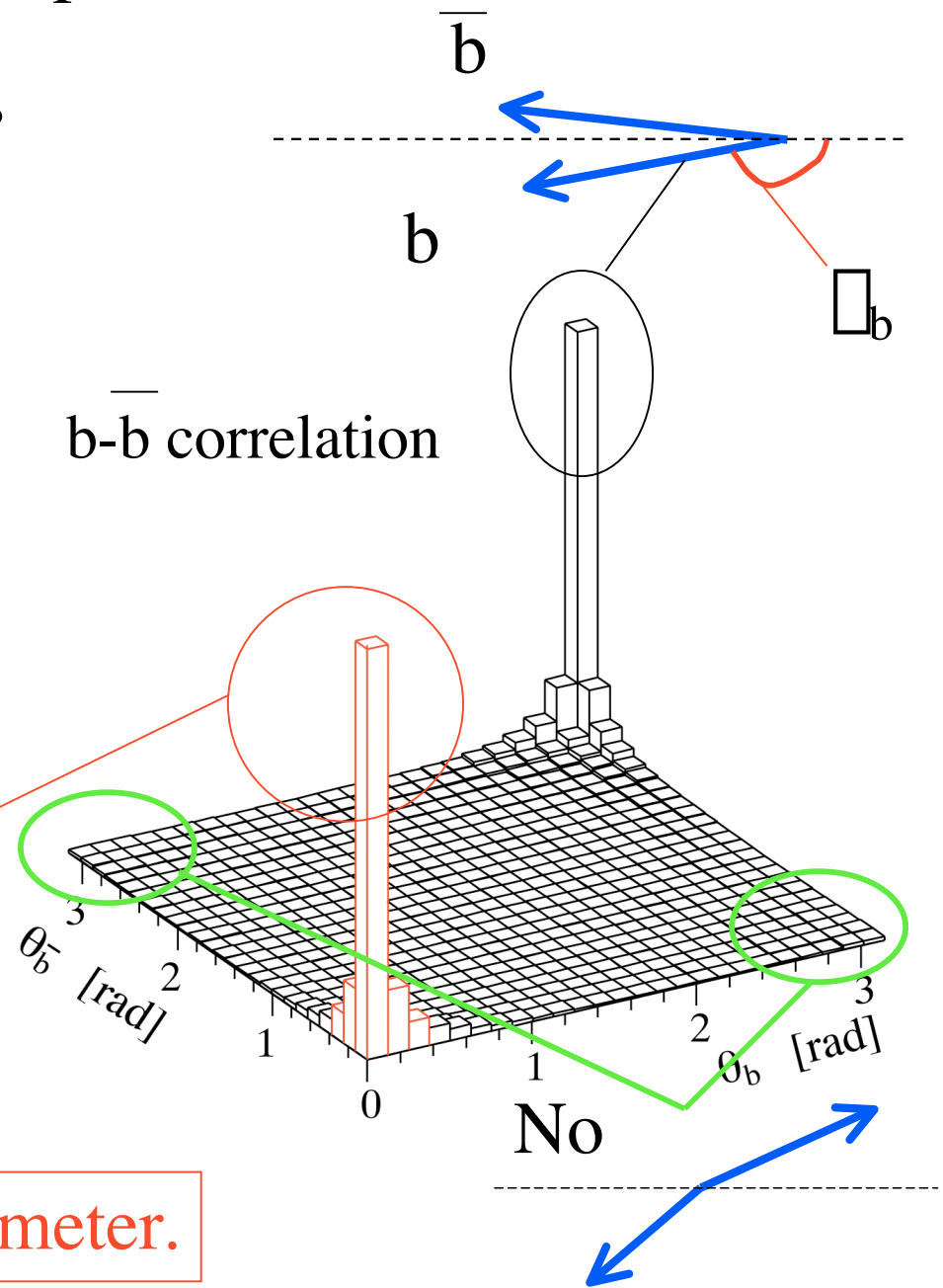
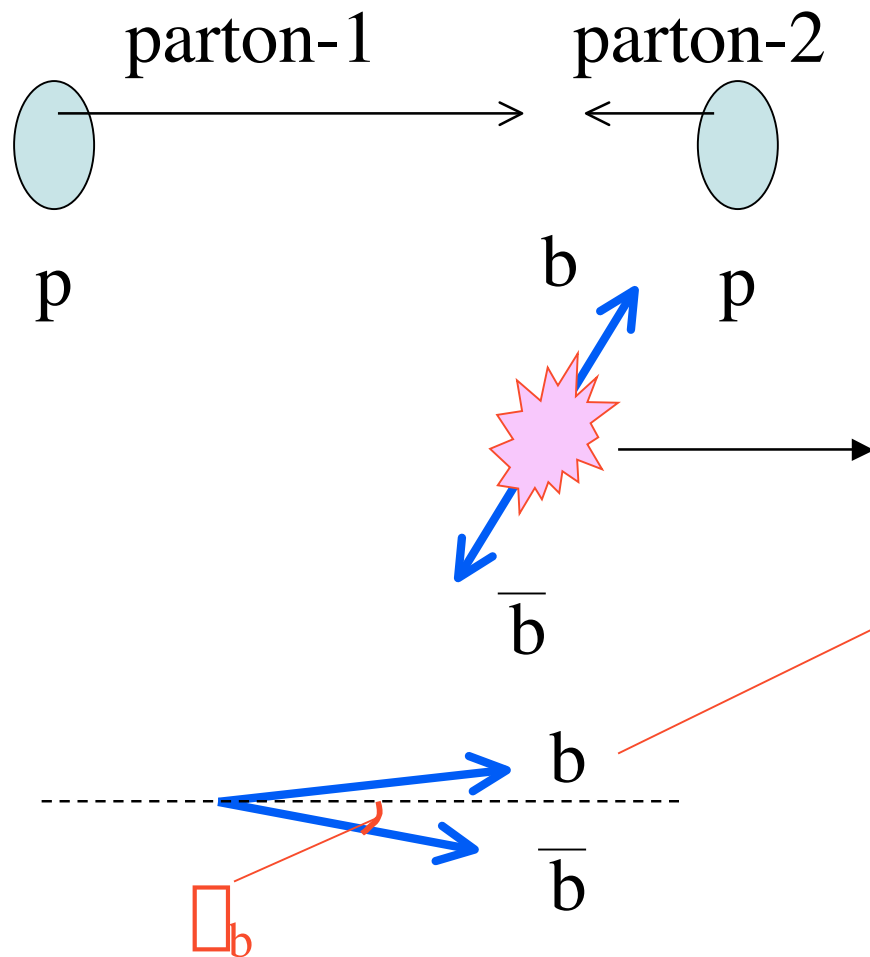


One inelastic interaction per bunch crossing dominates.

- Reconstruction easier (final state and tag)
- Lower radiation level

LHCb experiment: forward spectrometer (10-300 mrad)

$b\bar{b}$ pair production kinematics

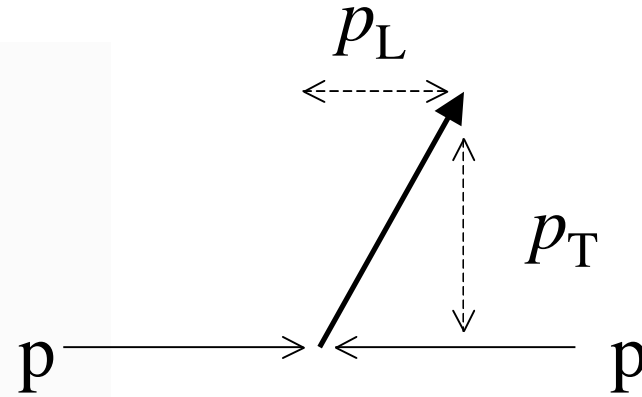
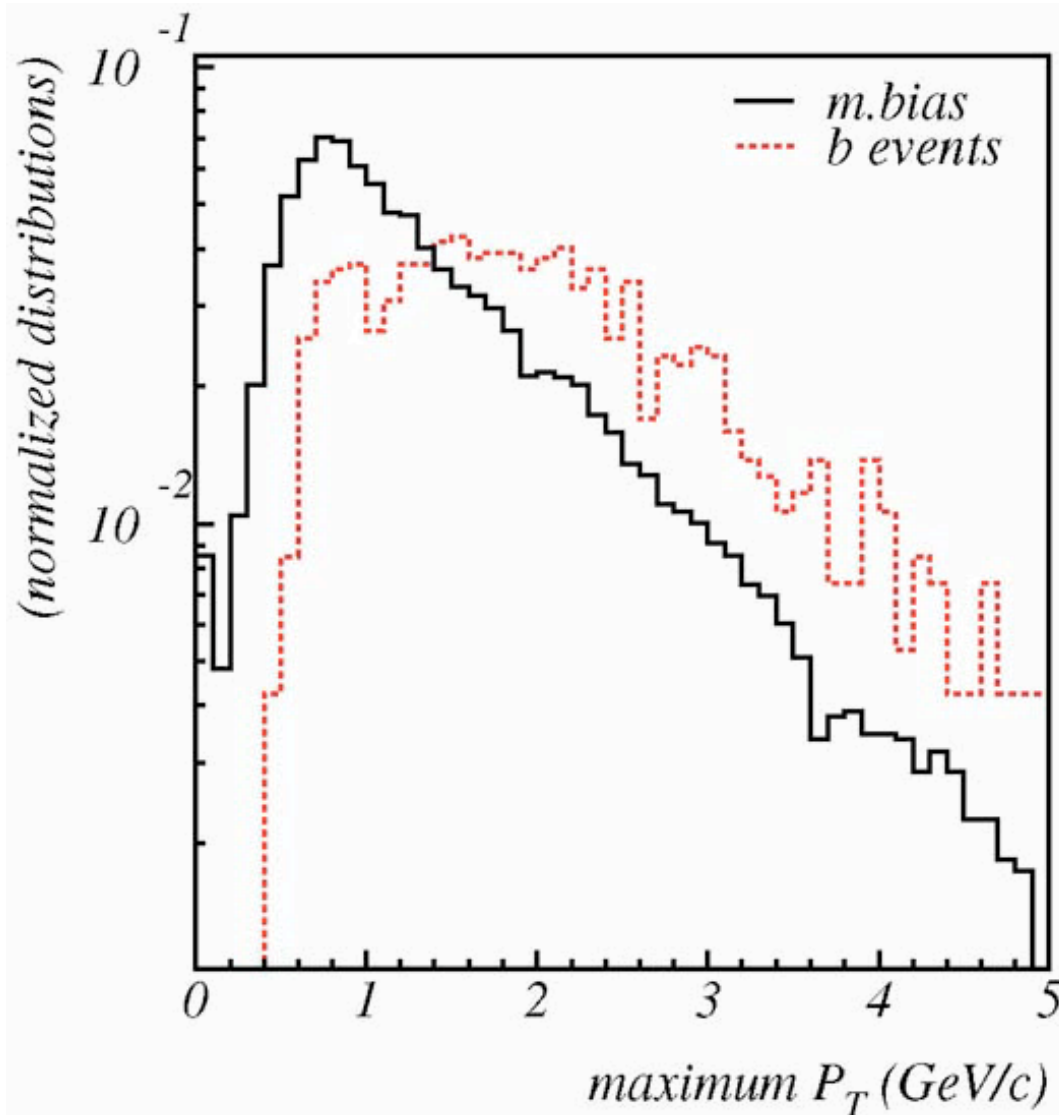


Both b and \bar{b} are in the spectrometer.

f_{pp} @ LHC = 40 MHz

□ simple first level trigger

Single p_T trigger for □, e, h



muon system:
low track density

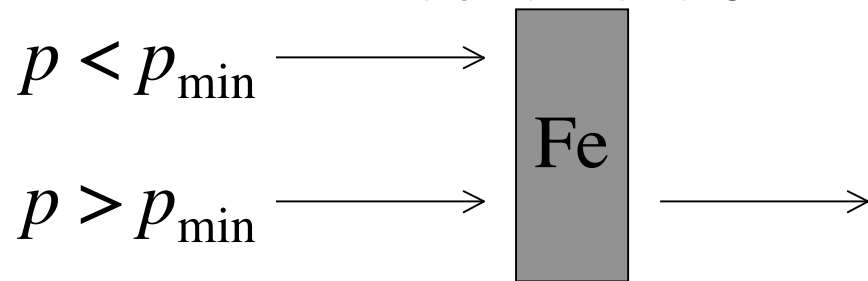
e and h:

calorimeter

E_T measurements

However.... $p > p_{\min}$

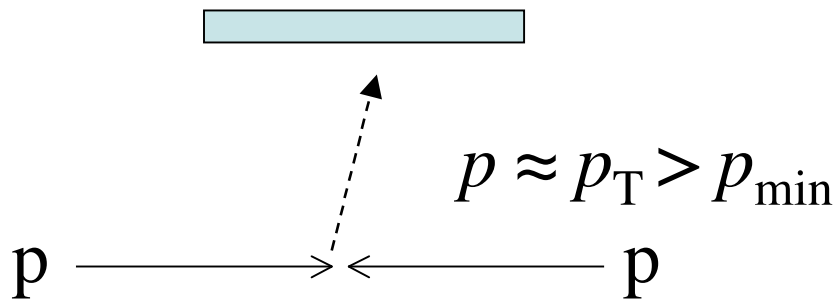
muon:
identification



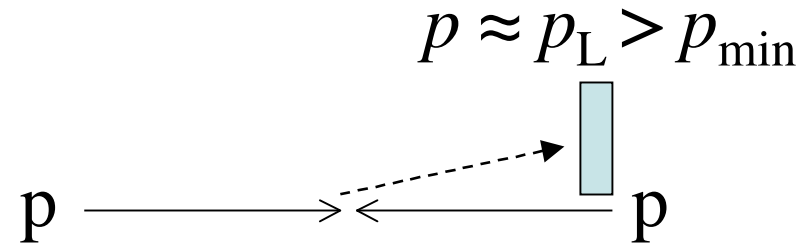
hadron:
energy resolution

$$\frac{\Delta E}{E} \approx \sqrt{70\%}/\sqrt{E}$$

central detector

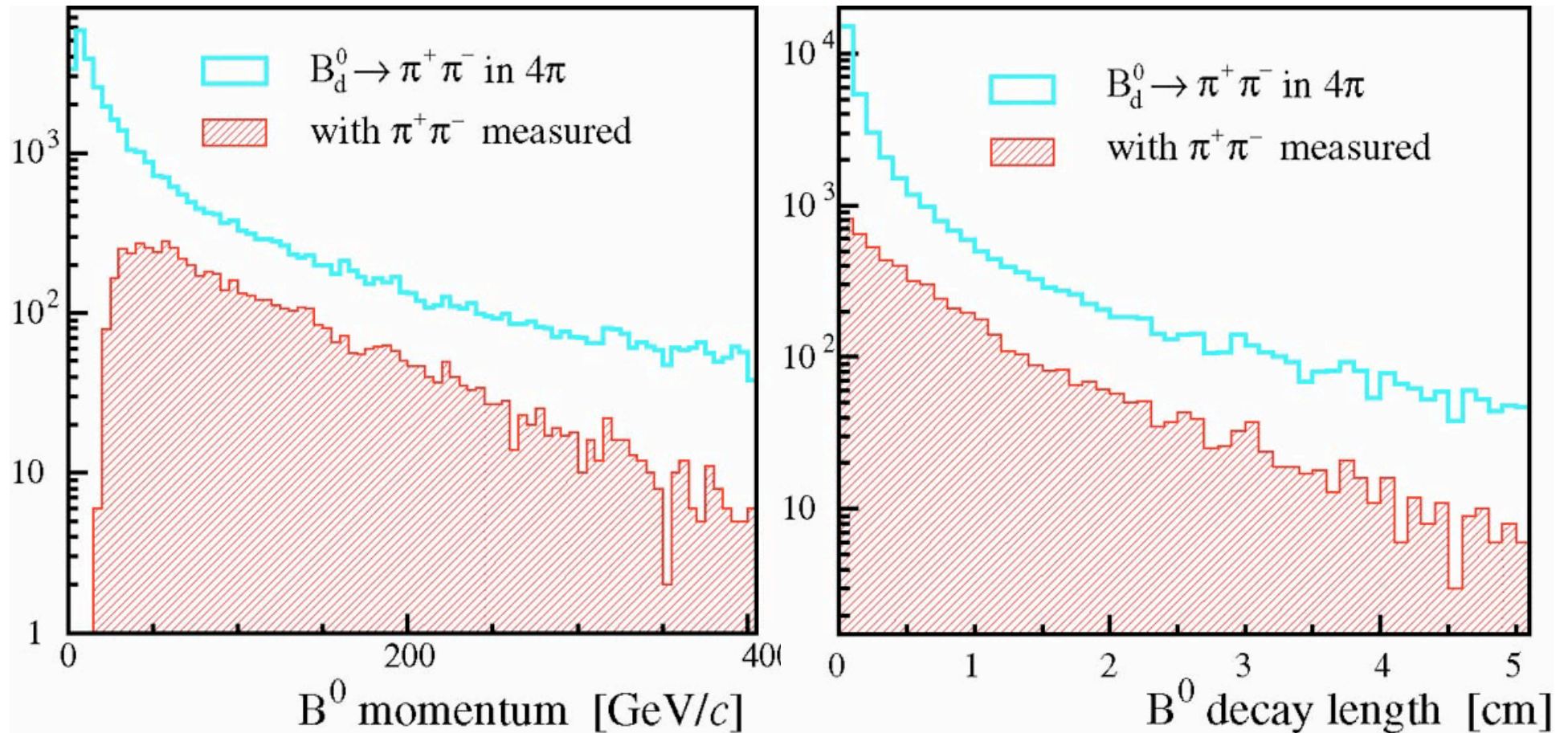


forward detector



p_T threshold can be set low:
□ high b efficiency

Momentum spectrum and decay distance for B mesons



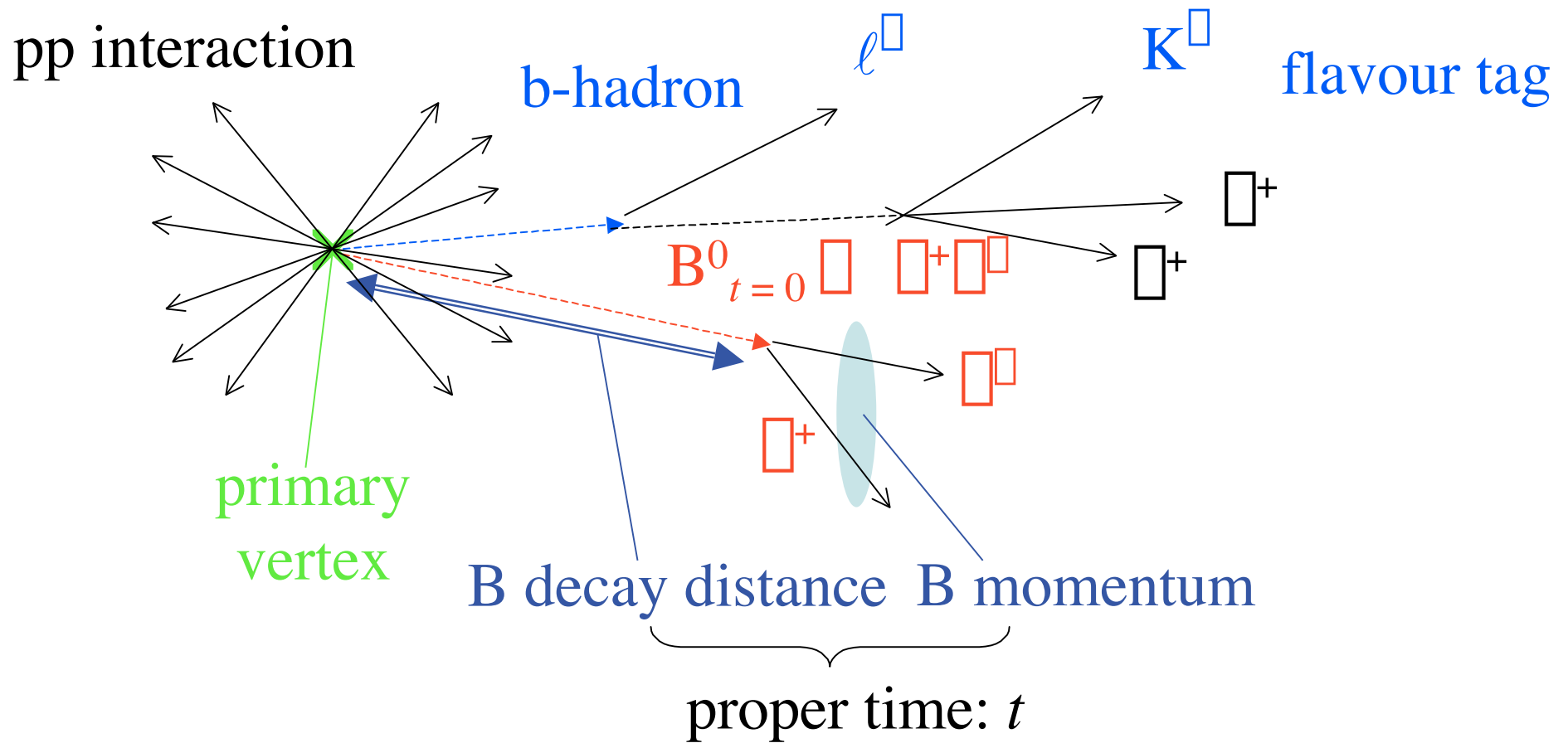
are larger in the forward region.

□ average B decay distance in the detector ~1cm

Good proper time resolution.

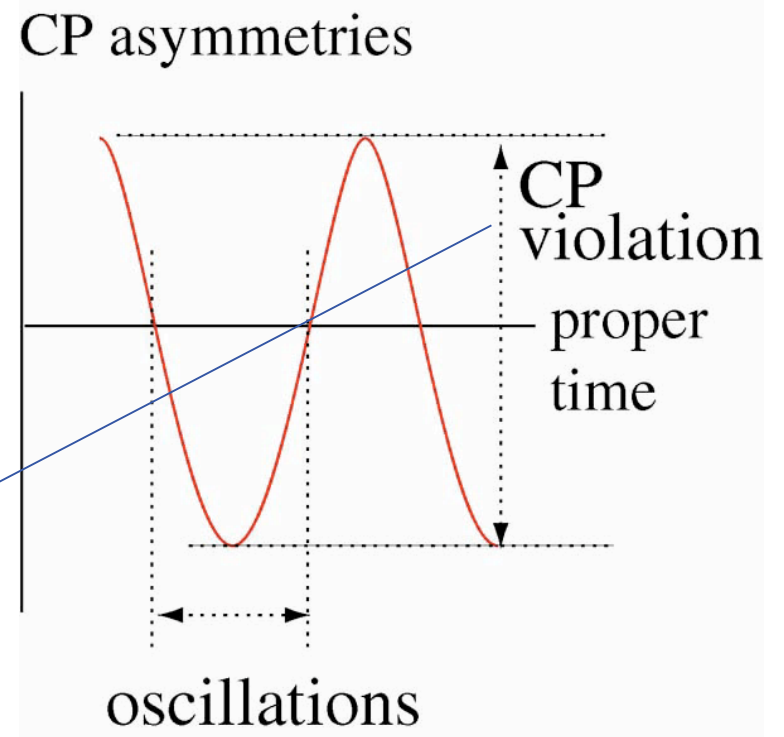
6) Some detector requirements

What do we measure? (an example)



Time dependent CP asymmetry

$$\frac{\overline{B}^0 \rightarrow \pi^+ \pi^0(t) - B^0 \rightarrow \pi^+ \pi^0(t)}{\overline{B}^0 \rightarrow \pi^+ \pi^0(t) + B^0 \rightarrow \pi^+ \pi^0(t)}$$



CP violating oscillation amplitudes are damped by

proper time resolution

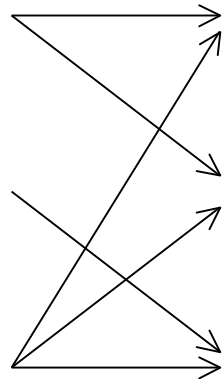
good decay vertex resolution

wrong flavour tag

good momentum resolution

background

good particle identification



7) LHCb Detector



VELO

Inner Tracker

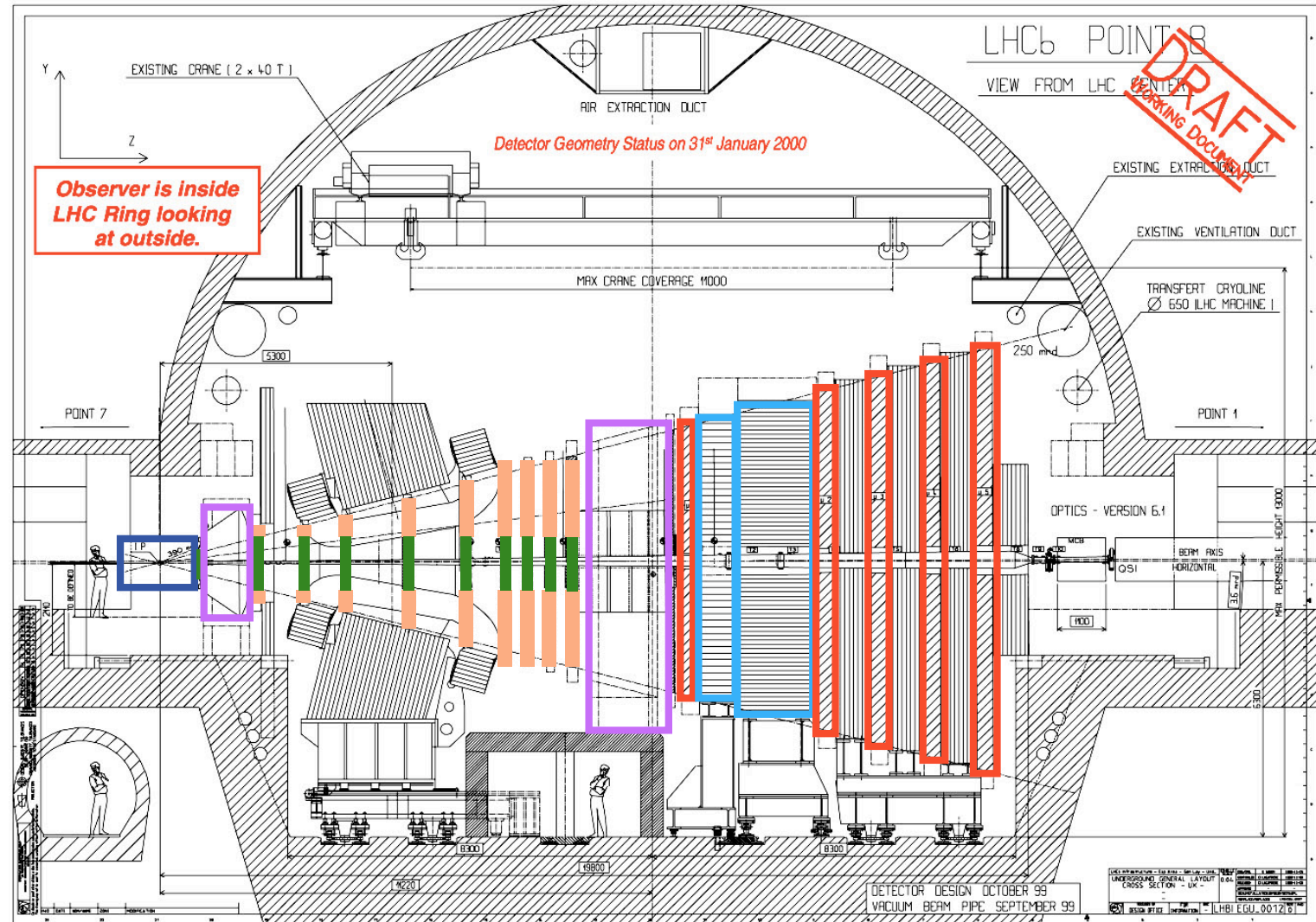
Outer Tracker

RICH-1

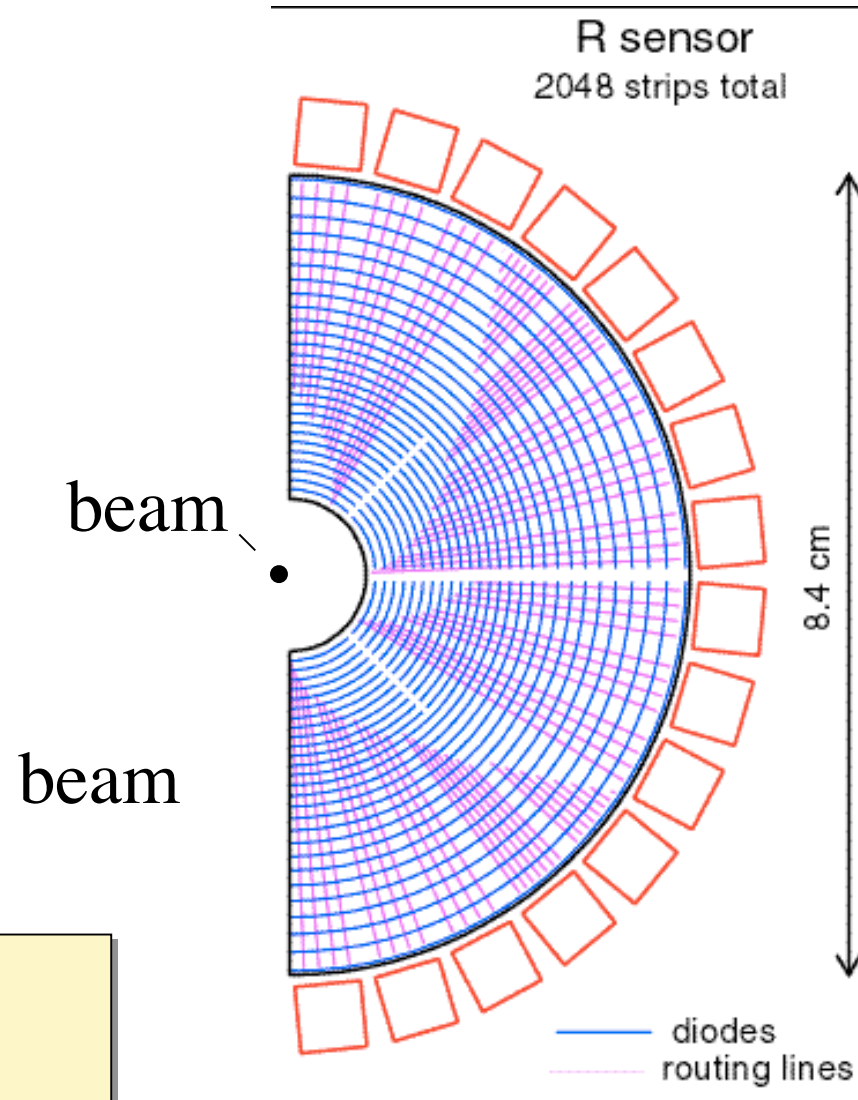
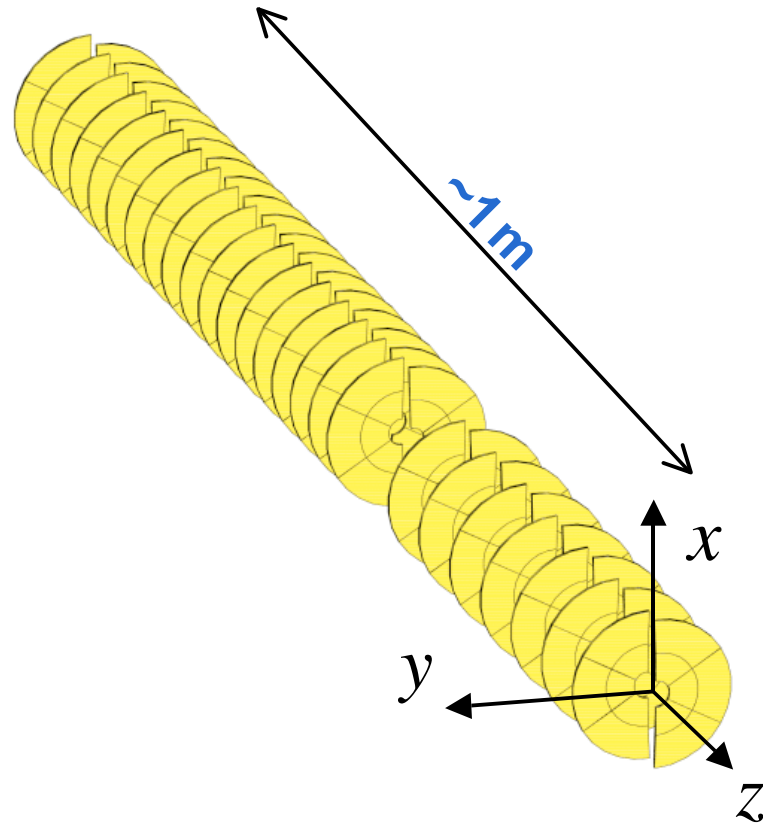
RICH-2

Calorimeters

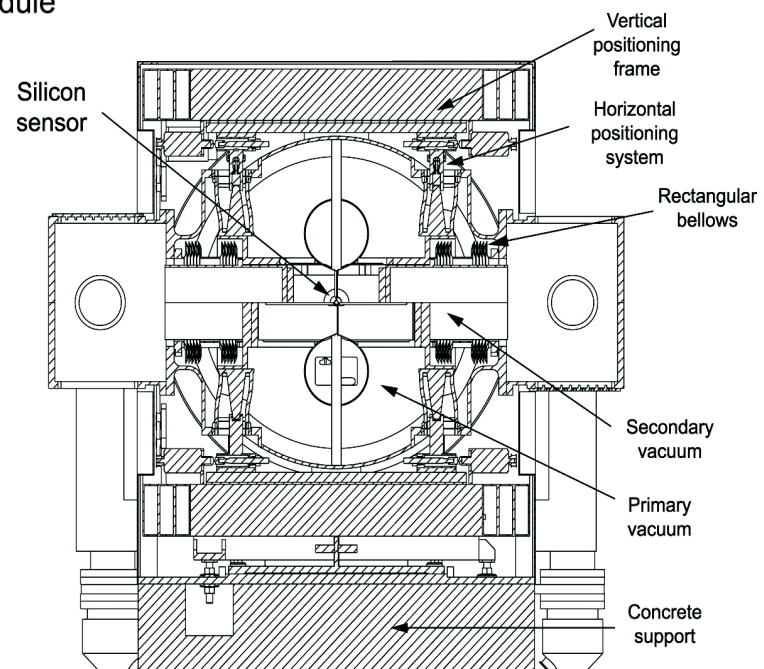
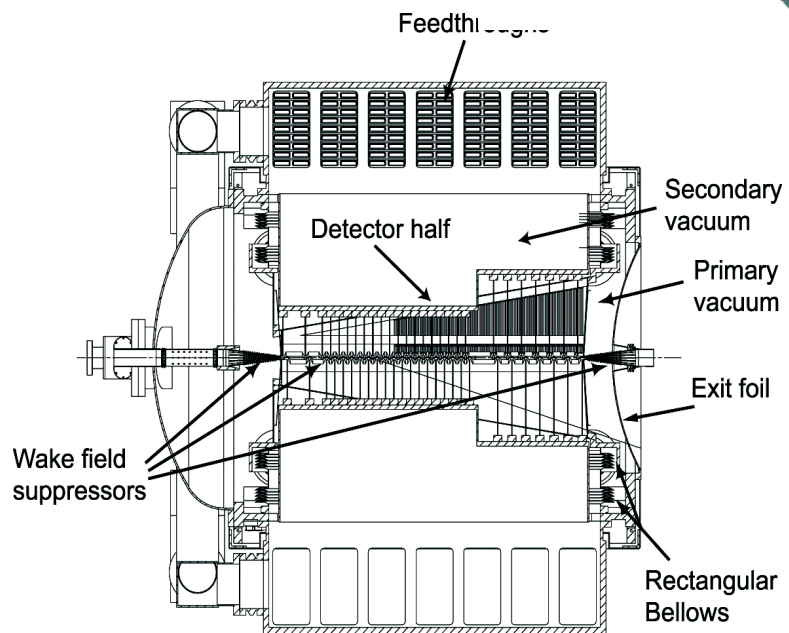
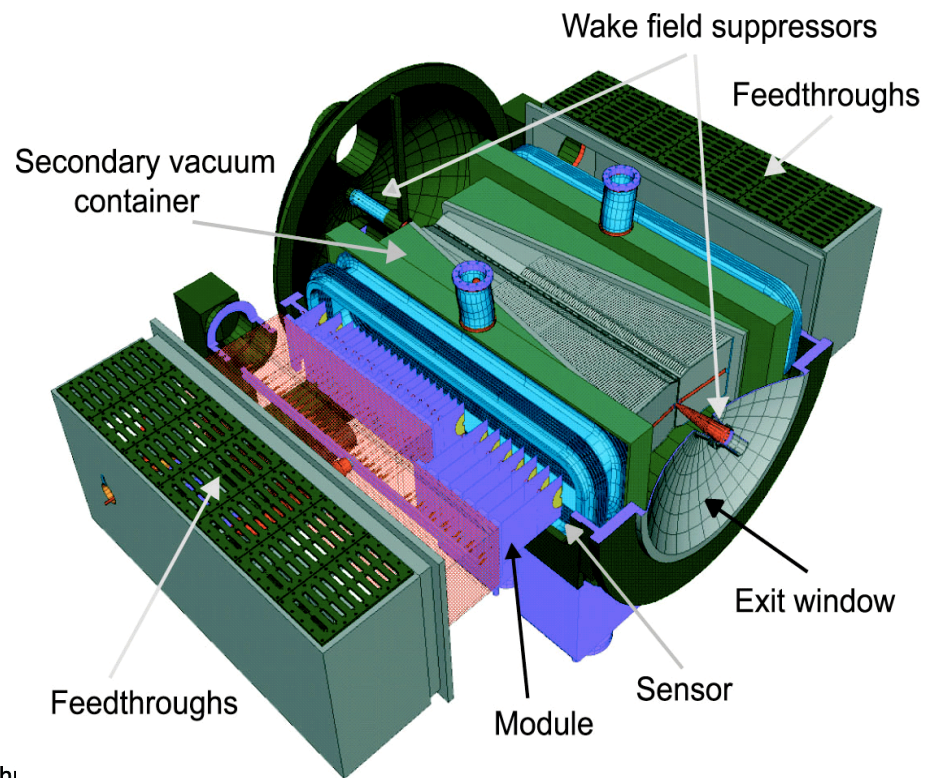
Muon system



VEetex LOcator(VELO)



Number of silicon sensors:	100
Area of silicon:	0.32m ²
Number of channels:	204,800

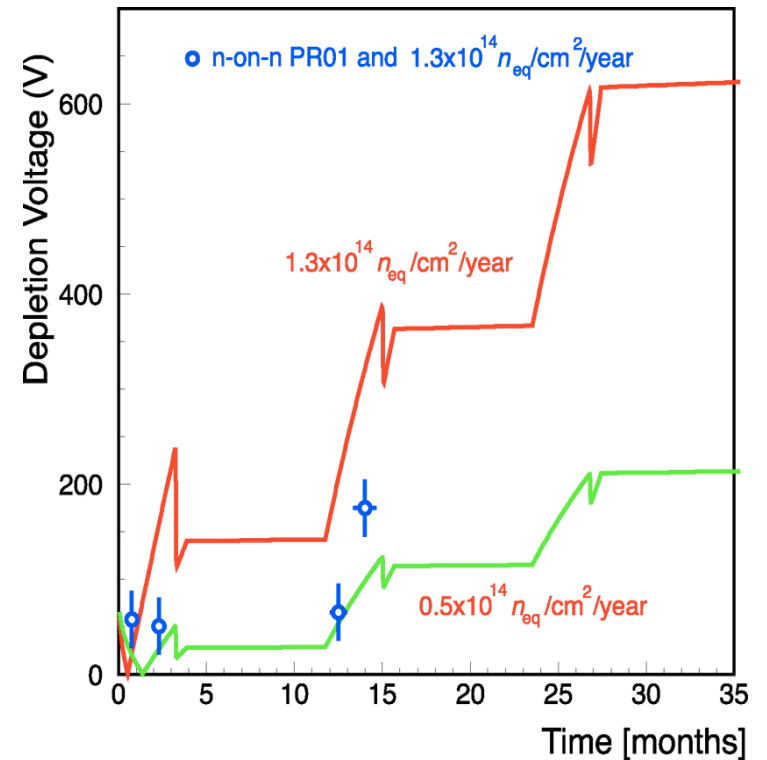
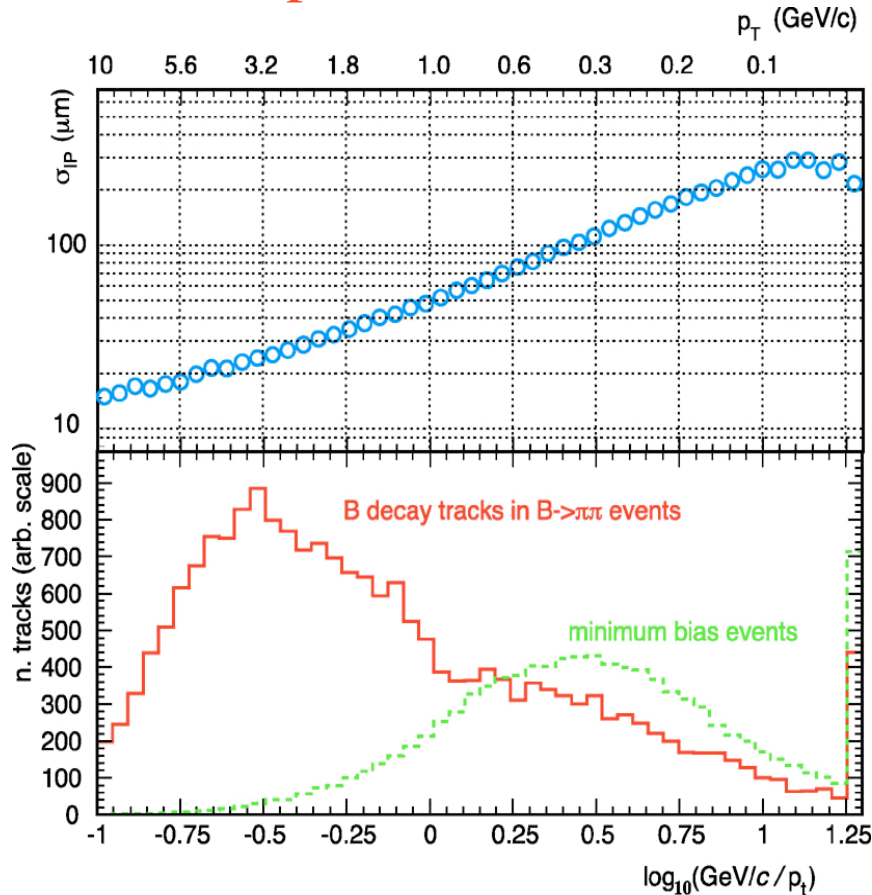
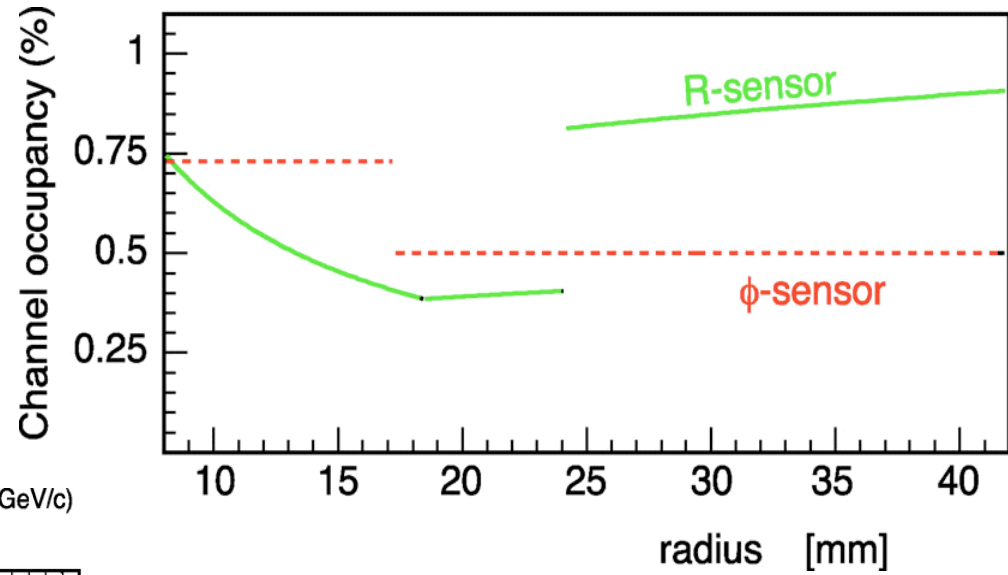


The VELO has

- low occupancy
- high resolution
- radiation hard

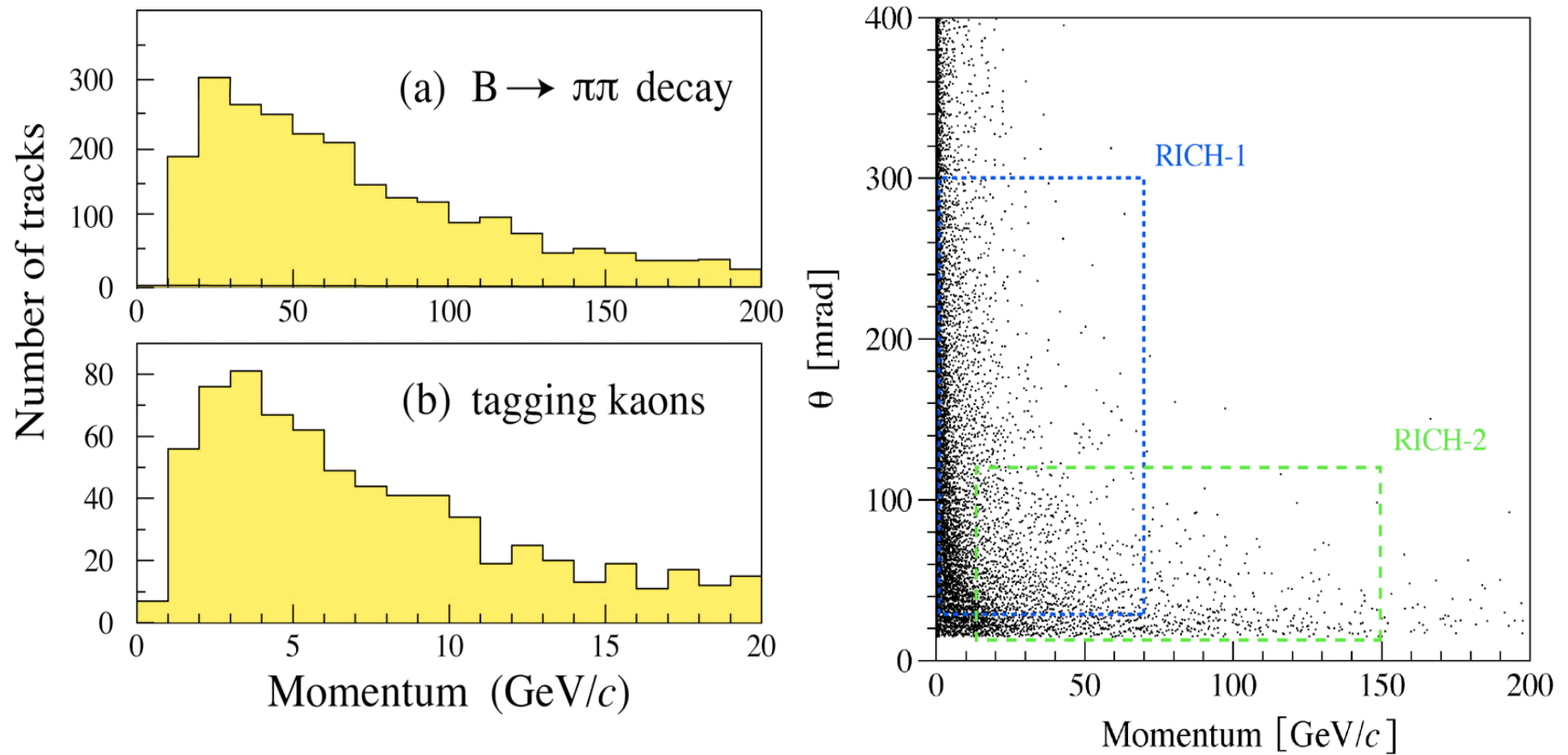
with only $\sim 200k$ channels:

□ optimal choice!

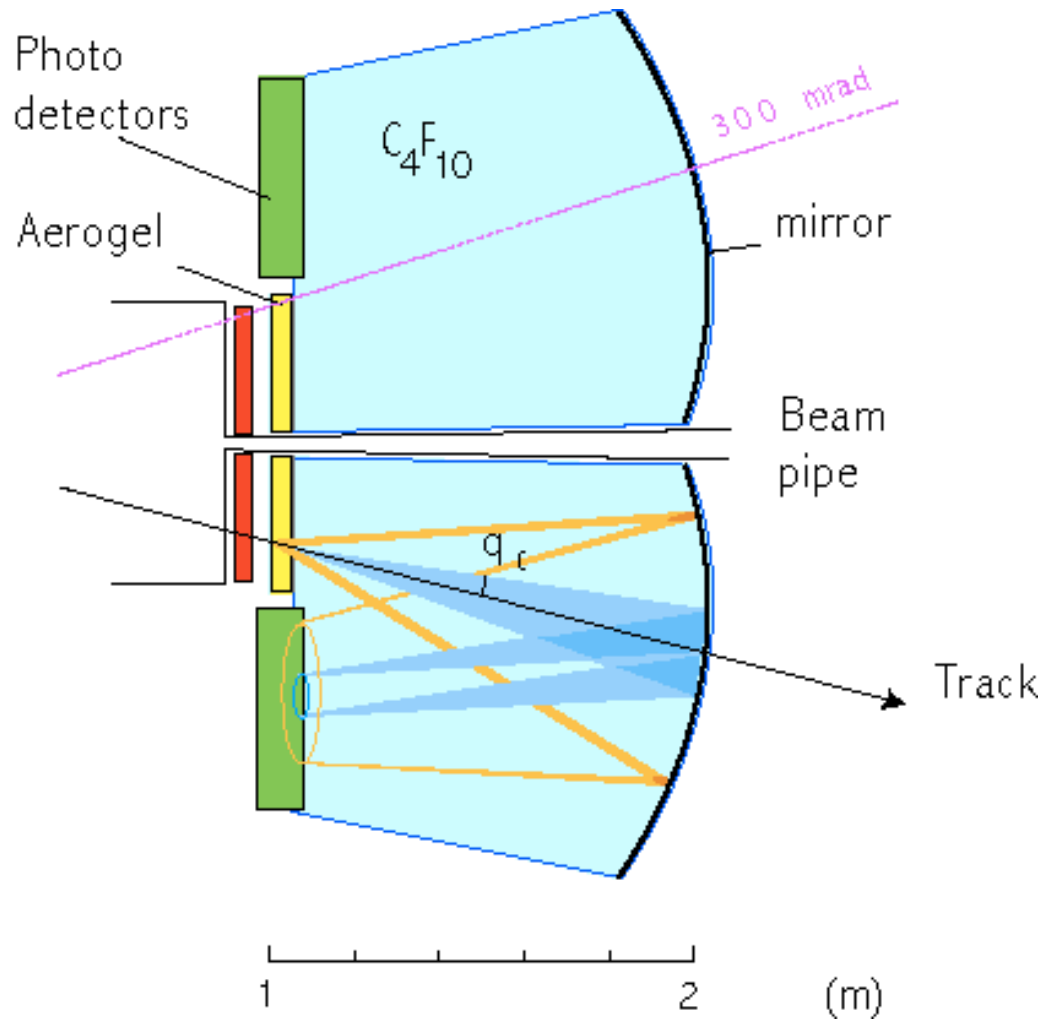


RICH

Required momentum range and angular coverage.

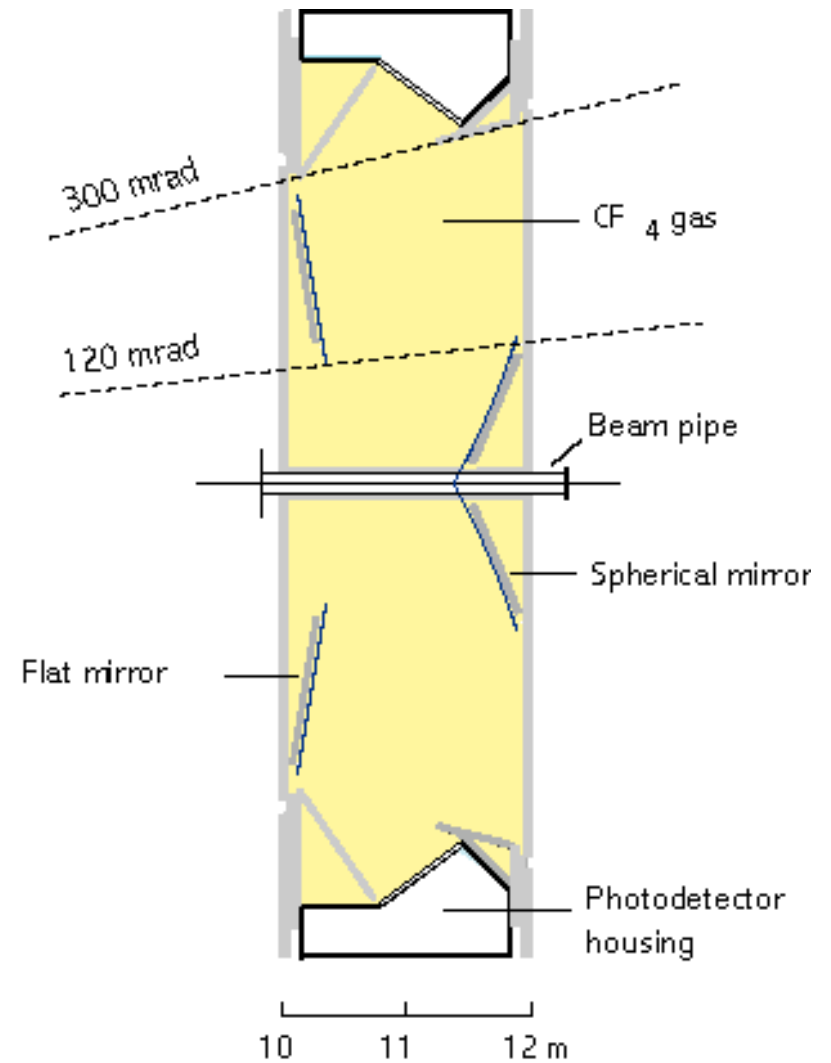


RICH1:



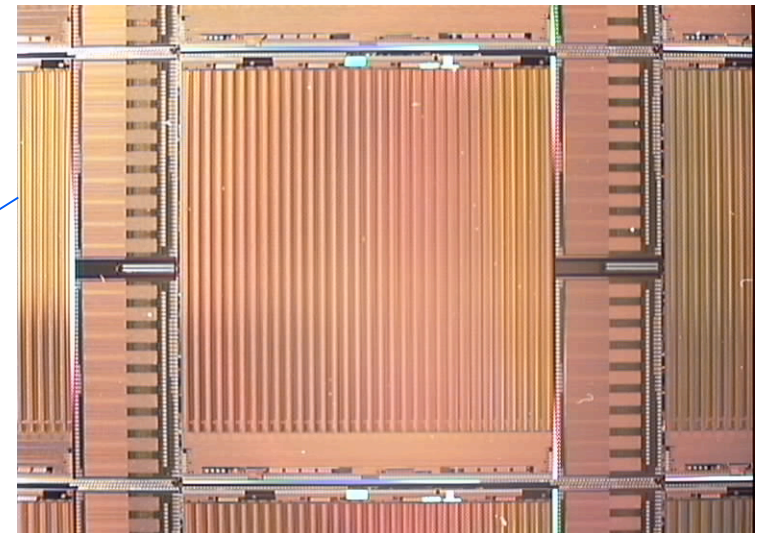
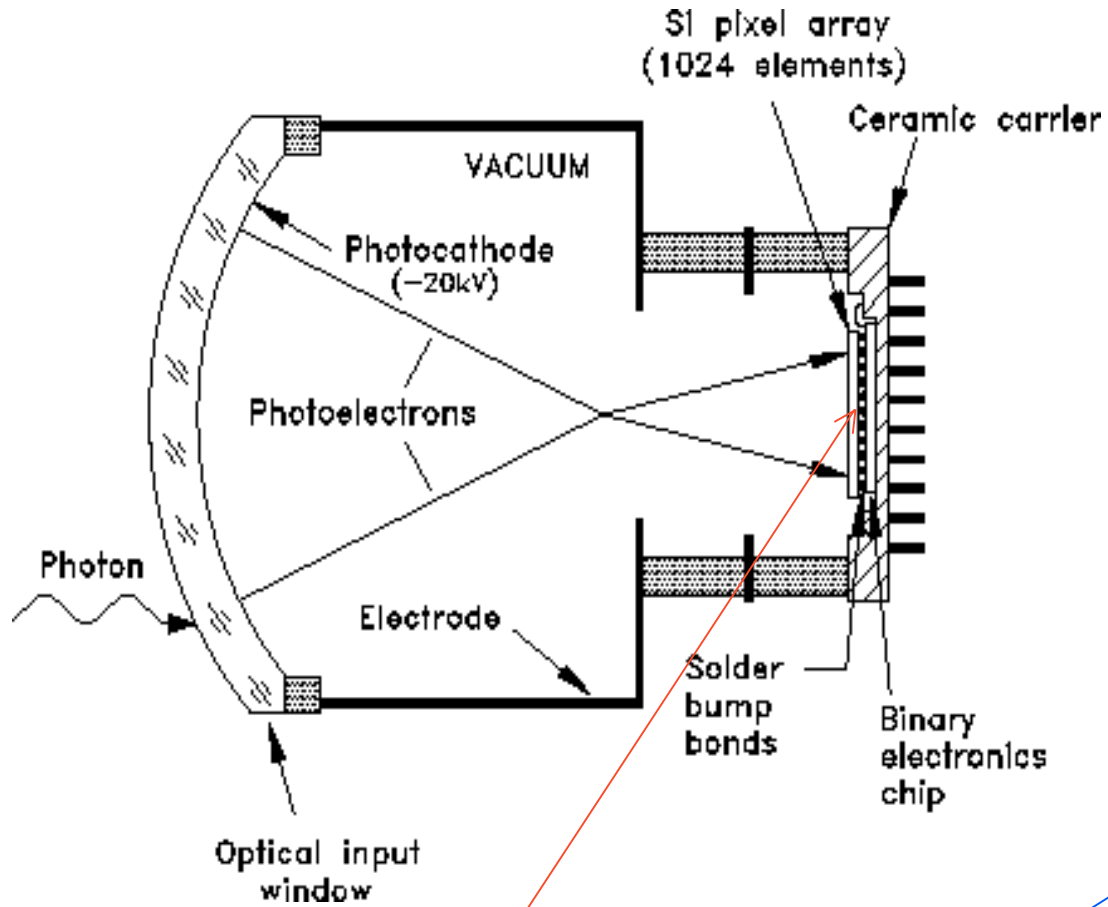
5cm aerogel, $n = 1.03$, 2-11 GeV
4 m³ C_4F_{10} , $n = 1.0014$, 10-70 GeV

RICH2:



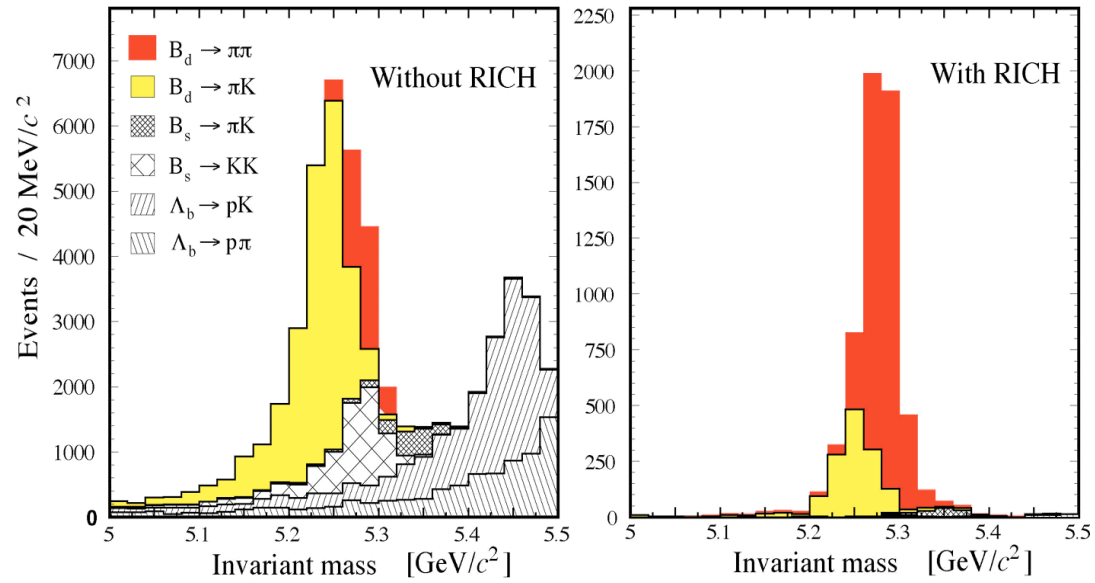
100 m³ CF_4 , $n = 1.0005$
17-150 GeV

Baseline option Pixel HPD

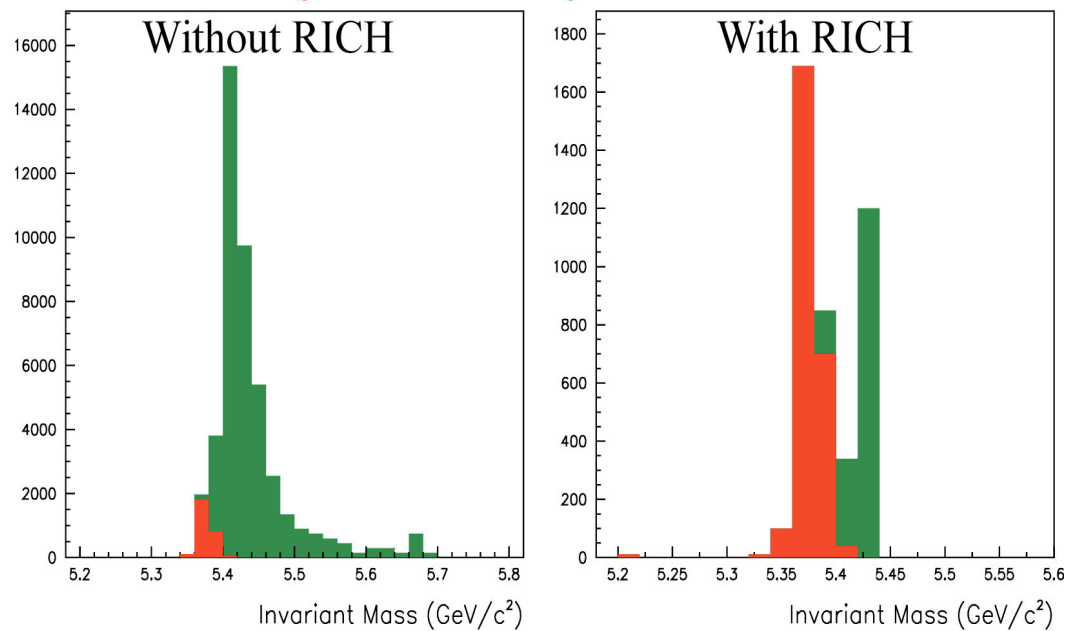


1024 pixel ($500\mu\text{m} \times 500\mu\text{m}$) detector and bump bonded **pixel readout electronics**.

$B \rightarrow \pi \pi + \pi \pi$



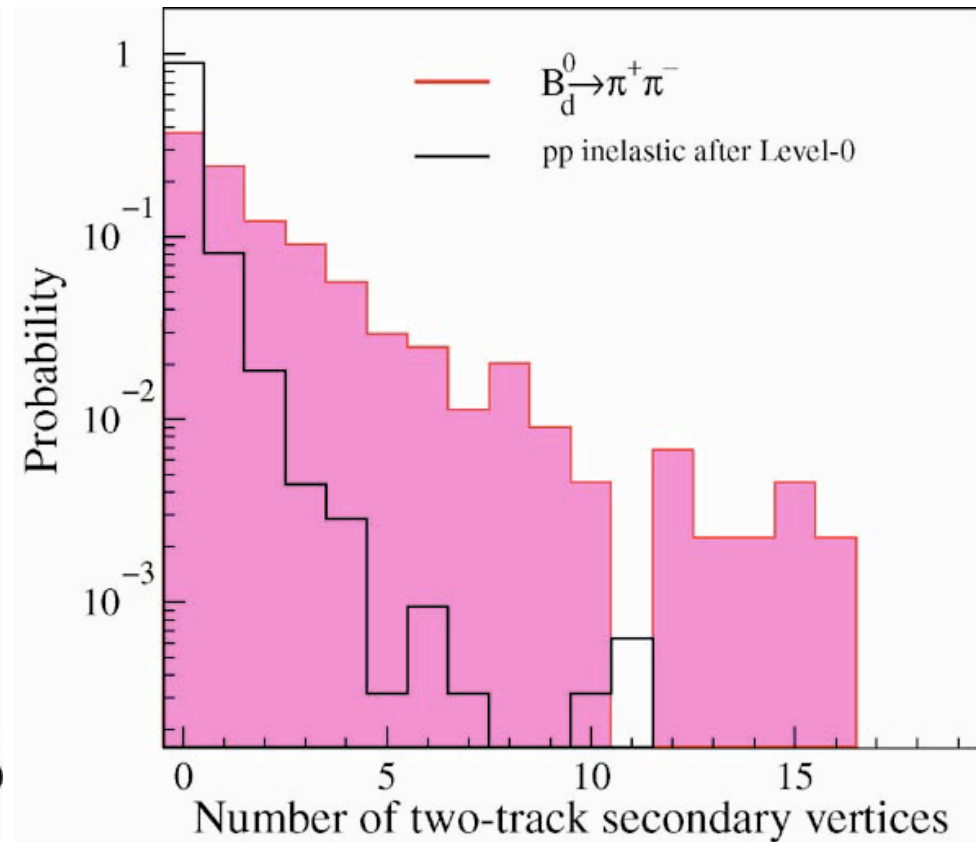
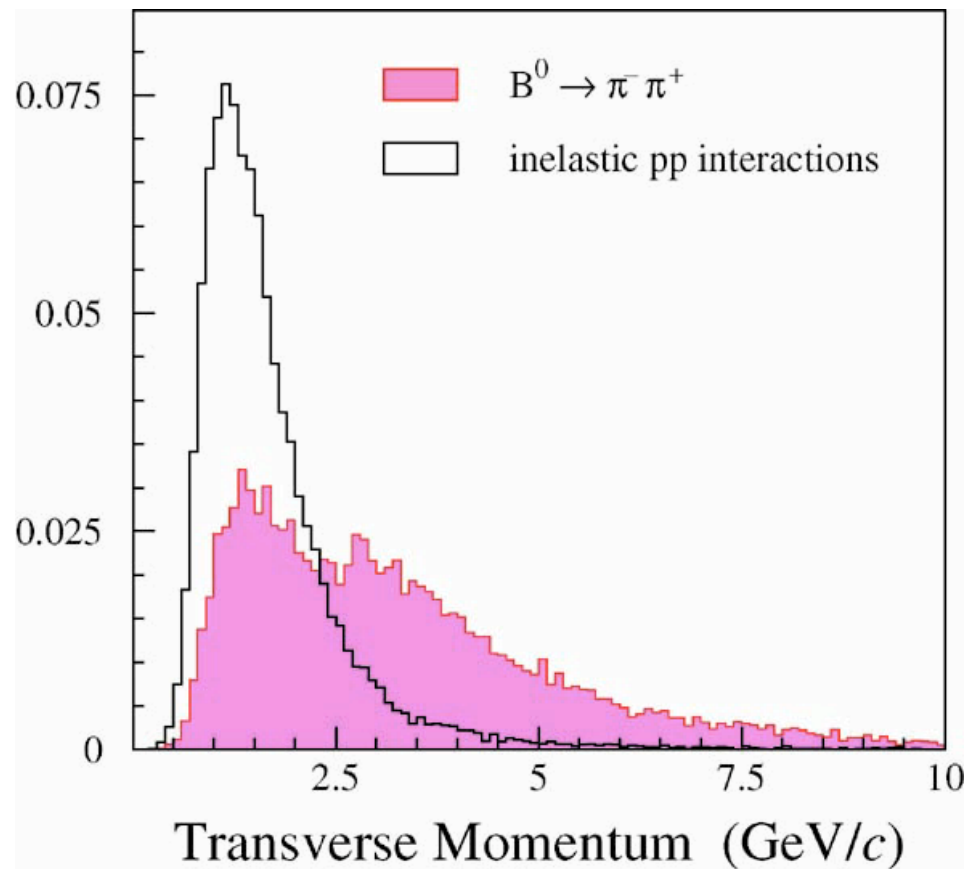
$D_s K$ signal $D_s \pi$ background



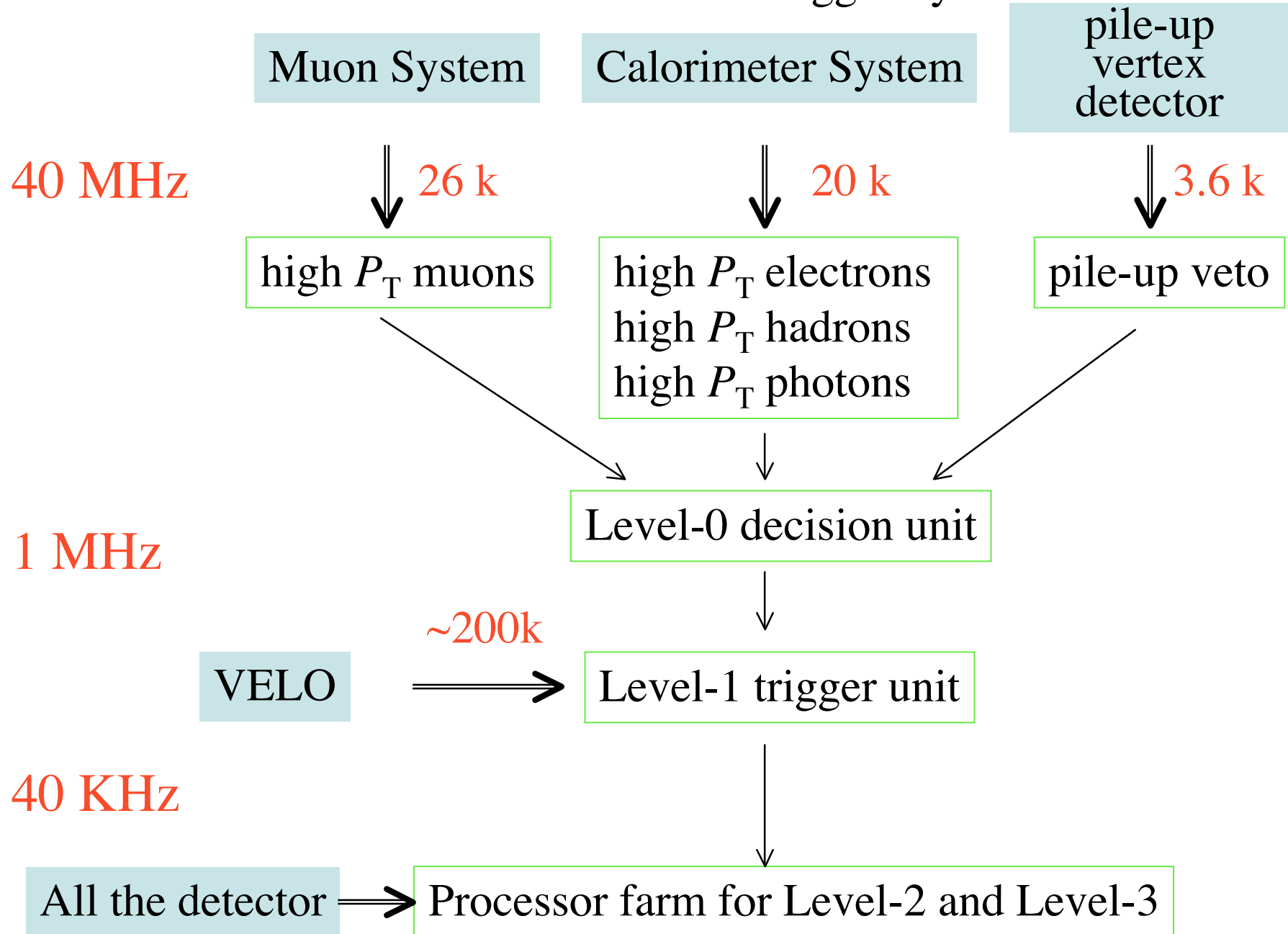
Trigger

Level-0: High p_T leptons and hadrons

Level-1: Detached decay vertices



The LHCb Trigger System



LHCb Trigger Efficiency

for reconstructed and correctly tagged events

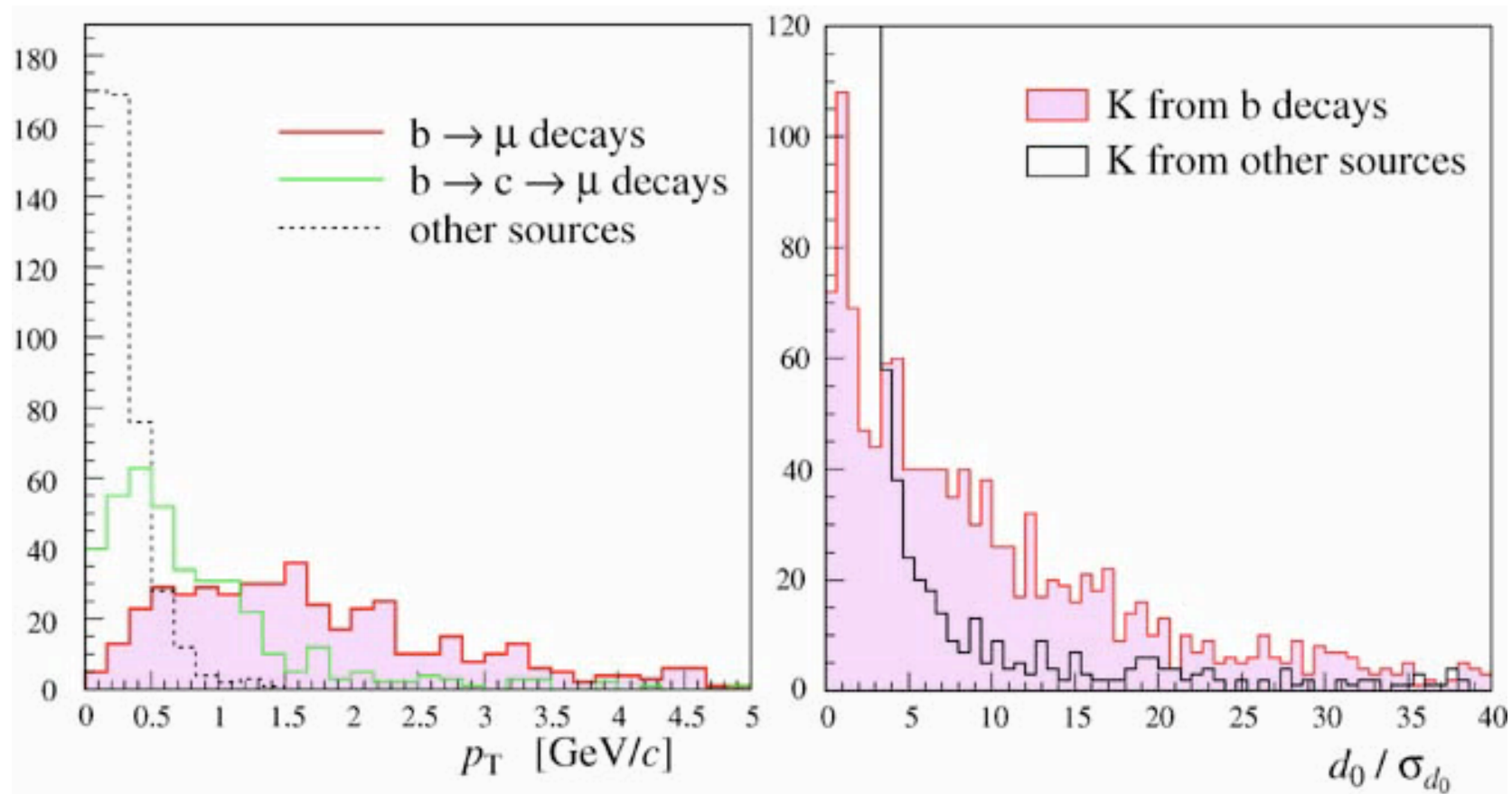
	L0(%)				L1(%)	L2(%)	Total(%)
	μ	e	h	all			
$B_d \rightarrow J/\psi(ee)K_S + \text{tag}$	17	63	17	72	42	81	24
$B_d \rightarrow J/\psi(\mu\mu)K_S + \text{tag}$	87	6	16	88	50	81	36
$B_s \rightarrow D_s K + \text{tag}$	15	9	45	54	56	92	28
$B_d \rightarrow DK^{\pm}$	8	3	31	37	59	95	21
$B_d \rightarrow \mu^+\mu^- + \text{tag}$	14	8	70	76	48	83	30

- trigger efficiencies are $\sim 30\%$
- hadron trigger is important for hadronic final states
useful only with the kaon tag
- lepton trigger is important for final states with leptons

evenly spread selectivity = robust and flexible

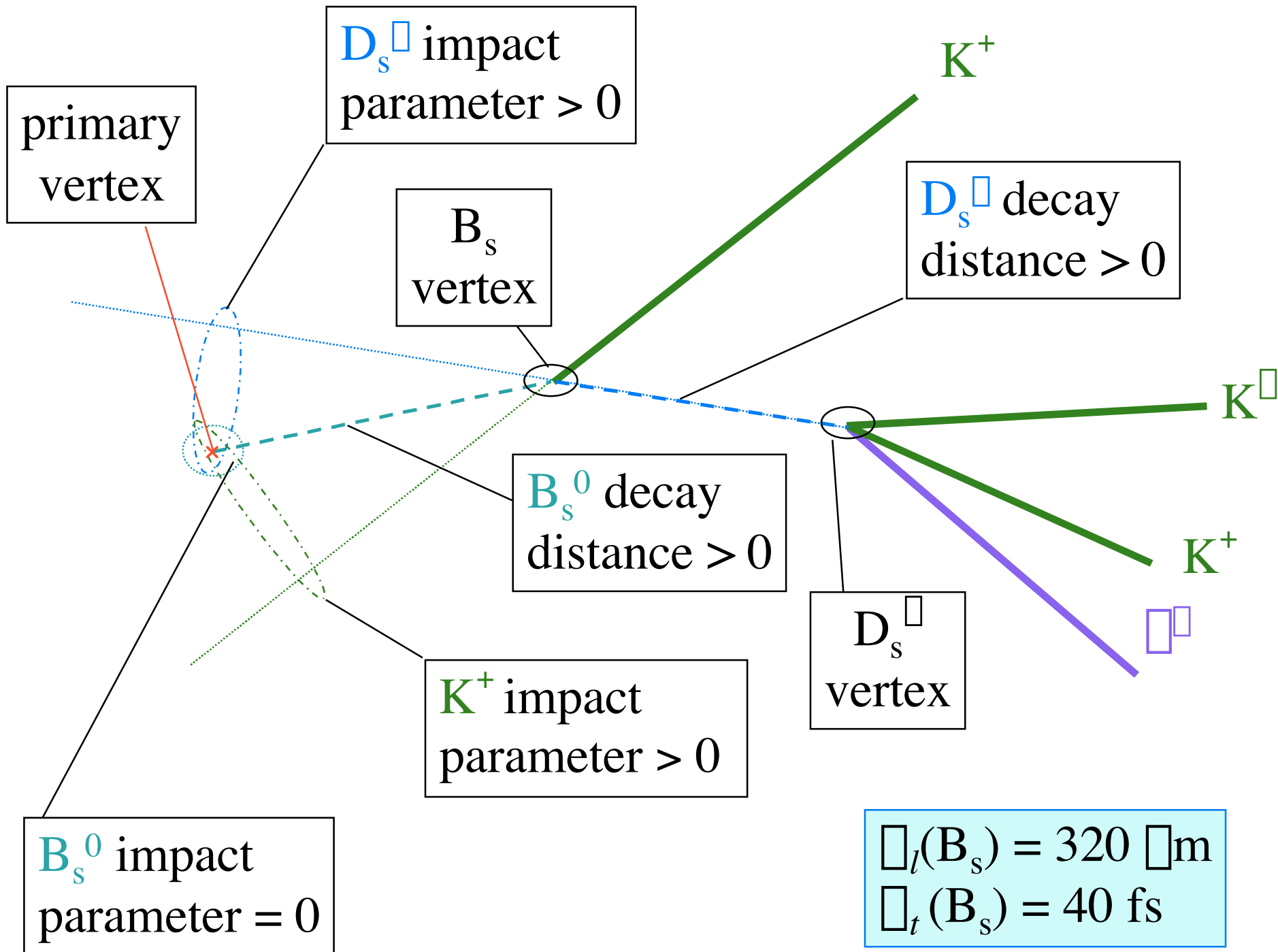
8) Something about reconstruction

Flavour tag with high p_T leptons and large impact parameter (d_0) kaons



$B-\bar{B}$ oscillations: irreducible wrong tag

$$\epsilon_{\text{tag}} = 0.4, \epsilon_{\text{tag}} = 0.2$$



9) Some LHCb Physics Performance

CP asymmetries in

$B_d \rightarrow J/\psi K_S$ (>40k tagged / 10^7 sec)

$\Delta_{\sin 2\beta} = 0.02 / 10^7$ sec

$B_s \rightarrow J/\psi \eta$ excellent Δ_t

$\Delta_{2\beta} = 0.04 - 0.06 / 10^7$ sec ($x_s = 20 - 40$)

B_s oscillations: hadron trigger, excellent Δ_t

$B_s \rightarrow D_s \eta$ measurable up to $x_s \approx 80$ (54 ps⁻¹) with 5%

CP violation in radiative decays

$B_d \rightarrow K^{*0} \eta$ s = 4.5k / 10^7 sec (Br=1.5% 10^6), s/b = 16

error on forward-backward asymmetry $\approx 0.03 / 10^7$ sec

$B_d \rightarrow K^{*0} \eta$ single photon trigger

s = 26k / 10^7 sec (Br=4.9% 10^5), s/b = 1

error on CP asymmetry $\approx 0.01 / 10^7$ sec

CP asymmetries in

$$B_d \rightarrow D^{\pm} \pi^{\mp}$$

hadron trigger

$$260 \text{ k tagged} / 10^7 \text{ sec} \quad \epsilon \approx 10^\circ$$

$$B_s \rightarrow D_s^{\mp} K^{\pm}$$

particle ID, hadron trigger, excellent ϵ_t

$$2.4 \text{ k tagged} / 10^7 \text{ sec} \quad \epsilon \approx 10^\circ$$

Time dependent Dalitz plot study: hadron trigger

$$B_s \rightarrow \pi^+ \pi^0 \pi^0$$

$$\epsilon_{\pi^+ \pi^0} \approx 5^\circ \text{ to } 10^\circ / 10^7 \text{ sec}$$

CP asymmetries in

$$B_d \rightarrow \pi^+ \pi^0$$

particle ID, hadron trigger

$$4.9 \text{ k tagged} / 10^7 \text{ sec} \quad (\text{Br} = 5 \times 10^{-6})$$

$$\text{and } B_s \rightarrow K^+ K^0$$

particle ID, hadron trigger, excellent ϵ_t

$$4.6 \text{ k tagged} / 10^7 \text{ sec} \quad (\text{Br} = 1.9 \times 10^{-5})$$

$$\epsilon \approx 5.4^\circ / 10^7 \text{ sec for } \Gamma_{m_s} = 20 \text{ ps}^{-1}$$

Rare decays

$$B_s \rightarrow \pi^+ \pi^0$$

$$s = 10 / 10^7 \text{ sec} \quad (\text{Br} = 3.5 \times 10^{-9}), \quad s/b = 3$$

10) Conclusions

Unique property of LHC

Large yield for different b hadrons ($10^{12}/10^7$ sec)



Unique properties of LHCb

Trigger: high P_T lepton or **hadron + secondary vertex**

Sensitive to final states: lepton only & leptons+hadrons & **hadron only**

RICH system: clean K/ π separation over wide range of p

Kaon tag (enhancing the importance of the hadron trigger),

Clean reconstruction of **hadronic final states**

Vertex detector: excellent decay time resolution

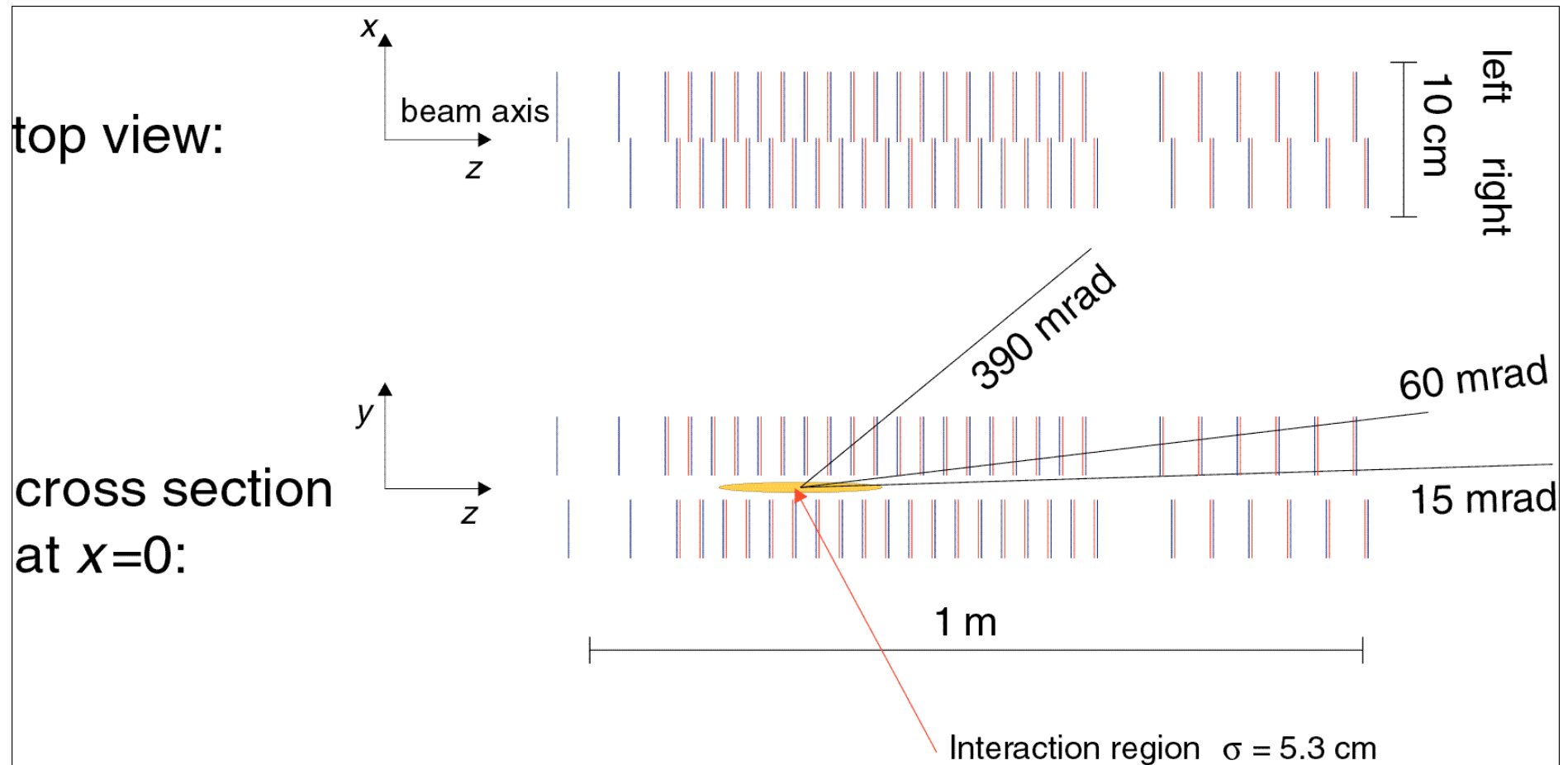
High CP and oscillation sensitivities for the **B_s system**



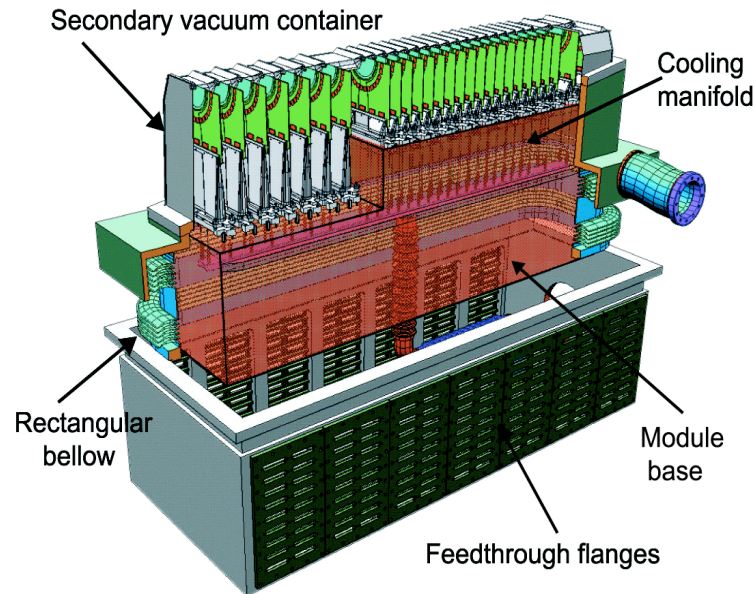
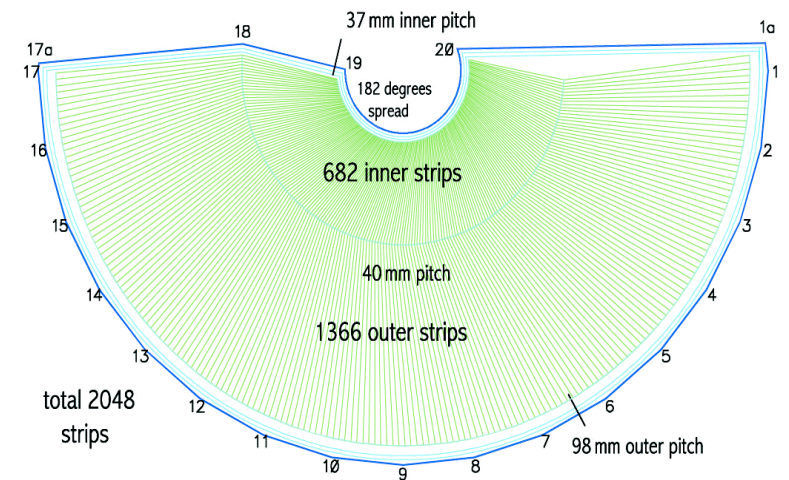
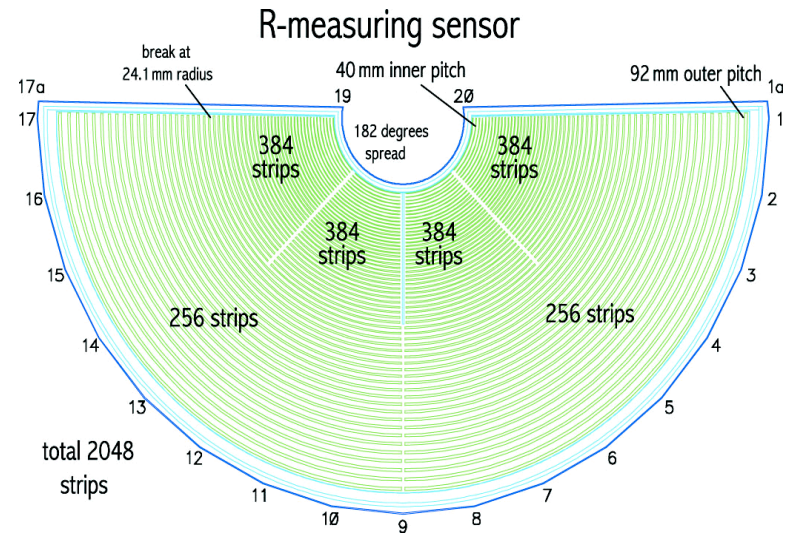
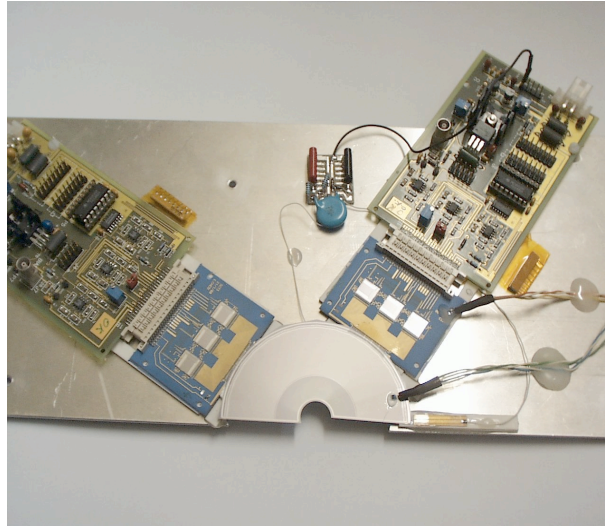
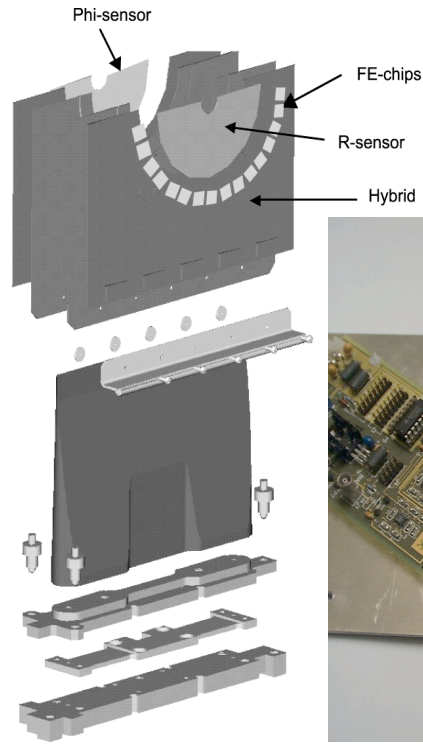
Unique opportunity to search for new physics

through CP violation and rare decays in B-meson decays!!

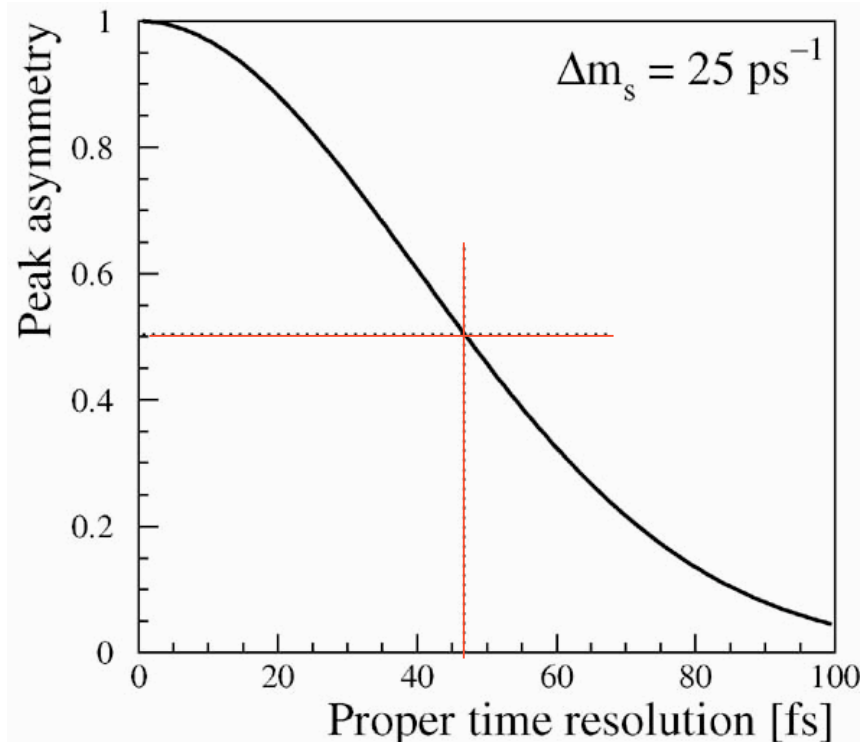
stations are overlapped



Small r - ϕ strip Si detector: 300 μ m n-on-n double metal layer



Decay time resolution...



Dilution of CP asymmetry for B_s due to Δt must be less than due to other sources: < 0.5

$$\Delta t < 50 \text{ fs} \approx 0.03 \lambda_B$$

$$\langle l_B \rangle \approx 1 \text{ cm}$$

$$\Delta l \approx 300 \text{ } \mu\text{m}$$

RICH Performance

Realistic simulation:

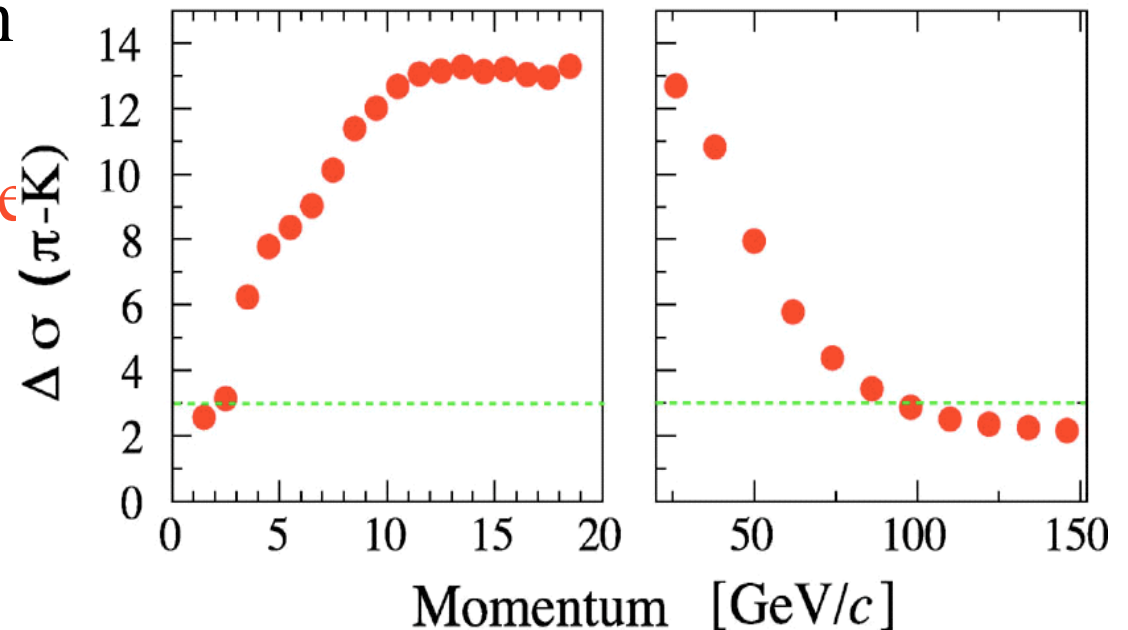
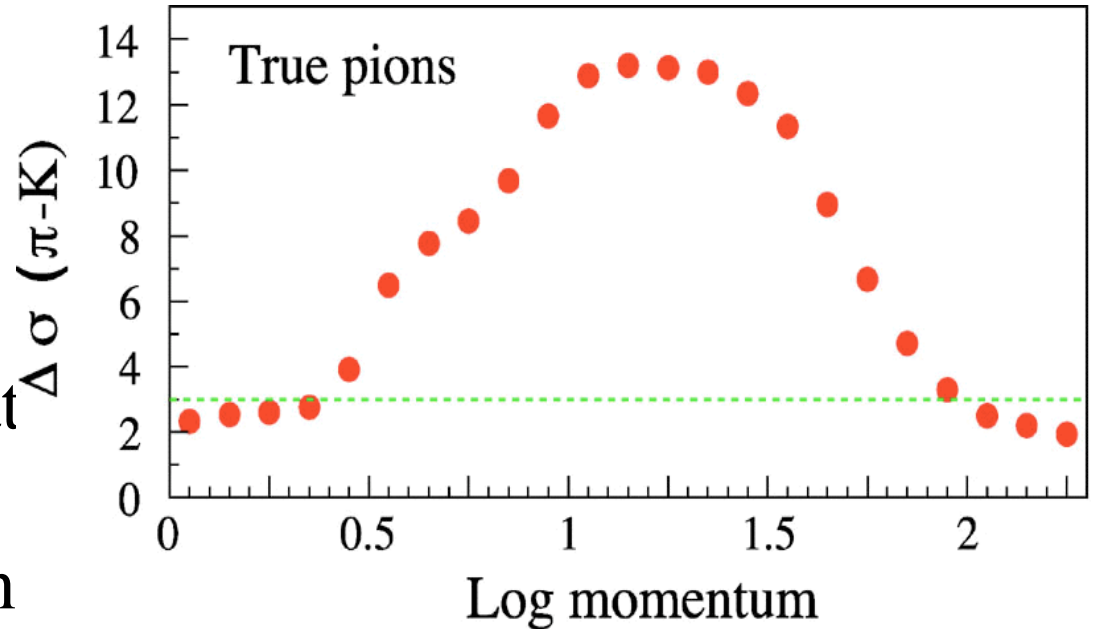
- tested by the test beam data
 - engineering design
 - measured HPD performance
 - all the background photons
 - pattern recognition
- (some can still improve)

No. of detected photons

6.6: RICH-1 aerogel

32.7: RICH-1 C_4F_{10}

18.4: RICH-2 CF_4

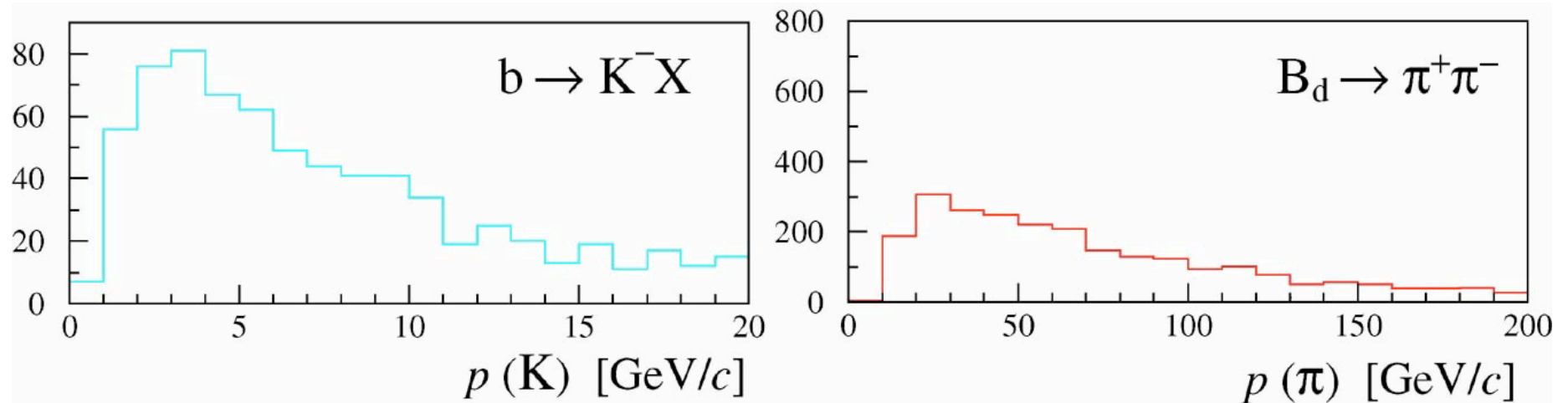


Particle identification is required in a momentum range of

$$p_{\min} = \sim 1 \text{ GeV}/c \quad (\text{Kaon tag})$$

to

$$p_{\max} = \sim 100 \text{ GeV}/c. \quad (\text{two-body B decay products})$$



Ring Imaging Cherenkov is a suitable technique.

Can new physics make CP violation in oscillation to be large?

$\epsilon_{bd(s)}^{NP}$ can increase $\arg M_{12}/\bar{M}_{12}$

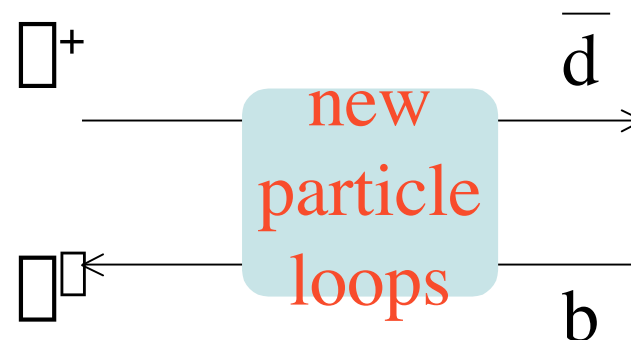
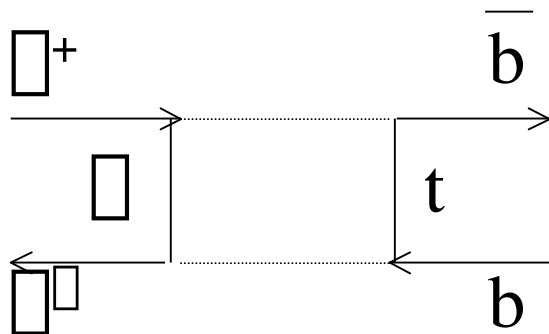
$$d \equiv \text{Im} \frac{\bar{M}_{12}}{M_{12}} = \epsilon \frac{\epsilon}{m} \sin(\arg M_{12}/\bar{M}_{12})$$

This may increase

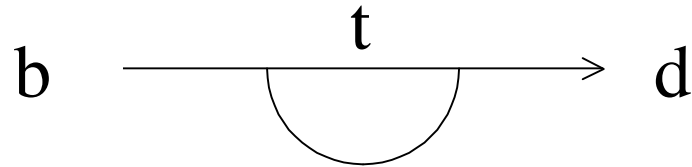
This can hardly change from $\sim 10^{-3}$.

$d > 10^{-3}$ is not really possible.

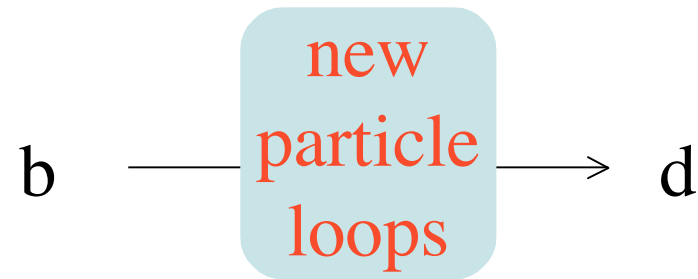
But could enhance $B_s \rightarrow \bar{b} + b$



Even better would be?



$$V_{bt} V_{td}^* \mu e^{i\phi}$$



$$A_{bd}^{NP} \mu e^{i\phi_{pbd}^{NP}}$$

$$|V_{bt} V_{td}^*| = \frac{1}{\sqrt{(1-\lambda)^2 + \lambda^2}} |V_{bt} V_{ts}^*|$$

more suppressed



New physics can have a larger influence.

$$B_d \approx \lambda^0, B_s \approx \lambda K_S$$