

Mesoscopic Fluctuations of the Pairing Gap

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**Workshop on
Level Density and
Gamma Strength in Continuum**

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Mesoscopic Fluctuations of the Pairing Gap¹⁾

I. Shell structure

II. Periodic orbit description of mean field

Fluctuations of shell energy – nuclear masses

Role of regular and chaotic dynamics

III. Periodic orbit description of pairing

Generic description of fluctuations of the pairing gap

IV. Fluctuations of nuclear pairing gap

V. Applications to other finite fermi systems

Nanosized metallic grains

Ultracold fermionic gases

¹⁾H. Olofsson, S. Åberg and P. Leboeuf, submitted



Mesoscopic Fluctuations of the Pairing Gap

”Fluctuations”:

**Variation of quantum properties at zero temperature,
(such as shell energy or pairing gap)
when shape, particle number, etc is varied.**



Many-body system

Hamiltonian:
$$H = \sum_i H_i + \frac{1}{2} \sum_{i,j} W(i, j)$$

$$H_i = T_i + V_i$$

V_i : external confinement potential or a mean field.

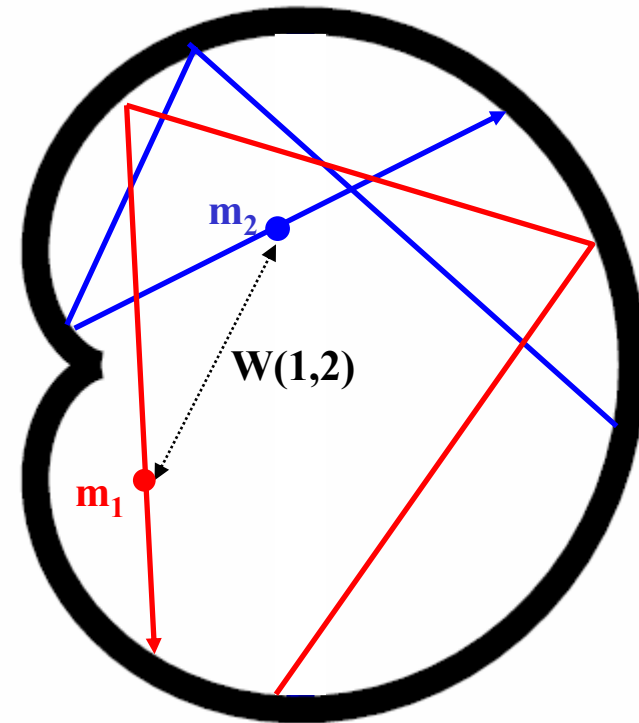
One-body dynamics in V_i can be regular, mixed or chaotic.

$W(i,j)$: residual two-body interaction.

$W(i,j)$ important for highly excited states.

Main interest here: ground states

⇒ Include only mean field

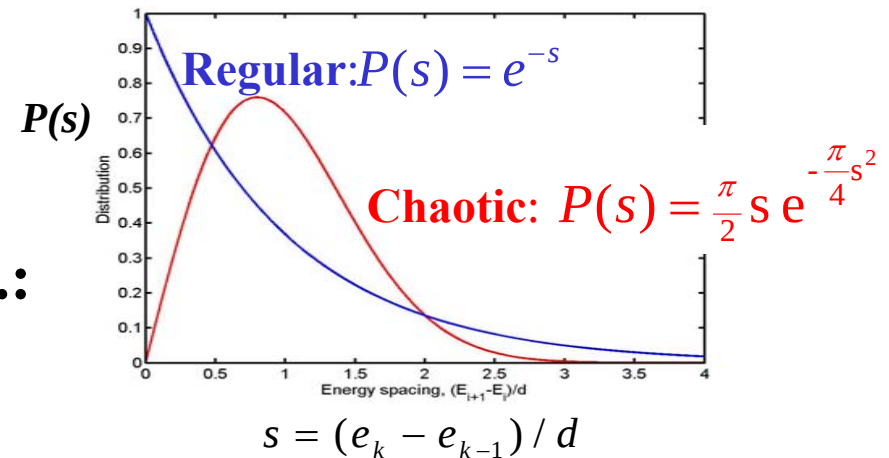


I. Shell structure - Mean field

One-body Hamiltonian: $\hat{H} = \hat{T} + \hat{V} = \sum_k e_k a_k^+ a_k$

Mean field $\hat{V}$ implies (classically) regular/chaotic one-body motion.
Can be distinguished in quantum mech.:

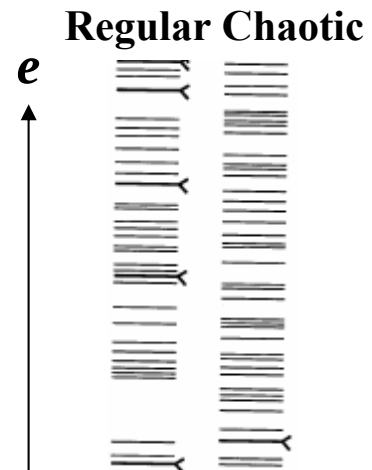
Distribution of nearest neighbor energy spacings:



One-body level density: $\rho(e) = \sum_{\lambda_{Fermi}^k} \delta(e - e_k) = \bar{\rho} + \tilde{\rho}$

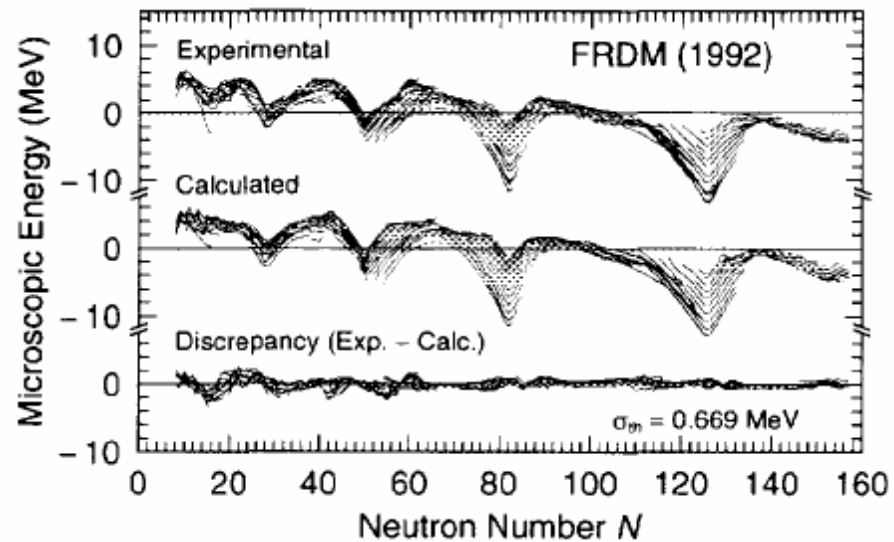
Ground-state energy: $E(N) = \int_{-\infty}^{\lambda_{Fermi}} e \rho(e) de = \bar{E} + \tilde{E}$

Shell structure seen in $\tilde{\rho}(e)$ or in $\tilde{E}(N)$
(suppressed in chaotic systems)



I. Shell structure - Nuclear masses

Shell energy $\tilde{E}$ versus neutron number



P. Möller et al, Atomic data and nucl data tables **59** (1995) 185



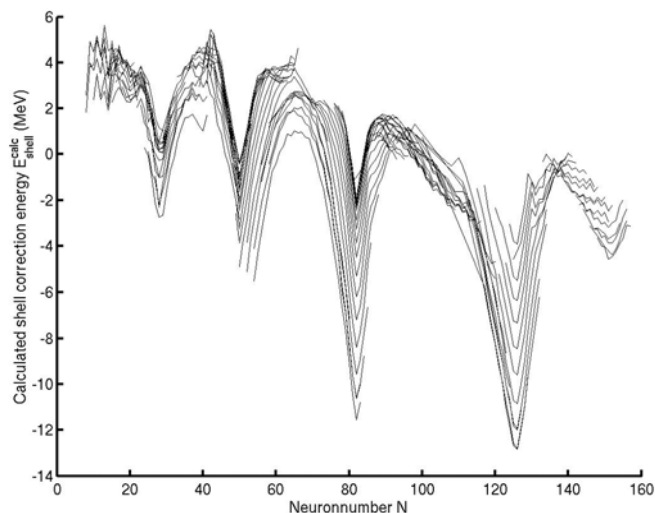
Nuclear masses

$$m(N, Z)/c^2 = E_{L.D.} + E_{shell} + E_{rest}$$

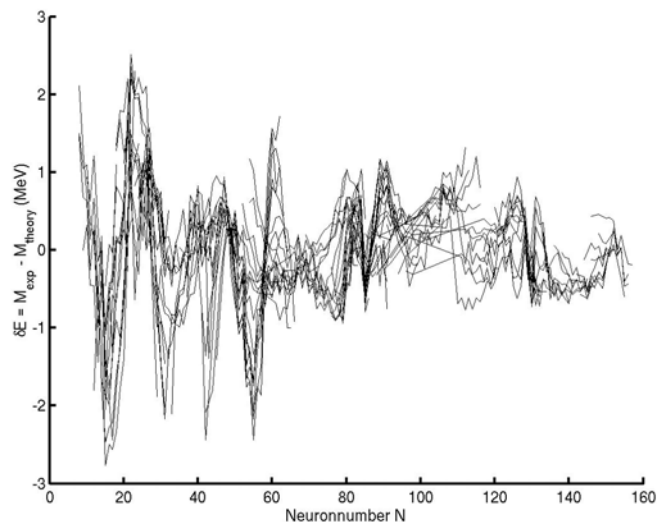
$$E_{L.D.} \approx 8A \text{ MeV}, \quad E_{shell}^{rms} \approx 3 \text{ MeV}, \quad E_{rest}^{rms} \approx 0.7 \text{ MeV}$$

Extensive mass calculation with Folded-Yukawa potential + finite range droplet
by Möller, Nix, Myers, Swiatecki, Atom. Data Nucl. Tables **59**, 185 (1995):

Calculated shell energy



Error in mass formula



II. Periodic orbit theory

The fluctuating part of the **level density**, $\rho(e) = \bar{\rho} + \tilde{\rho}$, is given by:

$$\tilde{\rho}(e) = \sum_{\substack{\text{periodic} \\ \text{orbits, } p}} \sum_{r=1}^{\infty} A_{p,r} \cdot \cos(rS_p / \hbar + \nu_{p,r})$$

$A_{p,r}$: stability amplitude

$S_p = \oint pdq$: action of periodic orbit p

$\nu_{p,r}$: Maslov index

$\tau_p = \partial S_p / \partial E$: period of p.o

The fluctuating part of the **total energy** for A particles then becomes [1]:

$$\tilde{E}(A) = \int_0^{e_F} \tilde{\rho}(e) \cdot e \cdot de = 2\hbar^2 \sum_p \sum_{r=1}^{\infty} \frac{A_{p,r}}{r^2 \tau_p^2} \cos(rS_p / \hbar + \nu_{p,r})$$

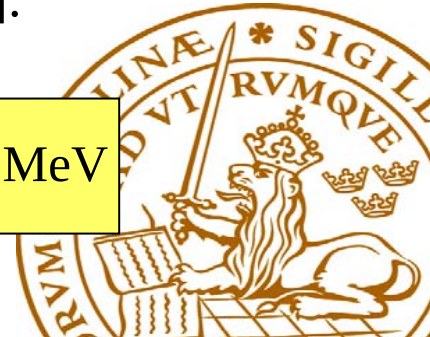
The second moment of $\tilde{E}$ can be evaluated and gives for nuclei [2]:

If regular:

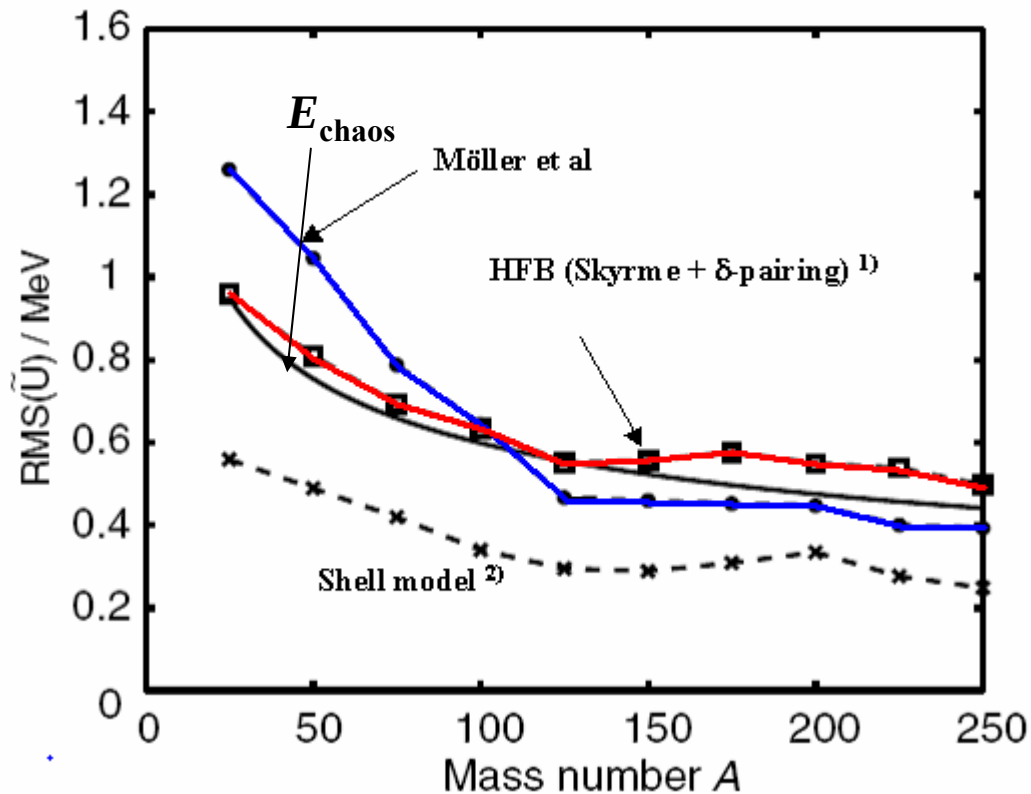
$$E_{\text{regular}}^{\text{rms}} = \sqrt{\langle E_{\text{regular}}^2 \rangle} \approx 2.8 \text{ MeV}$$

If chaotic:

$$E_{\text{chaos}}^{\text{rms}} = \sqrt{\langle E_{\text{chaos}}^2 \rangle} \approx \frac{2.78}{A^{1/3}} \text{ MeV}$$



Error in mass formulae compared to E_{chaos}



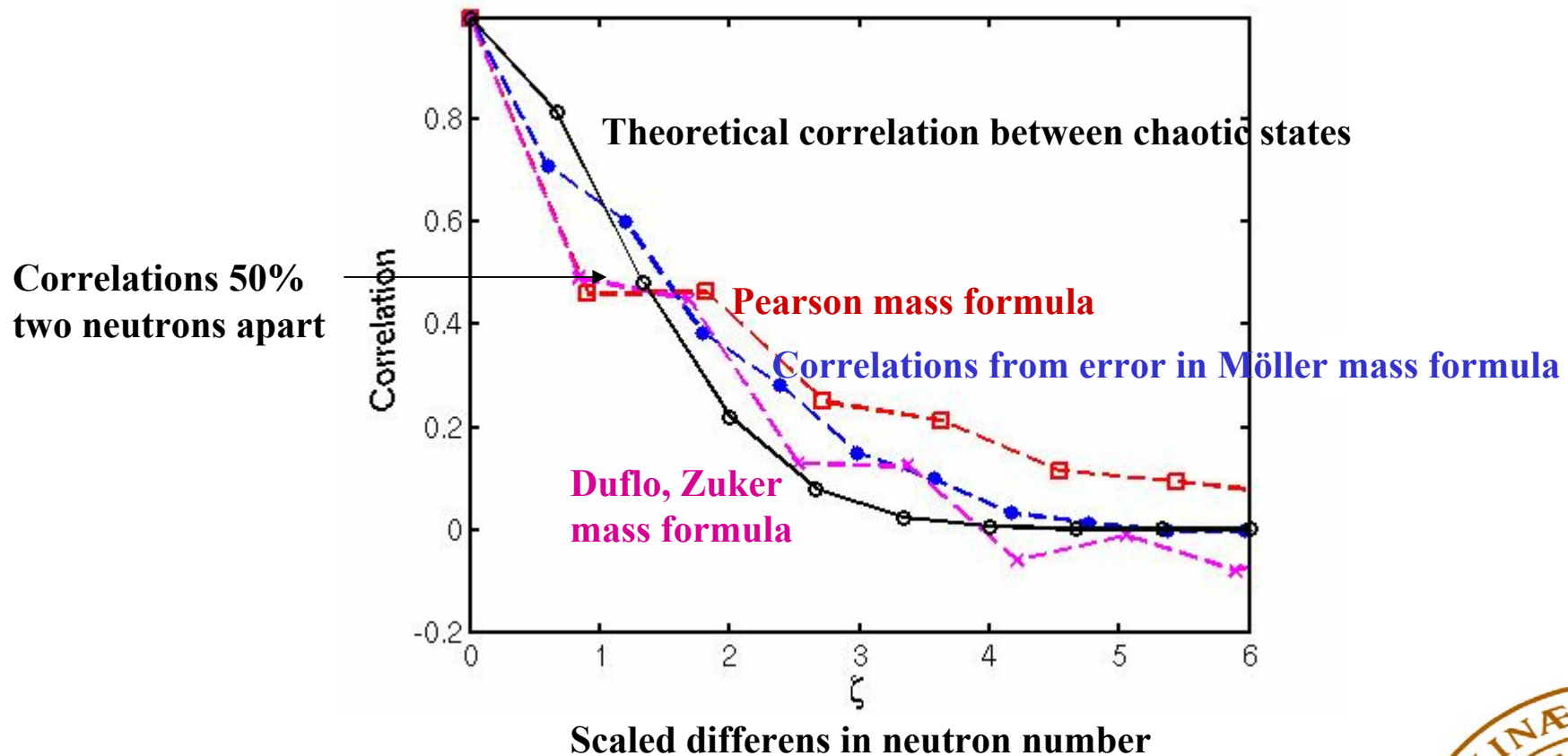
1) Samyn, Goriely, Bender, Pearson, PRC 70 (2004) 044309

2) Duflo, Zuker, PRC 52 (1995) R23



Autocorrelations between error in nuclear mass formulae [1]

$$C(x) = \langle E(A - x_N/2) E(A + x_N/2) \rangle$$



Notice:

- Good agreement between chaotic energy and error in nuclear mass!
- Chaotic energy is NOT random but strongly correlated!



III. Periodic orbit description of pairing

BCS theory

Hamiltonian:
$$H = \sum_k e_k a_k^+ a_k - G \sum_{kl} a_k^+ a_{\bar{k}}^+ a_{\bar{l}} a_l$$

Mean field approximation (in pairing space):

Pairing gap ("pairing deformation"):
$$\Delta = \left\langle G \sum_k a_k^+ a_{\bar{k}}^+ \right\rangle$$

is determined by
gap equation:

$$\frac{2}{G} = \sum_{\mu} \frac{1}{\sqrt{(e_{\mu} - \lambda)^2 + \Delta^2}}$$

$$\rightarrow \int_{-L}^L \frac{\rho(e) de}{\sqrt{e^2 + \Delta^2}}$$



III. Periodic orbit description of pairing

Divide pairing gap in smooth and fluctuating parts:

$$\Delta = \bar{\Delta} + \tilde{\Delta} \qquad \bar{\Delta} \approx 2L \exp\left(-\frac{1}{\bar{\rho}G}\right)$$

Insert into gap equation and assuming $\bar{\Delta} \ll L$

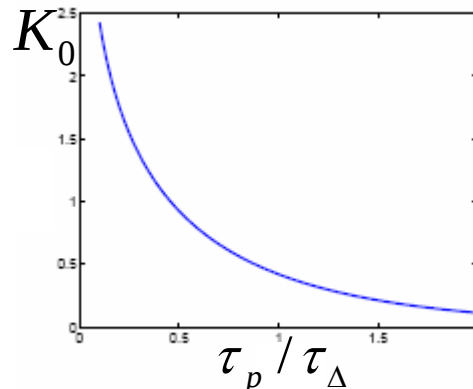
Expand to lowest order in fluctuations gives

$$\tilde{\Delta} = 2 \frac{\bar{\Delta}}{\bar{\rho}} \sum_{p,r} A_{p,r} K_0(r \tau_p / \tau_{\Delta}) \cos(r S_p(e) / \hbar + \nu_{p,r})$$

where

$\tau_{\Delta} = \frac{\hbar}{2\pi\bar{\Delta}}$ is characteristic time associated with pairing gap

$$K_0(x) = \int_0^{\infty} \frac{\cos(xt)}{\sqrt{1+t^2}} dt$$



i.e. no contribution from orbits with

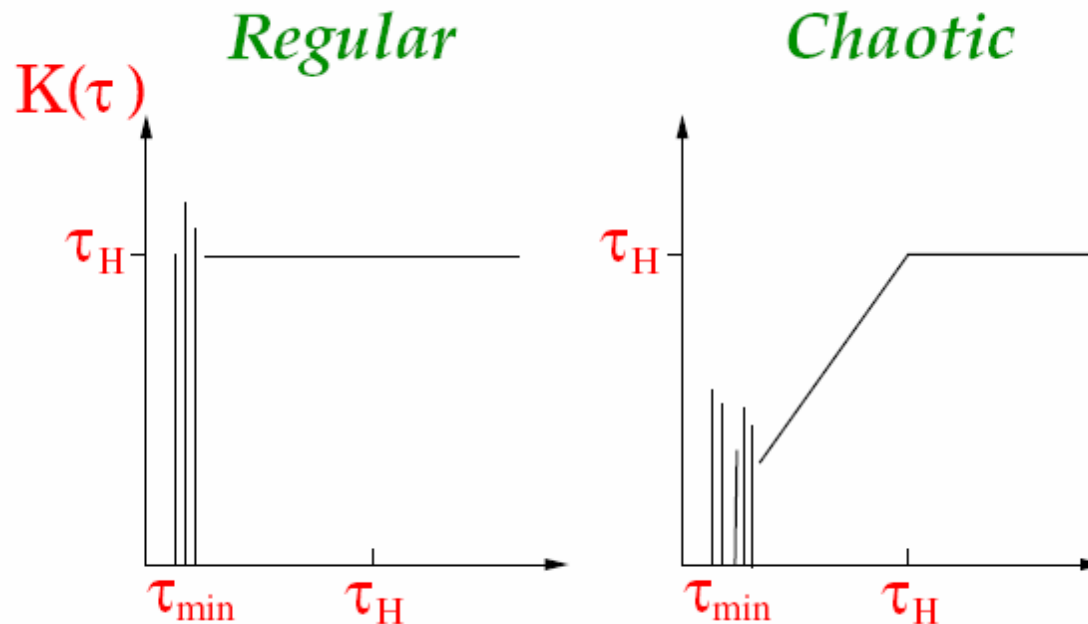
$$\tau_p \gg \tau_{\Delta}$$

III. Periodic orbit description of pairing

Fluctuations of pairing gap become

$$\langle \tilde{\Delta}^2 \rangle = 2 \frac{\overline{\Delta}^2}{\tau_H^2} \int_0^\infty d\tau K_o^2(\tau / \tau_\Delta) K(\tau)$$

where K is the spectral form factor (Fourier transform of 2-point corr. function):



$\tau_{\min}$ is shortest periodic orbit, $\tau_H = h\bar{\rho} = h / \delta$ is Heisenberg time



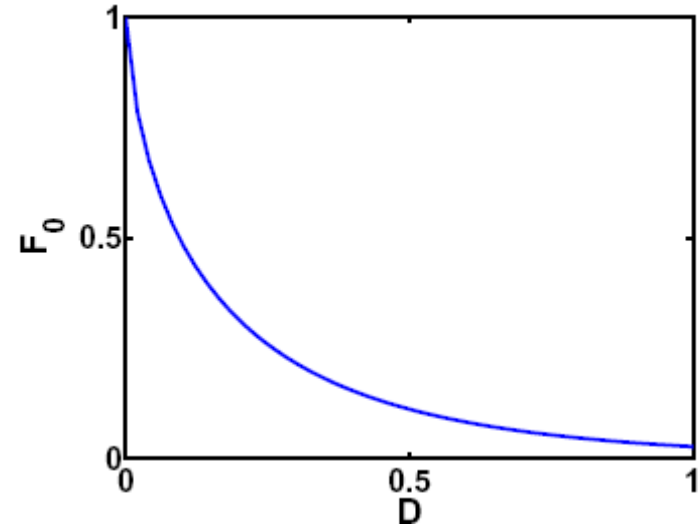
III. Fluctuations of pairing

Fluctuations of pairing, expressed in single-particle mean level spacing, δ :

$$\sigma = \sqrt{\langle \tilde{\Delta}^2 \rangle} / \delta$$

If regular:
$$\sigma_{reg}^2 = \frac{\pi}{4} \frac{\bar{\Delta}}{\delta} F_0(D)$$

If chaotic:
$$\sigma_{ch}^2 = \frac{1}{2\pi^2} F_1(D)$$



$$F_n(D) = 1 - \frac{\int_0^D x^n K_0^2(x) dx}{\int_0^\infty x^n K_0^2(x) dx}$$

$$D = \frac{\tau_{\min}}{\tau_\Delta} = \frac{2\pi}{g} \frac{\bar{\Delta}}{\delta}$$

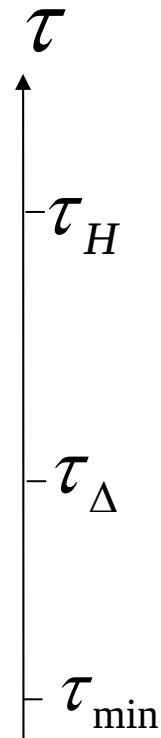
$$g = \frac{\tau_H}{\tau_{\min}} \quad \text{”dimensionless conductance”}$$

If $\tau_{\min} \rightarrow 0 \Rightarrow F_n = 1$ universal fluctuations (random matrix theory)

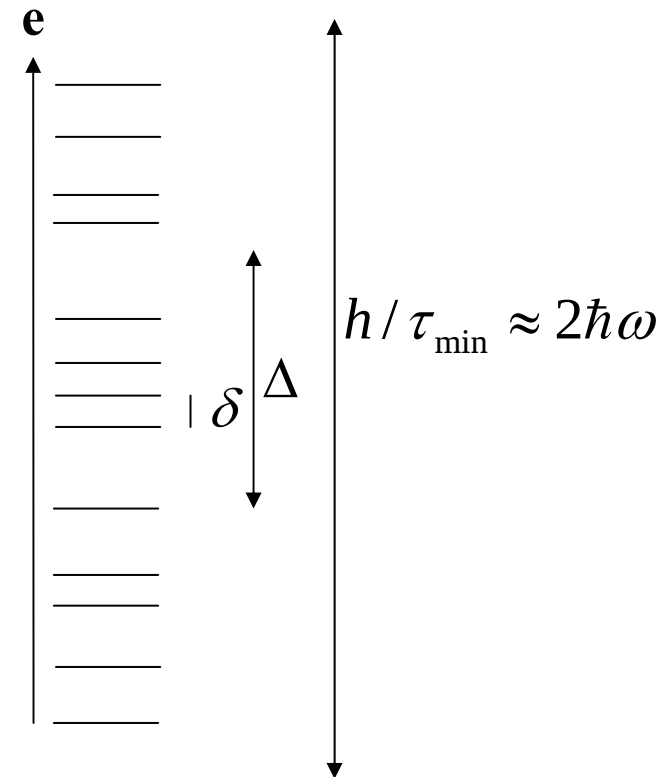


III. Fluctuations of pairing

Three time scales:



Corresponding energy scales:

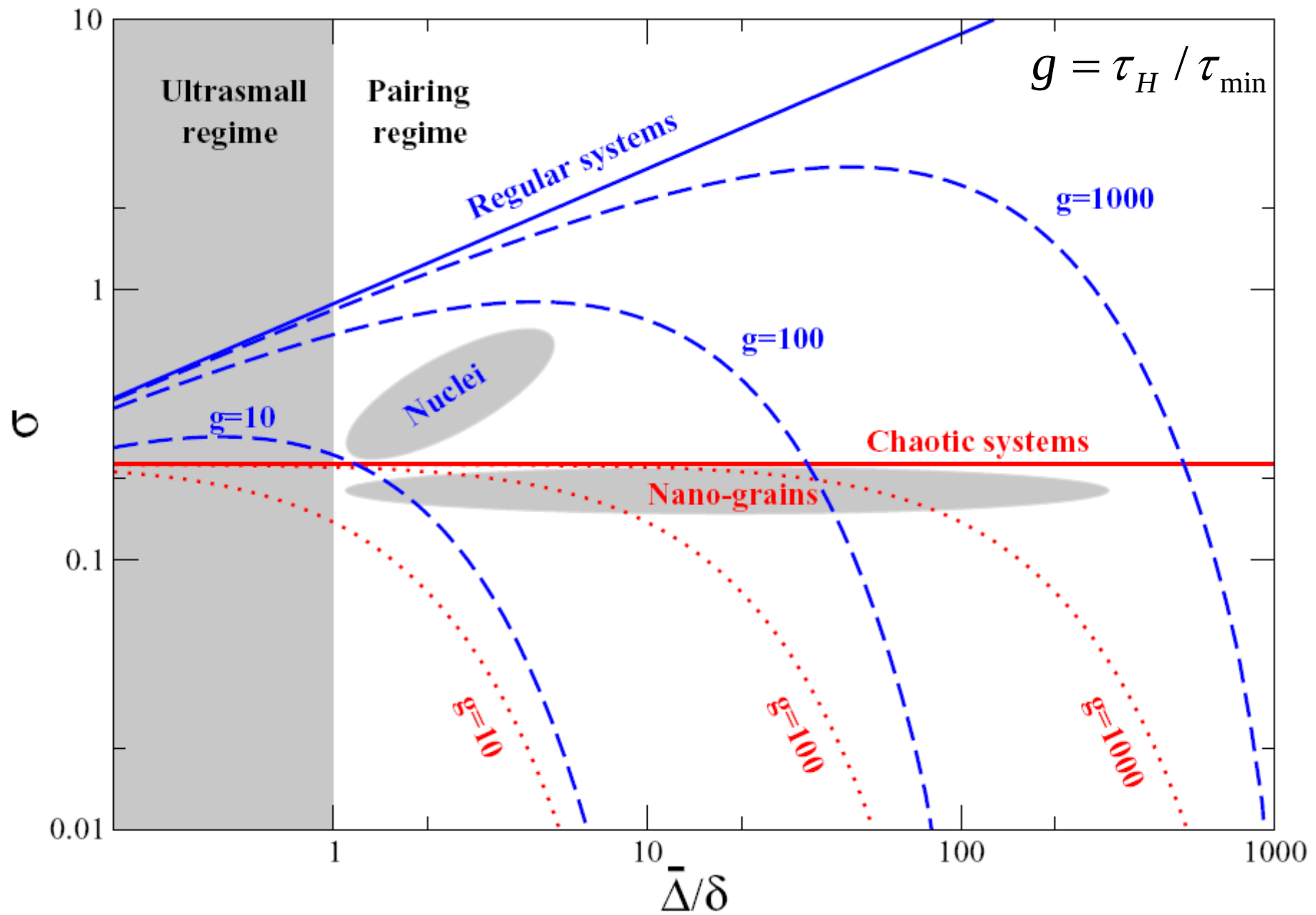


Anderson criterion:
No superconductivity if

$$\Delta > \delta \iff \tau_\Delta \ll \tau_H$$



Generic behavior of pairing fluctuations



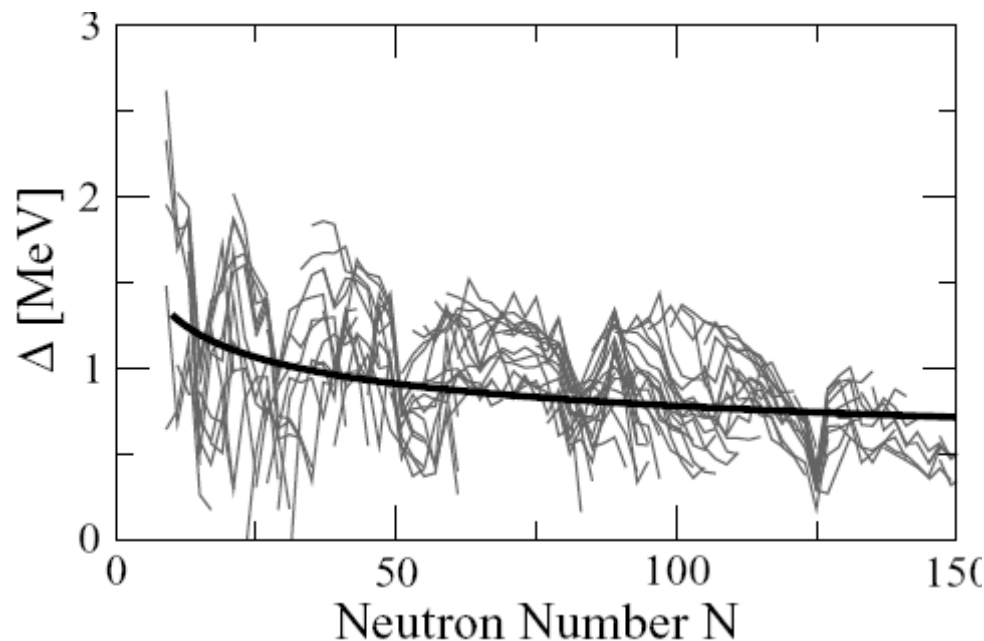
IV. Nuclear pairing gap

Pairing gap estimated from odd-even mass difference.

Three-point formula with N odd:

$$\Delta(N) = -0.5(B(N-1, Z) - 2B(N, Z) + B(N+1, Z))$$

gives minimum contribution from mean field (s.p. spectrum) [1].



Fit gives:

$$\bar{\Delta} = \frac{2.7}{A^{1/4}} \text{MeV}$$

If also even N included:

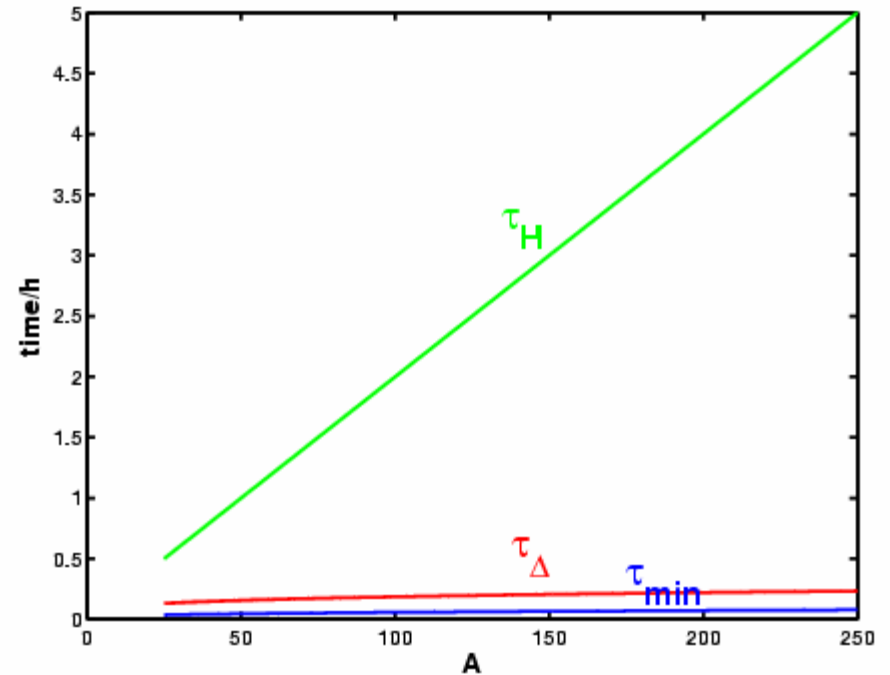
$$\bar{\Delta} = \frac{12}{A^{1/2}} \text{MeV}$$

But includes also
 A -dependence of
mean field

IV. Nuclear pairing gap

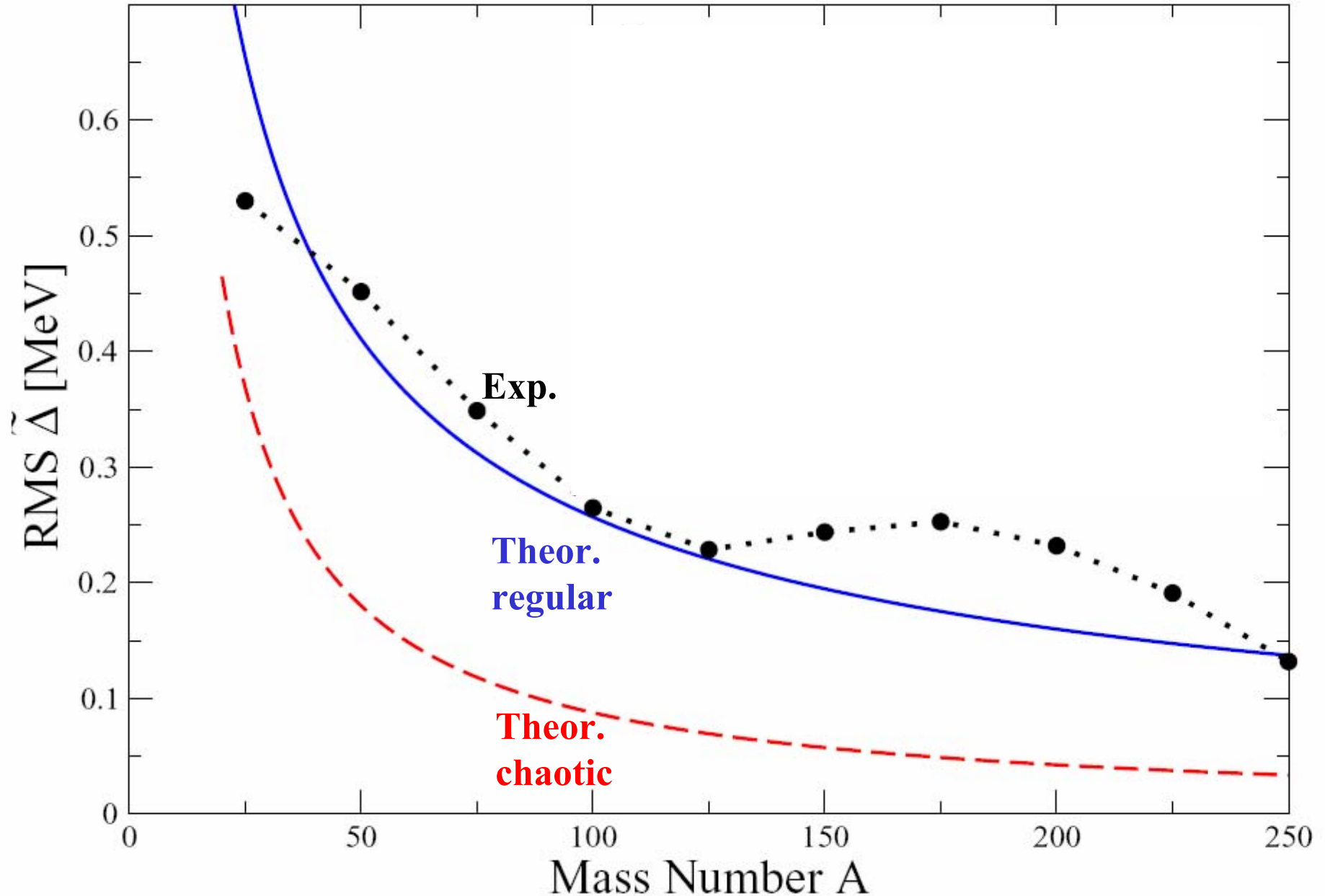
Estimates of time scales for nuclei

$$\begin{aligned}\hbar/\tau_{\Delta} &= \bar{\Delta} = \frac{2.7}{A^{1/4}} \text{ MeV} \\ h/\tau_H &= \delta = 50/A \text{ MeV} \\ h/\tau_{\min} &= 80A^{-1/3} \text{ MeV}\end{aligned}$$



$$\Rightarrow D = \tau_{\min} / \tau_{\Delta} = 0.21A^{1/12} = 0.27 - 0.33$$

IV. Fluctuations of nuclear pairing gap



V. Applications to other finite fermi systems



V.a Nanosized metallic grains

- Discrete exc. spectrum
- Irregular shape of grain $\Rightarrow$ chaotic dynamics
 - No symmetries – only time-rev. symm.
 - Energy level statistics described by GOE
- Excitation gap – pairing gap ($\gg \delta$) observed for even N
- Applied B-field $\Rightarrow$ gap disappears

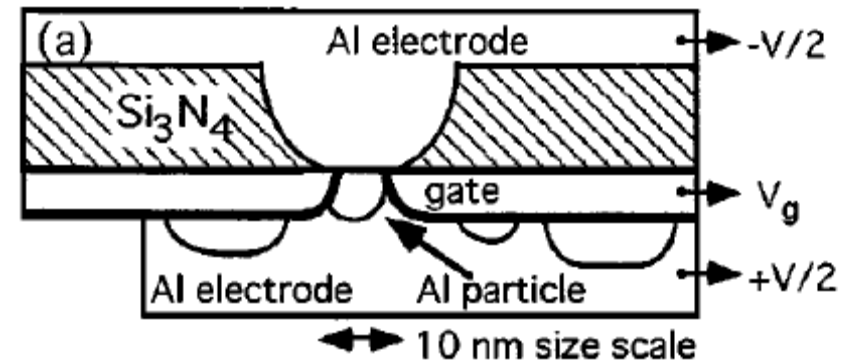
$$N \sim 10^3 - 10^5$$

$$\bar{\Delta} \approx 0.38 \times 10^{-3} \text{ eV} \quad \delta = 21/N \text{ eV} \quad g = 2.6N^{2/3}$$

$$D = \frac{2\pi}{g} \frac{\bar{\Delta}}{\delta} \approx 0.0004 \Rightarrow F_1 \approx 1$$

$\Rightarrow$ Universal pairing fluctuations:

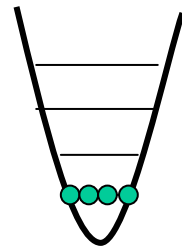
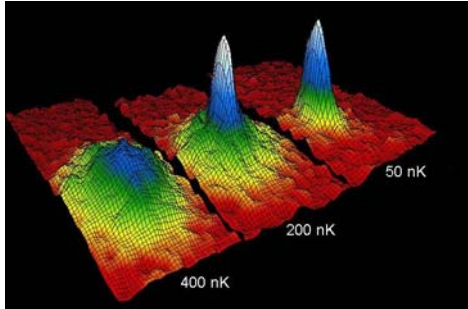
$$\sigma_{ch}^2 = \frac{1}{2\pi^2}$$



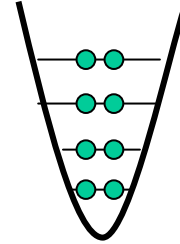
V.b Ultracold fermionic gases

Trapped atomic quantum gases of bosons or fermions

$T \approx 0$



Bose condensate



Degenerate fermi gas

gives possibilities to study new phenomena
in physics of finite many-body systems

Neutral atoms: # electrons = # protons
 $\Rightarrow$ # neutrons determines quantum statistics

e.g. : ${}^6\text{Li}_3$ fermionic
 ${}^7\text{Li}_4$ bosonic



V.b Ultracold fermionic gases

Atom-atom interaction is short-ranged (1-10 Å) and much smaller than interparticle range ($\sim 10^{-6}$ m) (dilute gas)

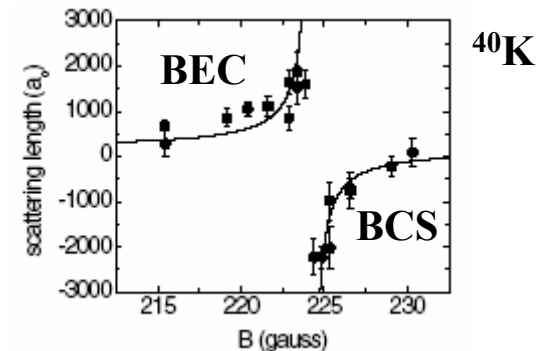
⇒ Approximate int. with:

$$V(r_1 - r_2) = 4\pi \frac{\hbar^2 a}{m} \delta^{(3)}(r_1 - r_2)$$

a =scattering length (s-wave)

Via Feshbach resonance one can experimentally control size and sign of interaction (via external magnetic field):

Two free experimental parameters:
Particle number and interaction strength



C.A. Regal, D.S. Jin,
PRL 90 (2003) 230404

V.b Ultracold fermionic gases

In dilute BCS region:

$$\bar{\Delta} = (2/e)^{7/3} E_F \exp\left(-\frac{\pi}{2k_F|a|}\right)$$
$$\delta = \frac{2E_F}{3N} \quad g = \frac{1}{3}(3N)^{2/3}$$

Recent experiments [1] using ${}^6\text{Li}$ reach $k_F|a| = 0.8$
and about 10^6 atoms gives:
negligible fluctuations of the pairing gap

However, for example, for $k_F|a| = 0.2$
and about 10^3 atoms gives:
large fluctuations of the pairing gap, $\sigma_{\text{reg}} \approx \bar{\Delta} / \delta$



SUMMARY

- Periodic orbit theory of pairing
- Generic description of fluctuations of pairing gap in finite fermi systems
- Accurate, parameter free description of fluctuations of *nuclear* pairing gaps
- Prediction of universal fluctuations in nanosized *metallic grains*
- Estimates of pairing fluctuations in ultracold *Fermi gases*

