

Masses and Fission Barriers of Atomic Nuclei

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Program:

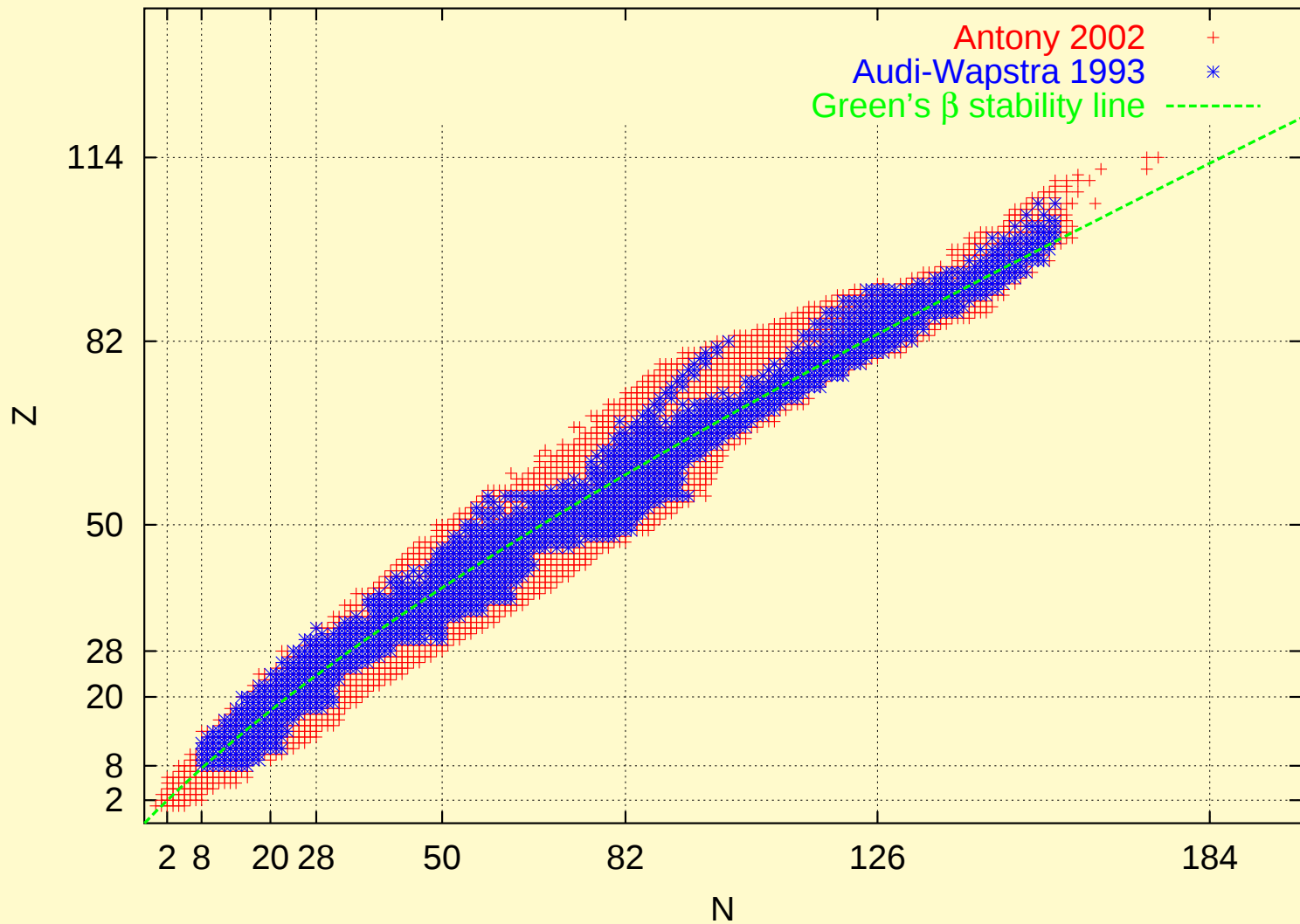
- Introduction
- Leptodermus expansion of energy functional
- Role of different components of the liquid drop model
- Lublin – Strasbourg Drop
- Parameterization of nuclear shapes
- Potential energy surfaces of fissioning nuclei
- Fission barriers and saddle point masses
- Shell and pairing energies
- Summary and conclusions

Liquid drop models:

- C.F. von Weizsäcker, Z. Physik **96**, 431 (1935).
- H.A. Bethe, R.F. Bacher, Rev. Mod. Phys. **8**, 82 (1936).
- L. Meitner, O.R. Frisch, Nature **143**, 239 (1939).
- N. Bohr, J.A. Wheeler, Phys. Rev. **56**, 426 (1939).
- D.L. Hill, J.A. Wheeler, Phys. Rev. **89**, 1102 (1953).
- W.D. Myers, W.J. Świątecki, Nucl. Phys. **81**, 1 (1966).
- H. von Groote, E. Hilf, Nucl. Phys. **A129**, 513 (1969).
- K. Pomorski, J. Dudek, Phys. Rev. **C67**, 044316 (2003).

Other models: droplet, Yukawa+exp., FRDM, TF, ETFSI

Chart of known isotopes



Audi-Wapstra – 1654, Antony – 2766 isotopes with $Z \geq 8$ and $N \geq 8$.

Nuclear energy functional

The one-body density of a nucleus can be expressed as

$$\rho = A \int \int \dots \int \Psi^* \Psi d\tau_2 \dots d\tau_A$$

and the corresponding energy-density as

$$\eta = \int \int \dots \int \Psi^* \hat{H} \Psi d\tau_2 \dots d\tau_A .$$

The total energy of the nucleus is

$$B = \int_V \eta d^3r$$

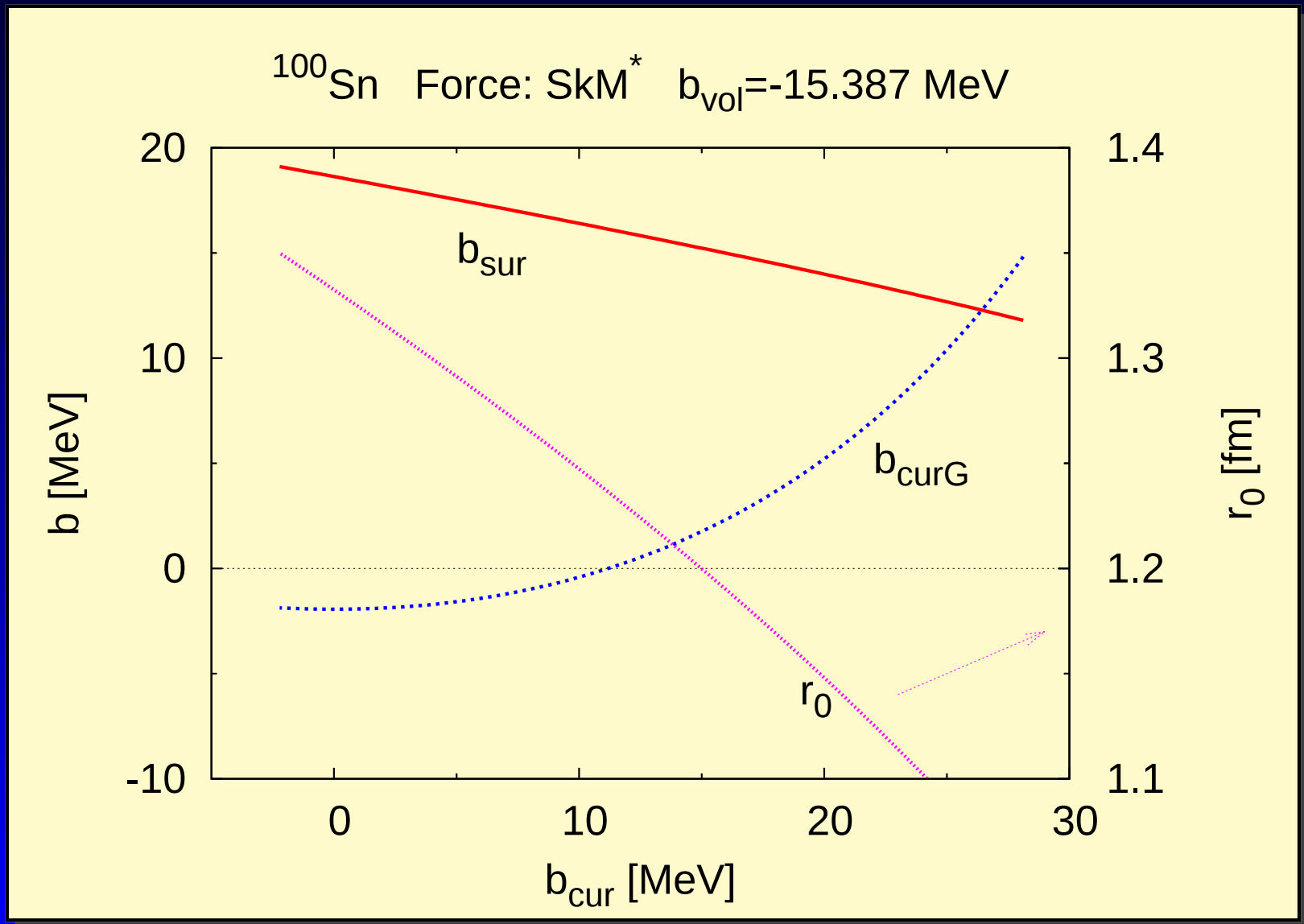
Leptodermous expansion

The total energy integral can be decomposed around diffused around a nuclear surface Σ :

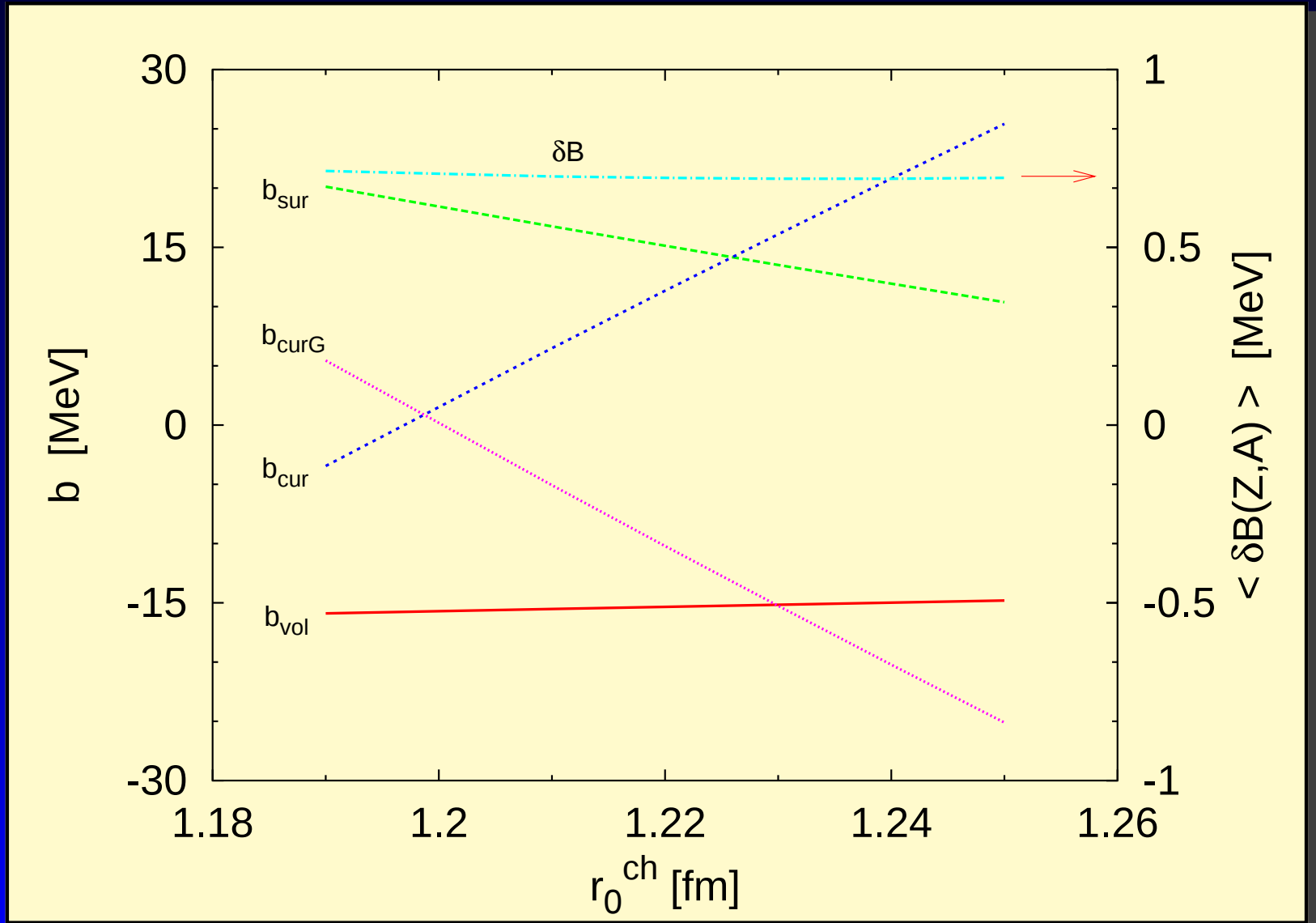
$$\begin{aligned} B &= b_{\text{vol}} A + \oint_{\Sigma} d\sigma \int_0^{\infty} (\eta - b_{\text{vol}} \rho) dr_{\perp} \\ &= b_{\text{vol}} A + \gamma^{(0)} \oint_{\Sigma} d\sigma + \gamma'_{\kappa} a \oint_{\Sigma} \kappa d\sigma \\ &\quad + \frac{1}{2} \gamma''_{\kappa\kappa} a^2 \oint_{\Sigma} \kappa^2 d\sigma + \gamma'_{\Gamma} a^2 \oint_{\Sigma} \Gamma d\sigma + \dots, \end{aligned}$$

where $\kappa = \frac{1}{R_1} + \frac{1}{R_2}$ and $\Gamma = \frac{1}{R_1 \cdot R_2}$ are the 1st and the 2nd order (Gauss) curvatures and R_1 and R_2 are the local main radii of the surface.

Leptodermous expansion of the Skyrme energy functional:



Fit of the liquid drop parameters to experimental binding energies of all known isotopes as function of the charge radius:



Lublin – Strasburg Drop model:

$$\begin{aligned} M(Z, N; \text{def}) = & ZM_{\text{H}} + NM_{\text{n}} - b_{\text{elec}} Z^{2.39} \\ & + b_{\text{vol}} (1 - \kappa_{\text{vol}} I^2) A \\ & + b_{\text{surf}} (1 - \kappa_{\text{surf}} I^2) A^{2/3} B_{\text{surf}}(\text{def}) \\ & + b_{\text{cur}} (1 - \kappa_{\text{cur}} I^2) A^{1/3} B_{\text{cur}}(\text{def}) \\ & + \frac{3}{5} e^2 \frac{Z^2}{r_0^{ch} A^{1/3}} B_{\text{Coul}}(\text{def}) - C_4 \frac{Z^2}{A} \\ & + E_{\text{micr}}(Z, N; \text{def}) + E_{\text{cong}}(Z, N) \end{aligned}$$

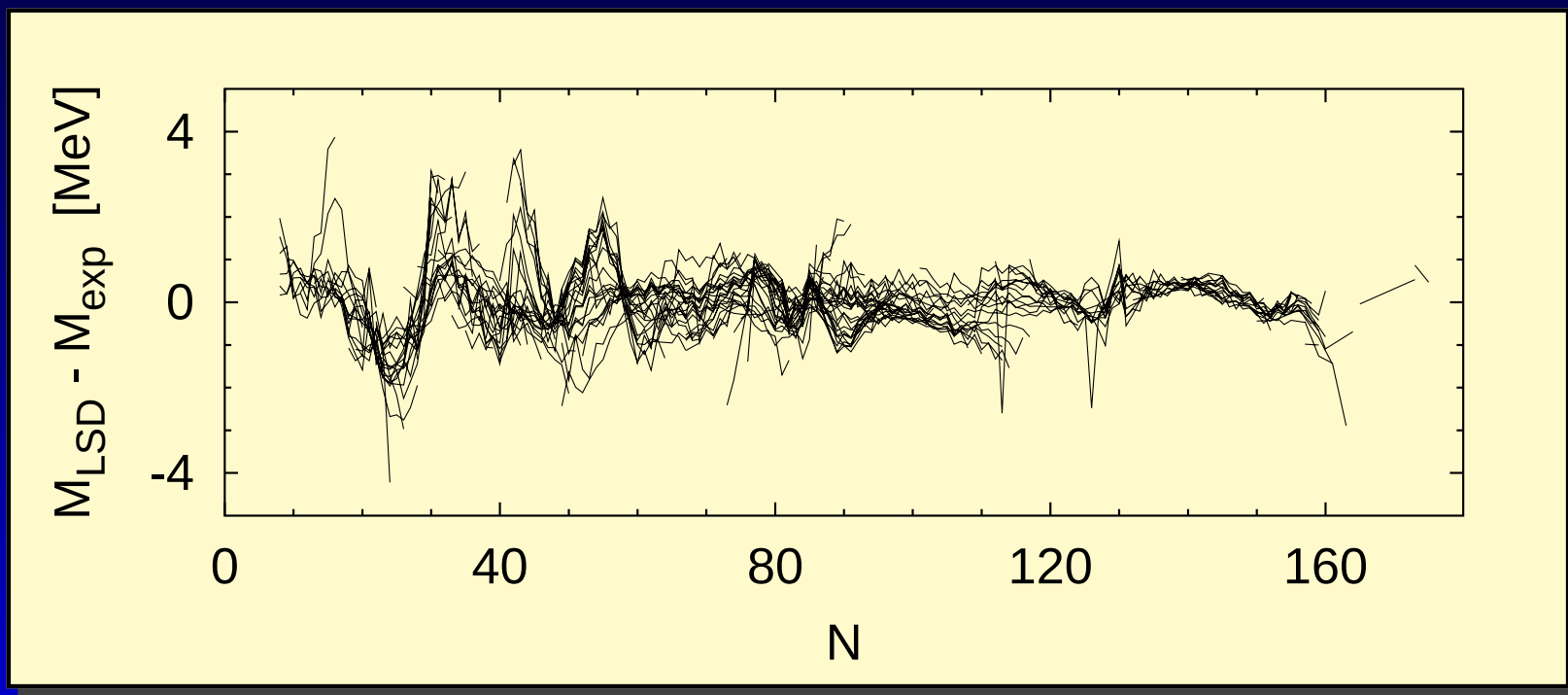
Fit to the 2766 experimental masses:

Term	Units	LDM	LSD
b_{vol}	MeV	-15.8484	-15.4920
κ_{vol}	-	1.8475	1.8601
b_{surf}	MeV	19.3859	16.9707
κ_{surf}	-	1.9830	2.2938
b_{cur}	MeV	0	3.8602
κ_{cur}	-	0	-2.3764
r_0	fm	1.18995	1.21725
C_4	MeV	1.1995	0.9181
δM	MeV	0.732	0.698
$\delta V_{B_{Z>70}}$	MeV	5.58	0.88

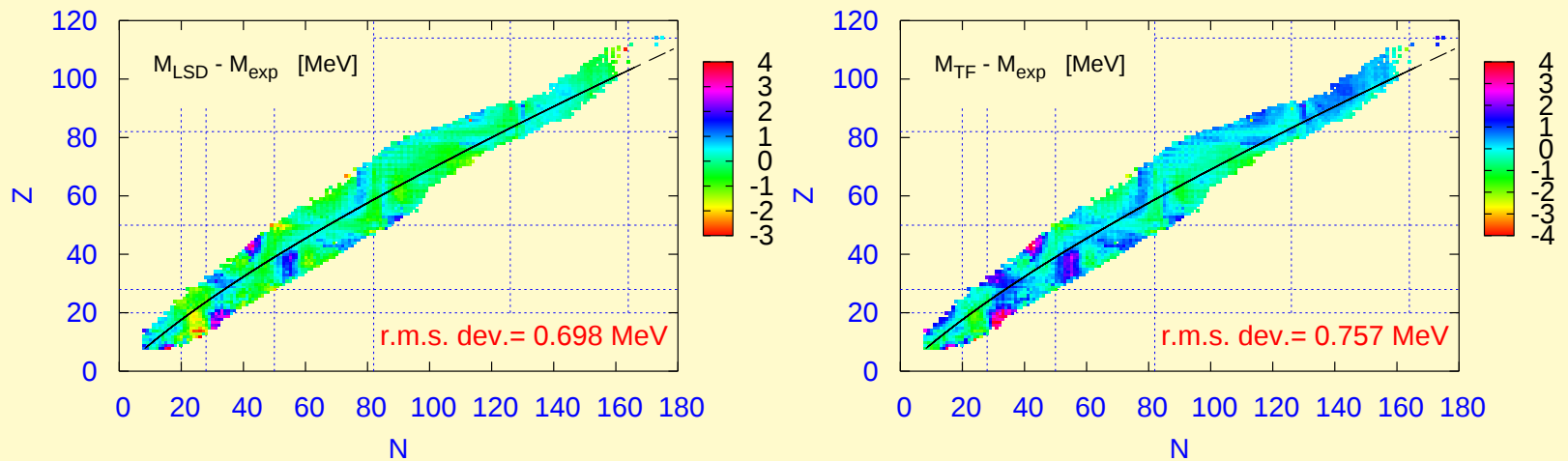
$M_{\text{H}}=7.289034$ MeV; $M_{\text{n}}=8.071431$ MeV; $b_{\text{elec}}=1.433$ eV

K. Pomorski , J. Dudek, Phys. Rev. C67, 044316 (2003).

LSD+YF mass deviation:



Masses of isotopes:

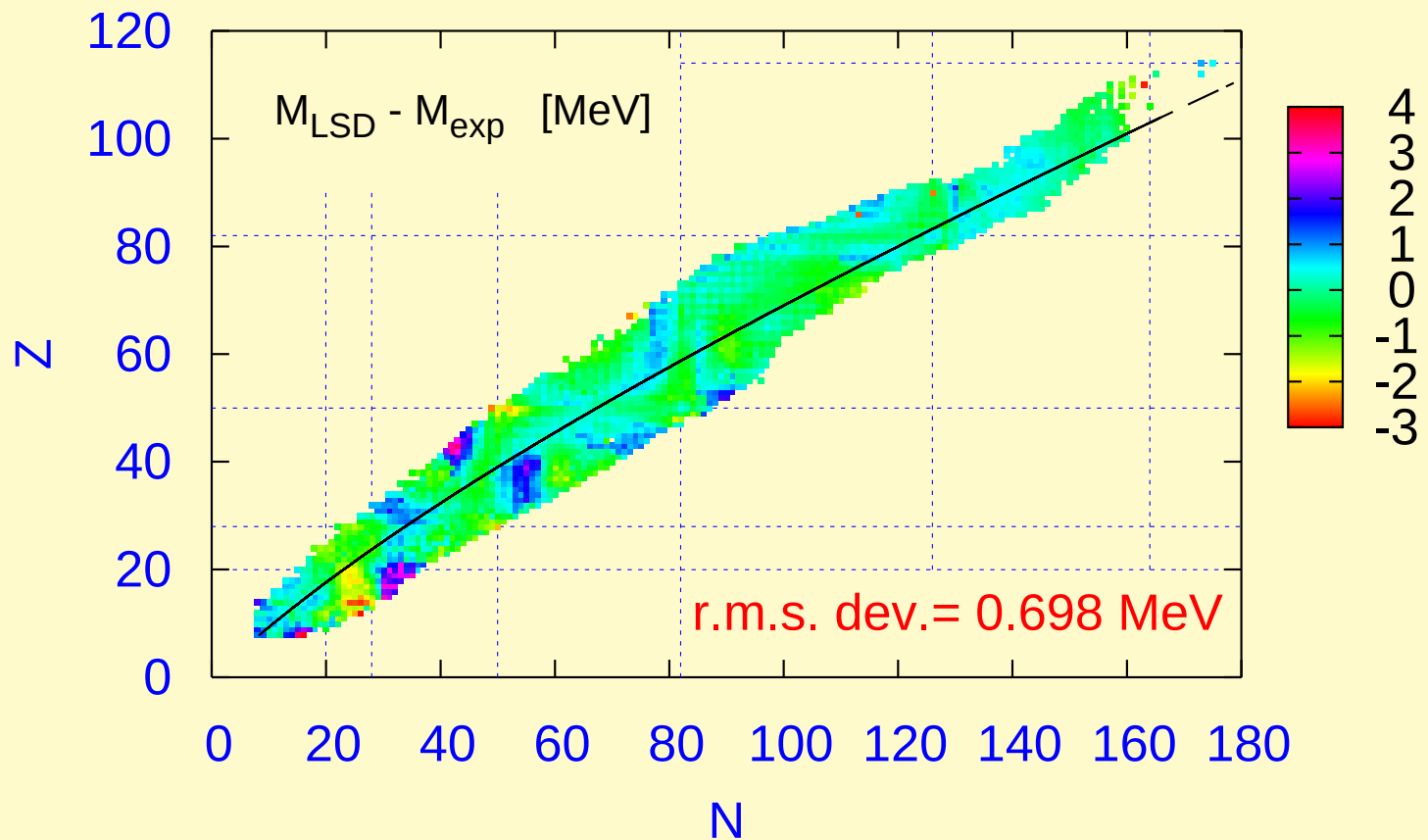


exp: *The mass deviations are plotted for 2766 isotopes taken from the Chart of Nuclides by M.S. Antony, Strasbourg, 2002.*

LSD: *K. Pomorski, J. Dudek, Phys. Rev. **C67**, 044316 (2003).*

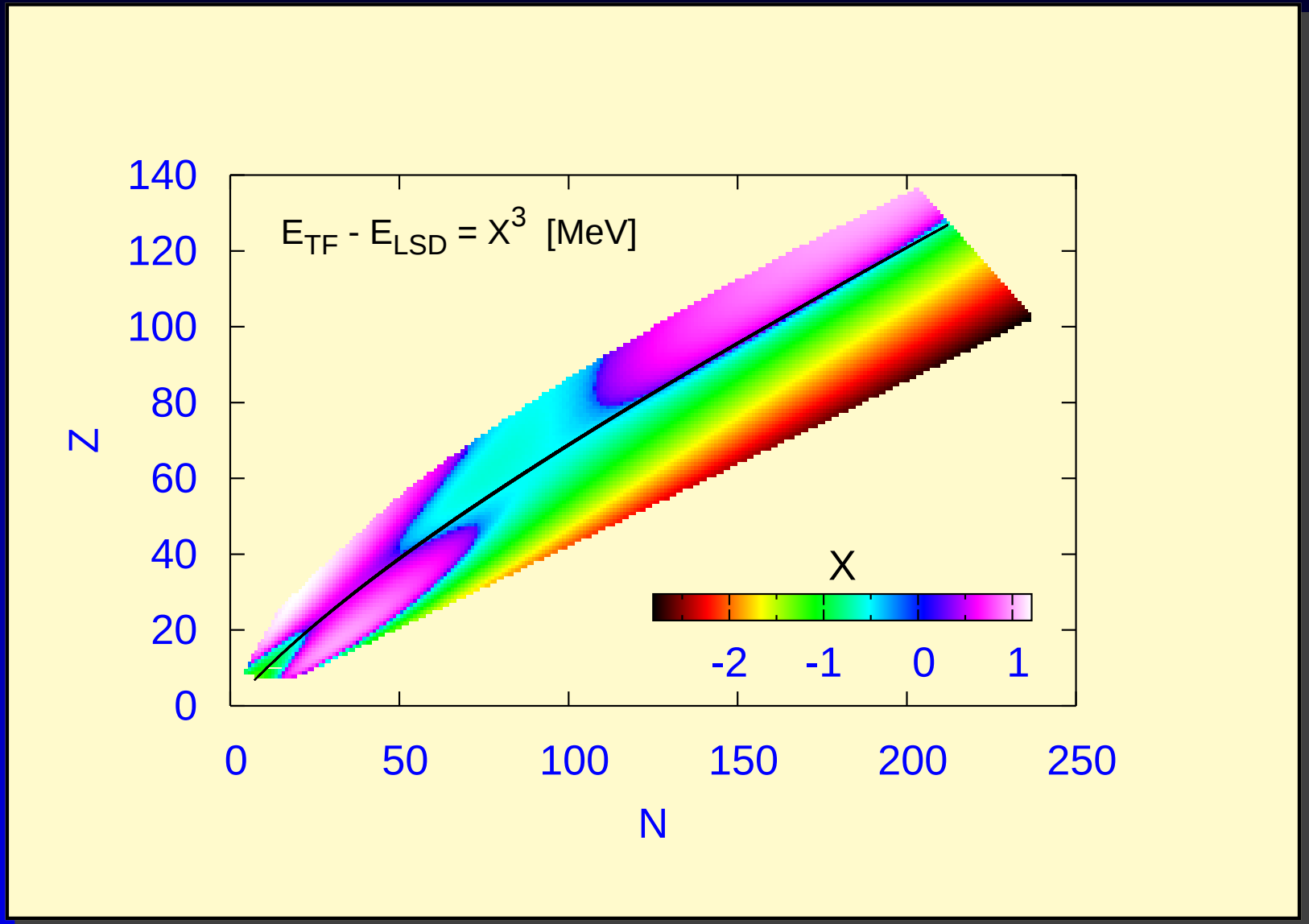
TF: *P. Möller, J.R. Nix, W.D. Myers, W.J. Świątecki, At. Data Nucl. Data Tables **59**, 185 (1995).*

LSD masses:



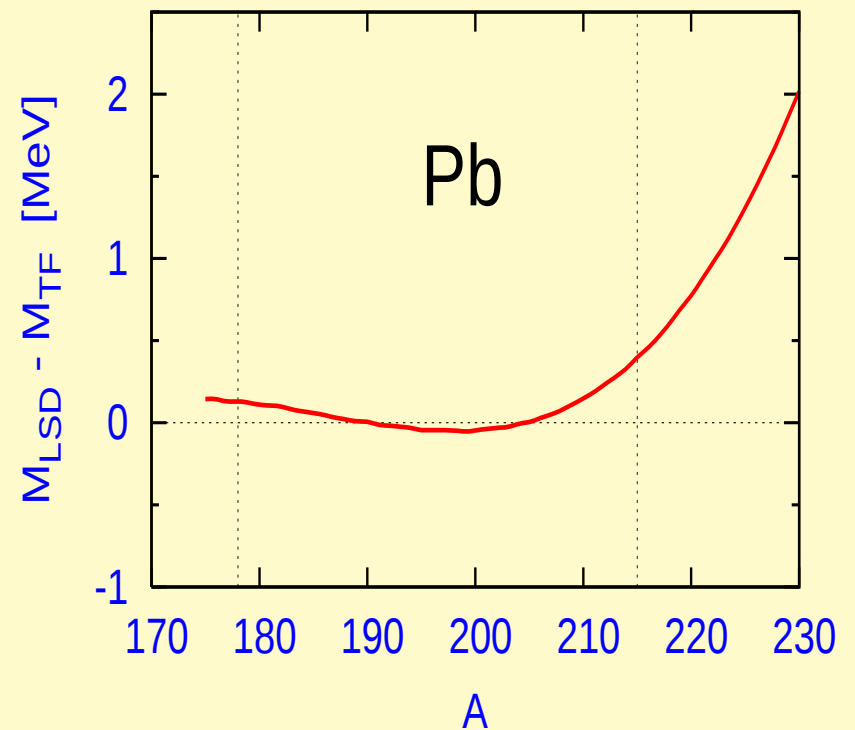
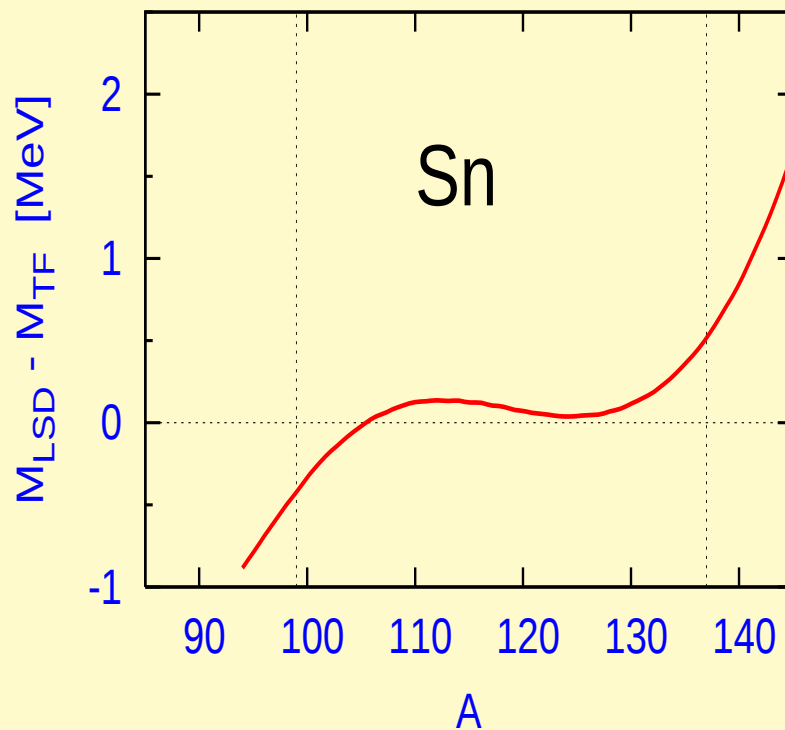
K. Pomorski, J. Dudek, Phys. Rev. C67, 044316 (2003).

TF and LSD mass differences:



Data plotted for 8757 isotopes taken from the mass table of P. Möller et al., *At. Data Nucl. Data Tables* **59**, 185 (1995).

TF and LSD mass differences:



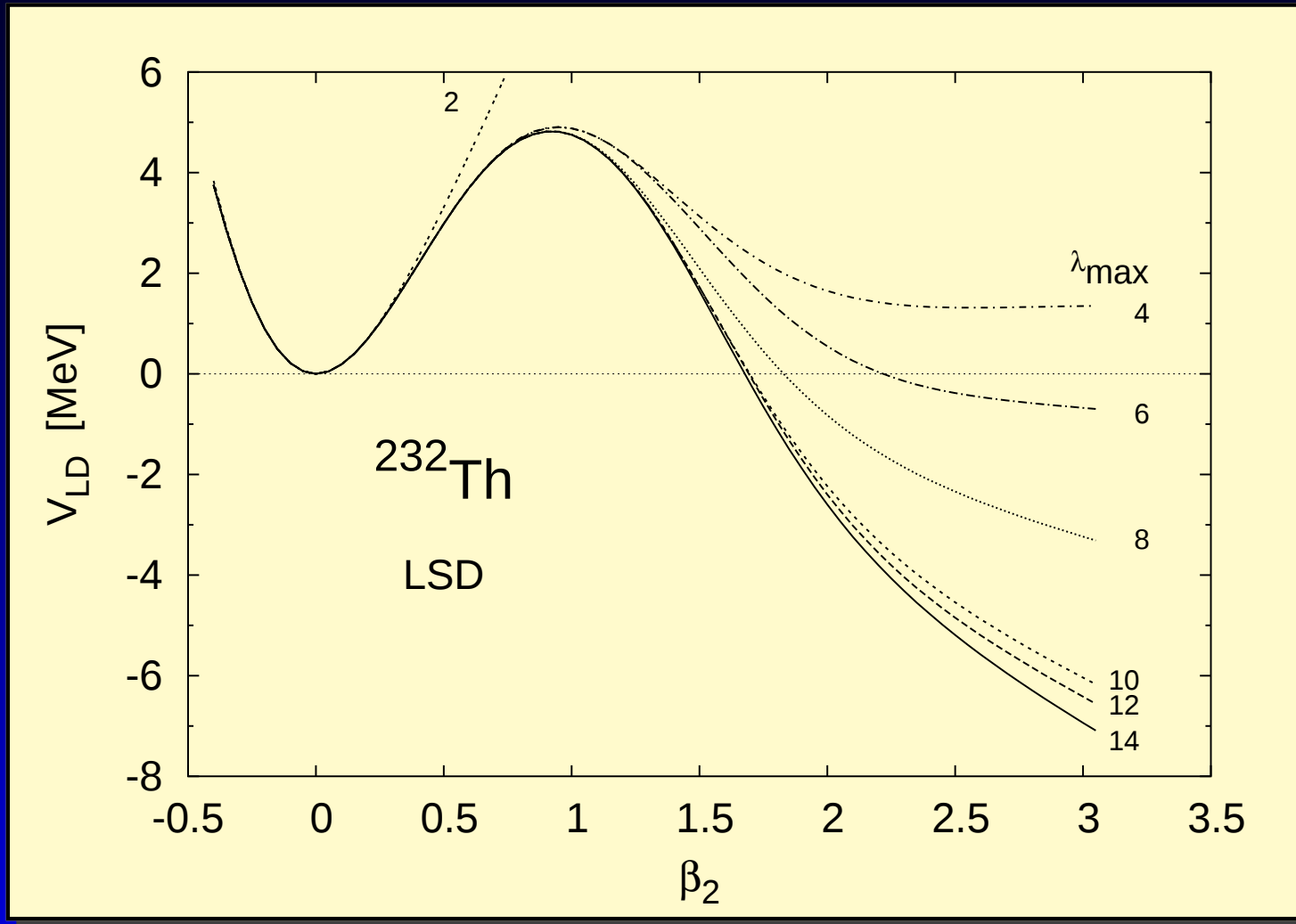
Predictive power of mass formulae

Root means square deviations (δM in MeV) of theoretical and experimental binding energies of 1654 isotopes taken from the Audi-Wapstra tables and 2766 isotopes from the Antony chart:

Model	$\delta M(2766)$	$\delta M(1654)$
LSD(2766)	0.698	0.610
LSD(1654)	0.711	0.600
TF-MS(1654)	0.757	0.655
ETFSI(1888)	0.828	0.738

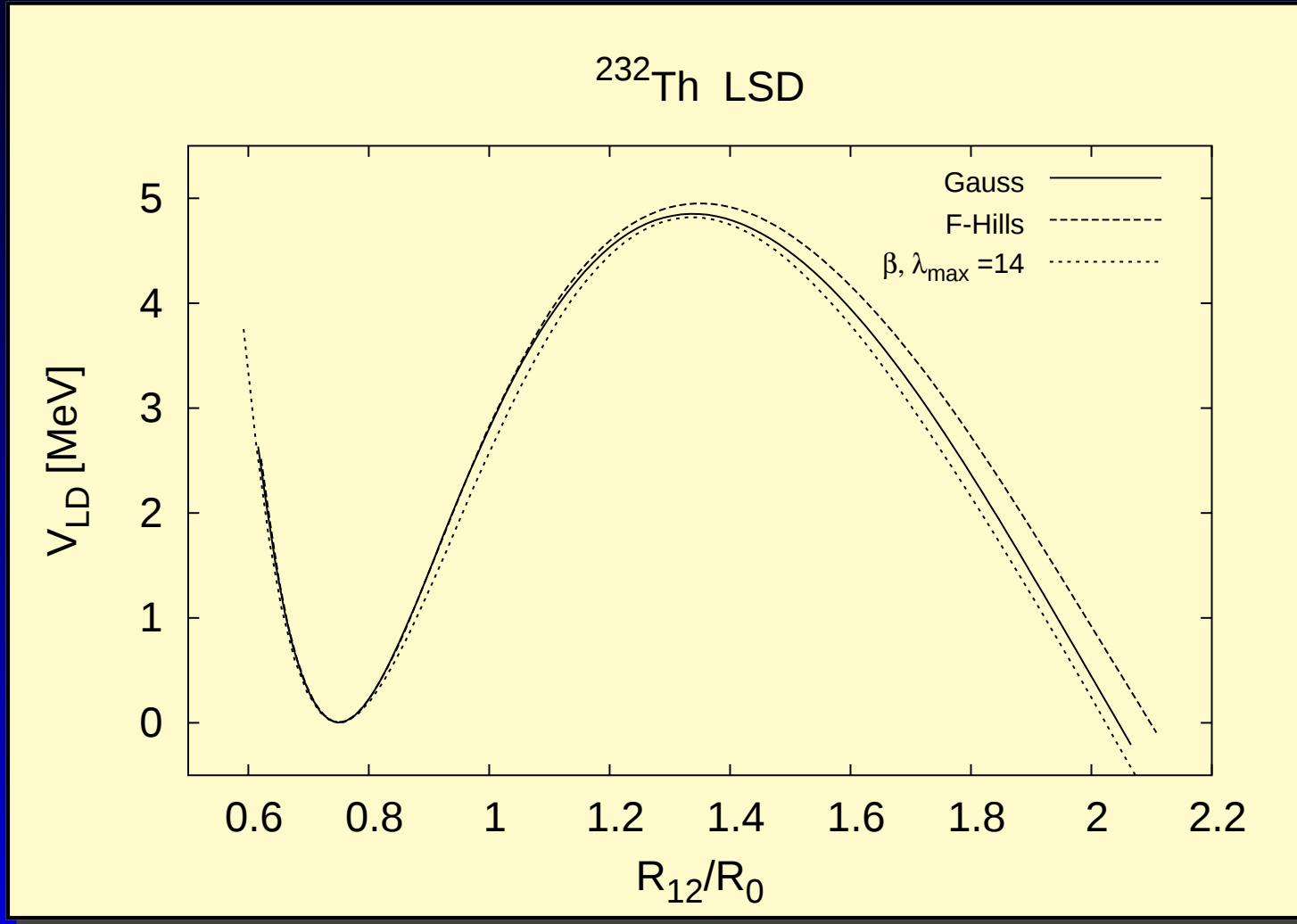
ETFSI: S. Goriely, J.M. Pearson, and F. Tondeur, At. Data and Nucl. Data Tables **77**, 311 (2001).

Spherical harmonics expansion:



$$R(\theta) = R_0 \sum_{\lambda=0}^{\infty} \beta_{\lambda} P_{\lambda}(\cos \theta)$$

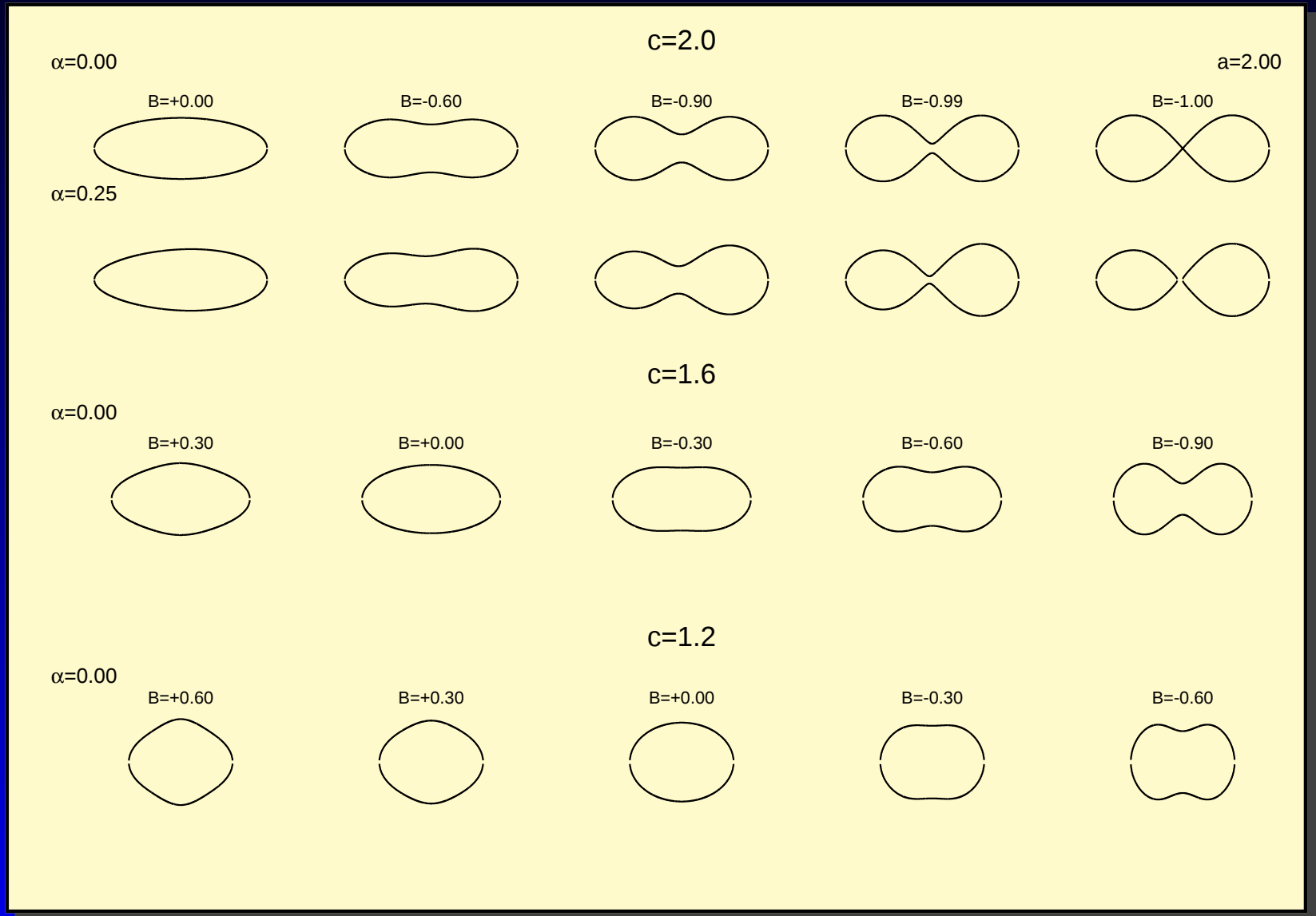
Modified Funny-Hills:



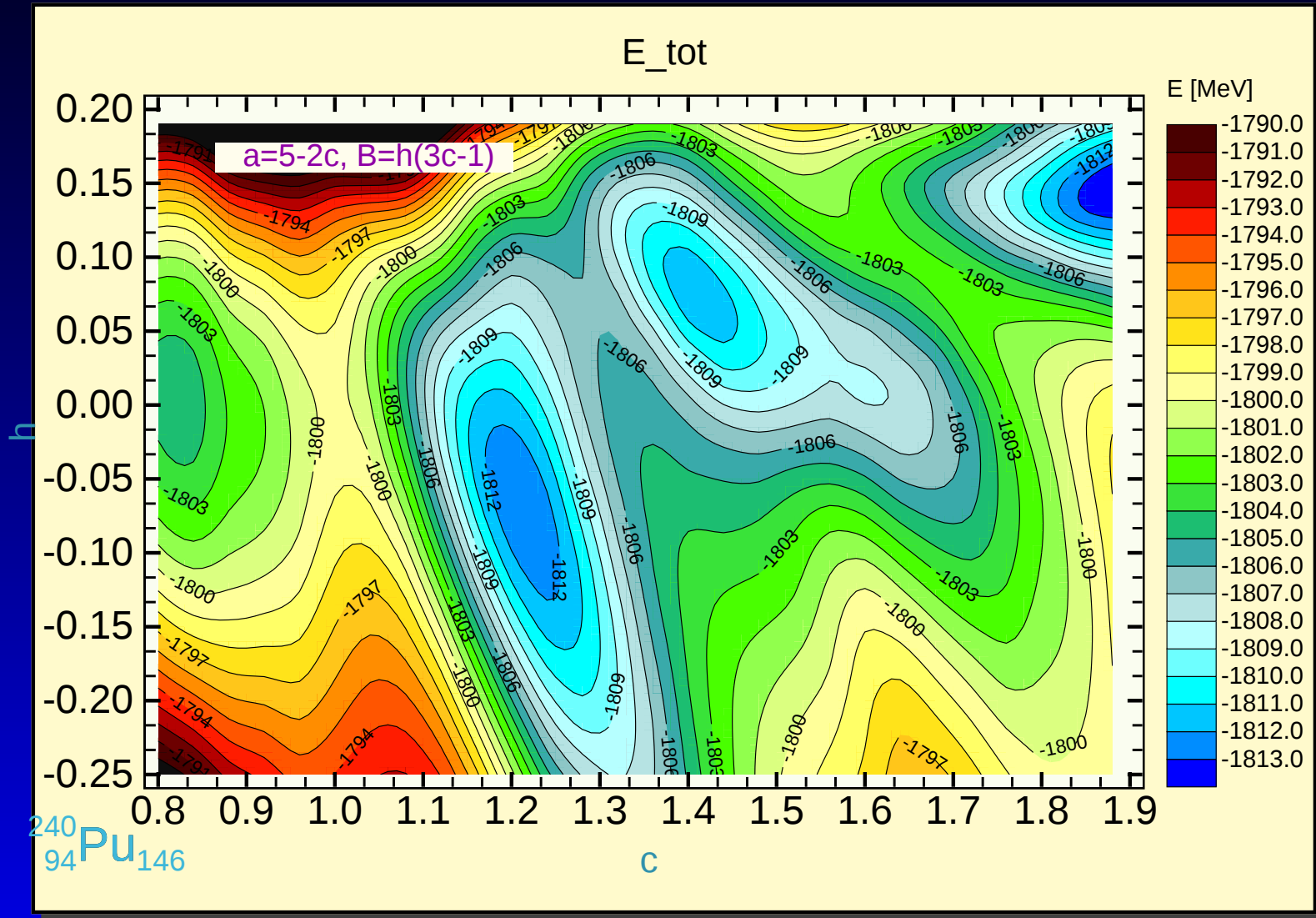
$$\tilde{\rho}_s^2(z) = \frac{R_0^2}{c f(a, B)} (1 - u^2) (1 + \alpha u - B e^{-a^2 u^2}) ,$$

where $u = (z - z_{sh}) / (c R_0)$.

MFH shapes examples:

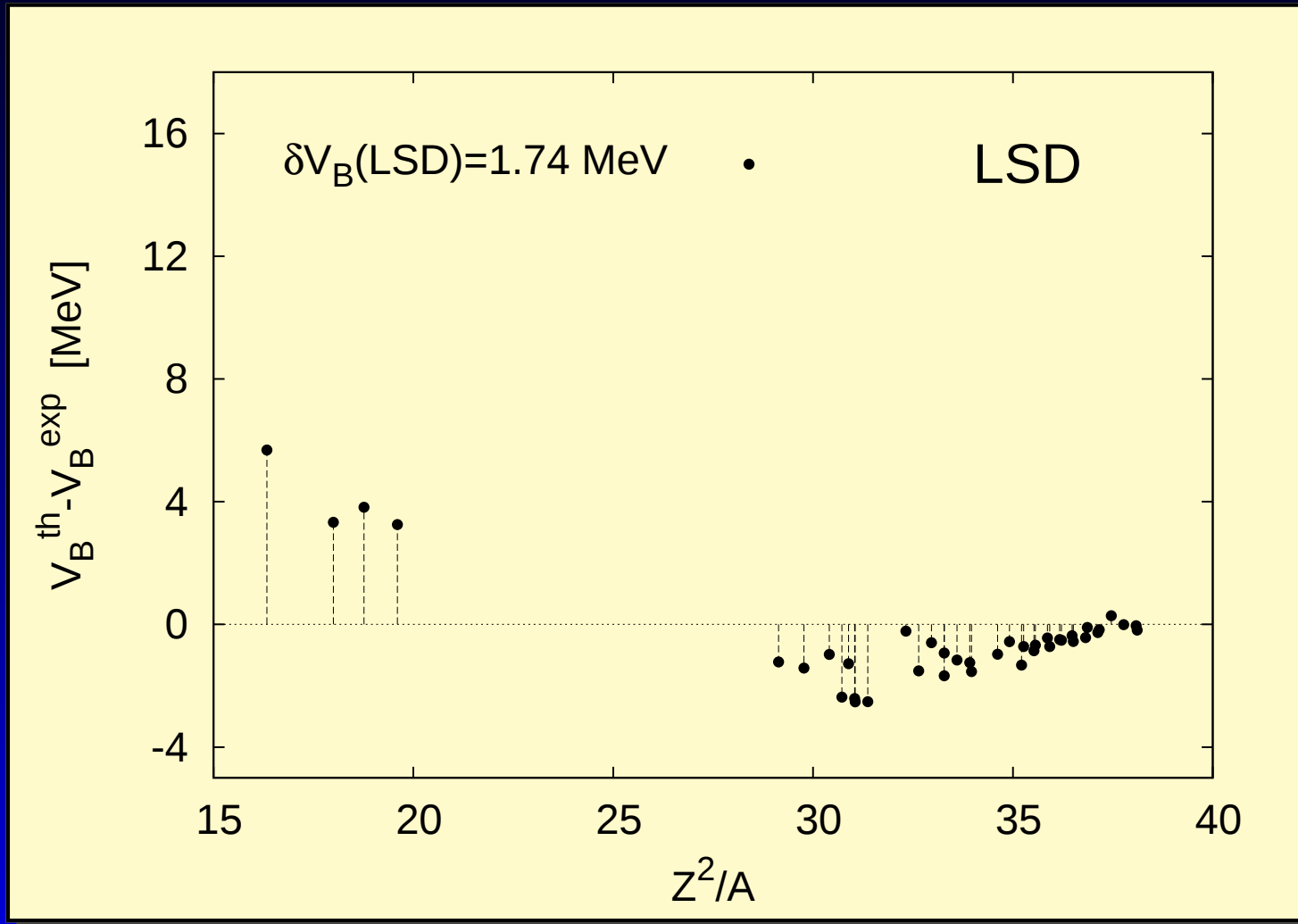


PES of ^{240}Pu in the MFH space:



LSD with shell and pairing corrections obtained with the Yukawa-folded single particle potential.

Fission barrier heights:

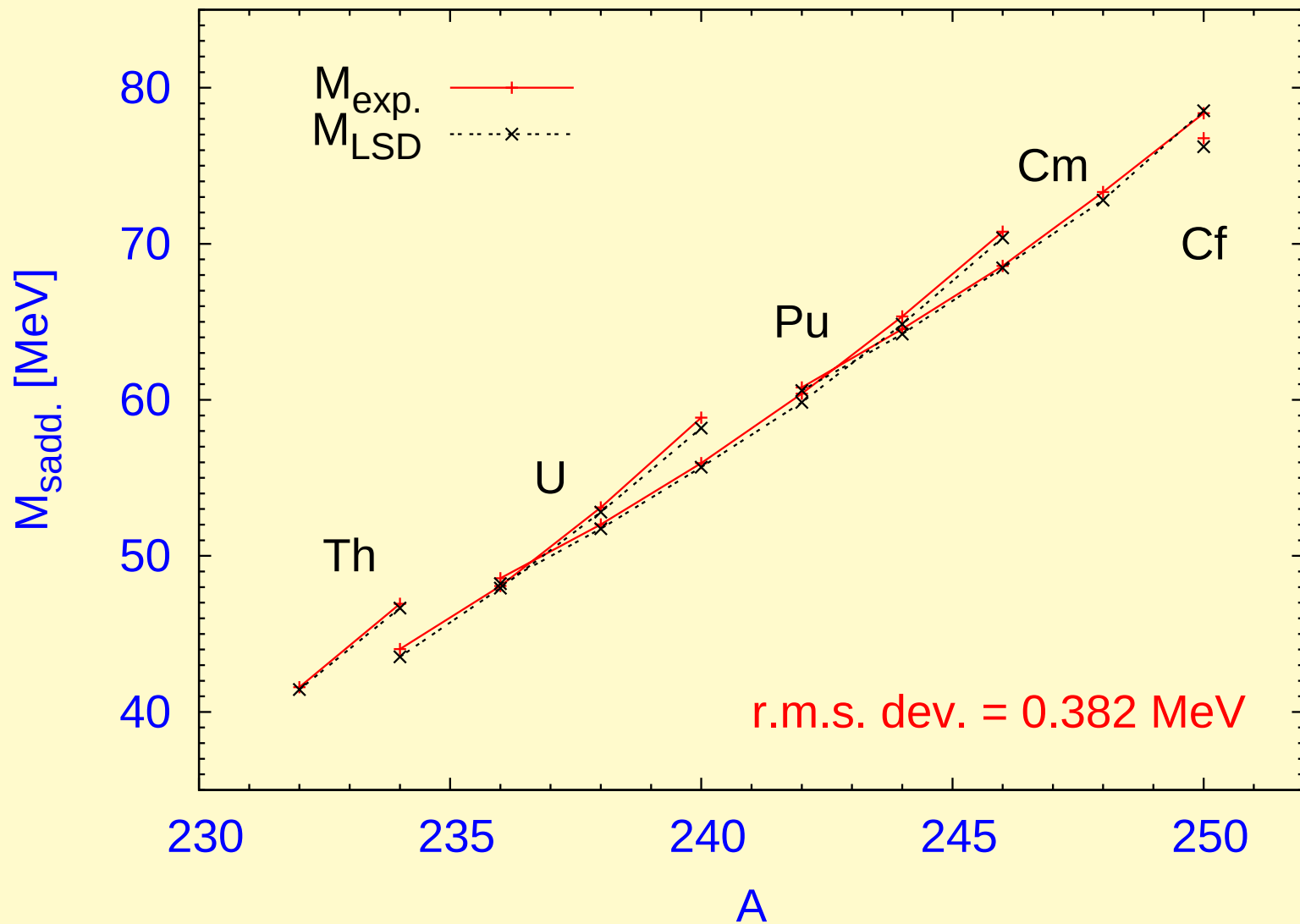


Deformation dependent congruence energy term is included here according to:

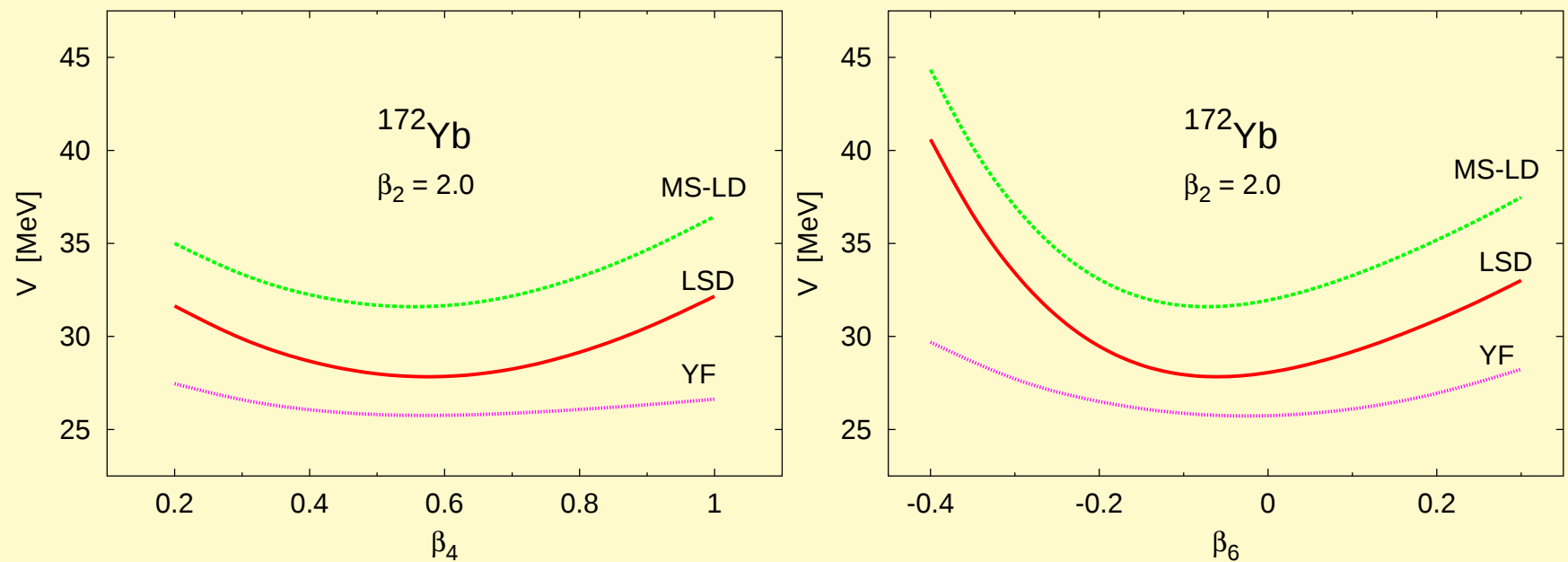
W.D. Myers, J. Świątecki, Nucl. Phys. **A612** (1997) 249.

K. Pomorski, J. Dudek, Int. Journ. Mod. Phys. E13 (2004) 107.

Saddle point masses:



Stiffness at the saddle point:



Shell and pairing corrections:

$$E_{\text{tot}} = E_{\text{macro}} + \delta E_{\text{shell}} + \delta E_{\text{pair}}$$

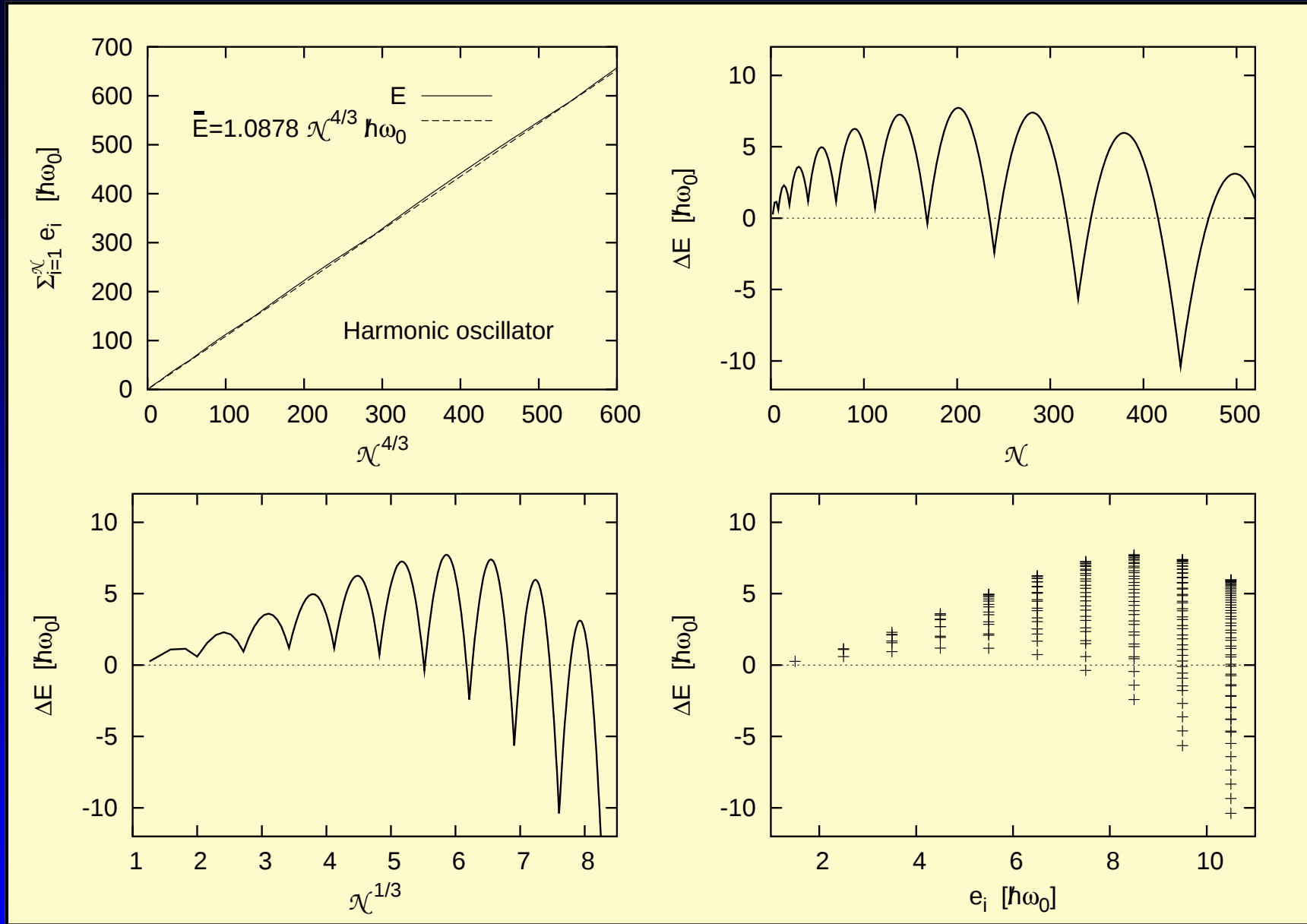
$$\delta E_{\text{shell}} = \sum_{\text{occ}} 2e_{\nu} - \langle \sum_{\text{occ}} 2e_{\nu} \rangle$$

$$E_{\text{pair}} = E_{\text{BCS}} - \sum_{\text{occ}} 2e_{\nu}$$

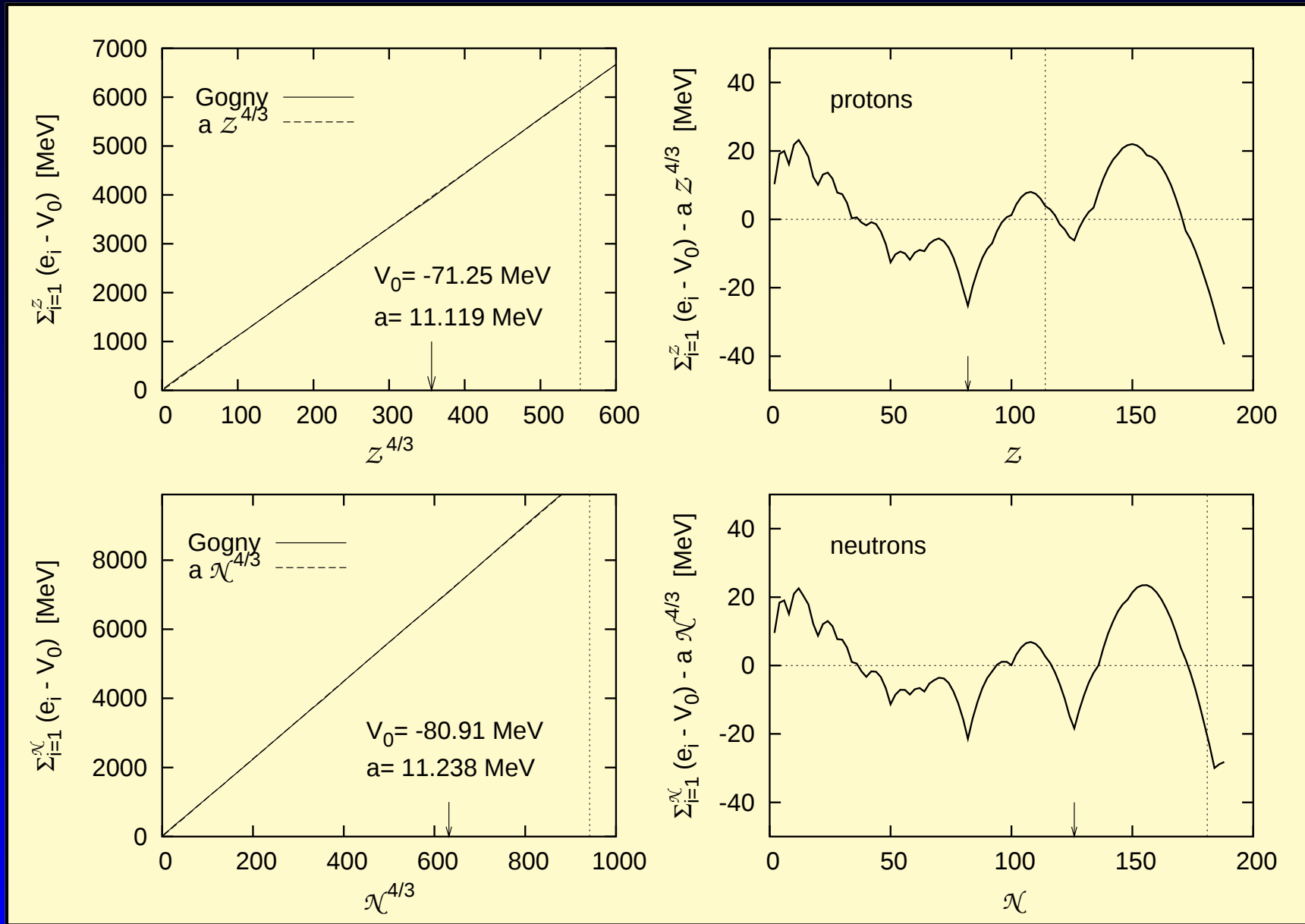
$$\delta E_{\text{pair}} = E_{\text{pair}} - \langle E_{\text{pair}} \rangle$$

$$\delta E_{\text{shell}} + \delta E_{\text{pair}} = E_{\text{BCS}} - \langle E_{\text{BCS}} \rangle$$

Harmonic oscillator s.p. energy sum



Gogny mean-field Hamiltonian:



$\mathcal{N}$ -folding of the s.p. energy sum

$$E_{\text{sum}}(n) = \sum_{\nu=1}^n e_{\nu}$$

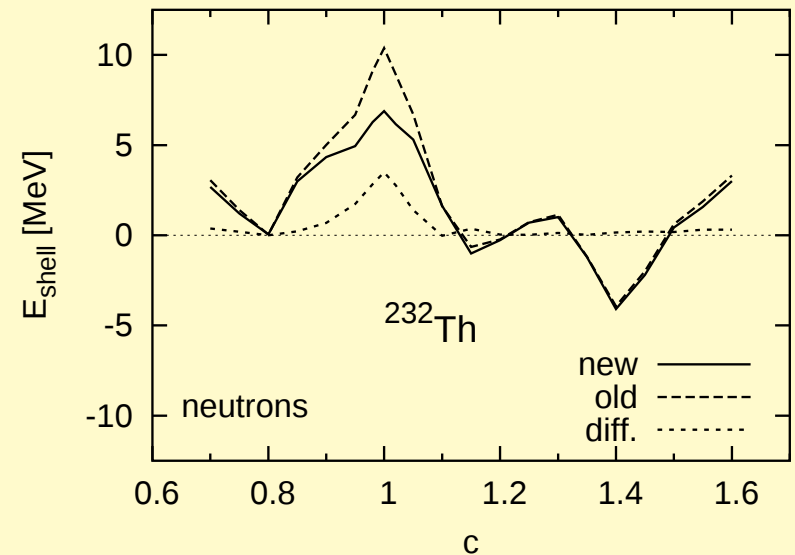
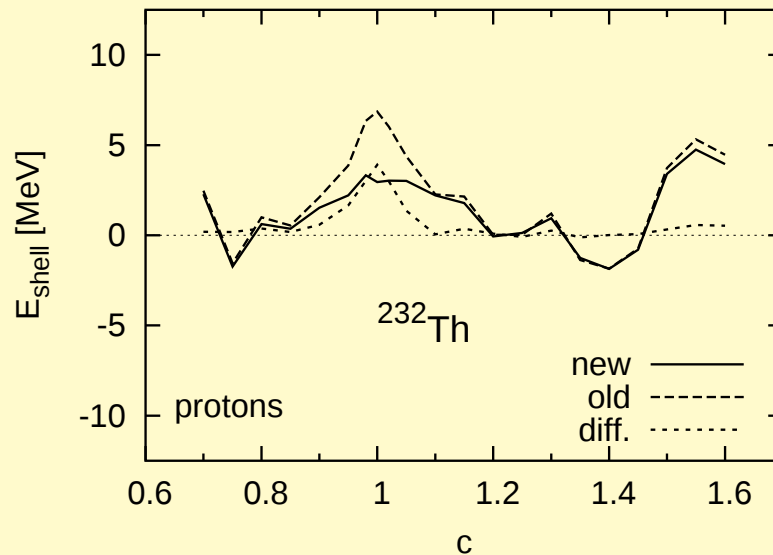
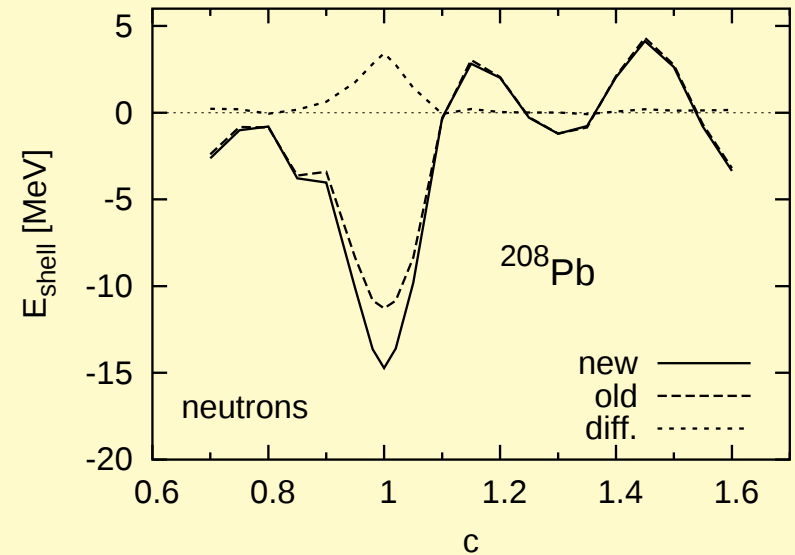
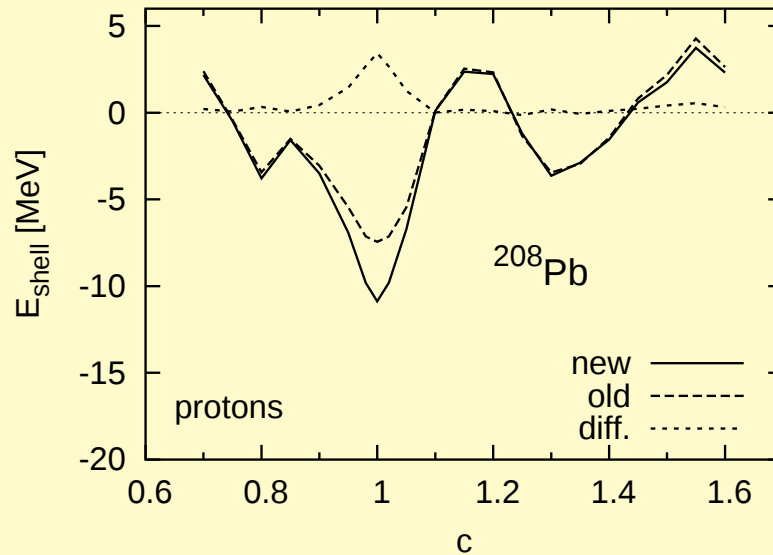
The energy sum smoothed in $\mathcal{N}^{1/3}$ -space is given by:

$$\langle E_{\text{sum}}(\mathcal{N}) \rangle = \sum_{\mathcal{N}_{\min}}^{\mathcal{N}_{\max}} \frac{2 E_{\text{sum}}(n)}{3 n^{2/3}} j\left(\frac{\mathcal{N}^{1/3} - n^{1/3}}{\gamma}\right)$$

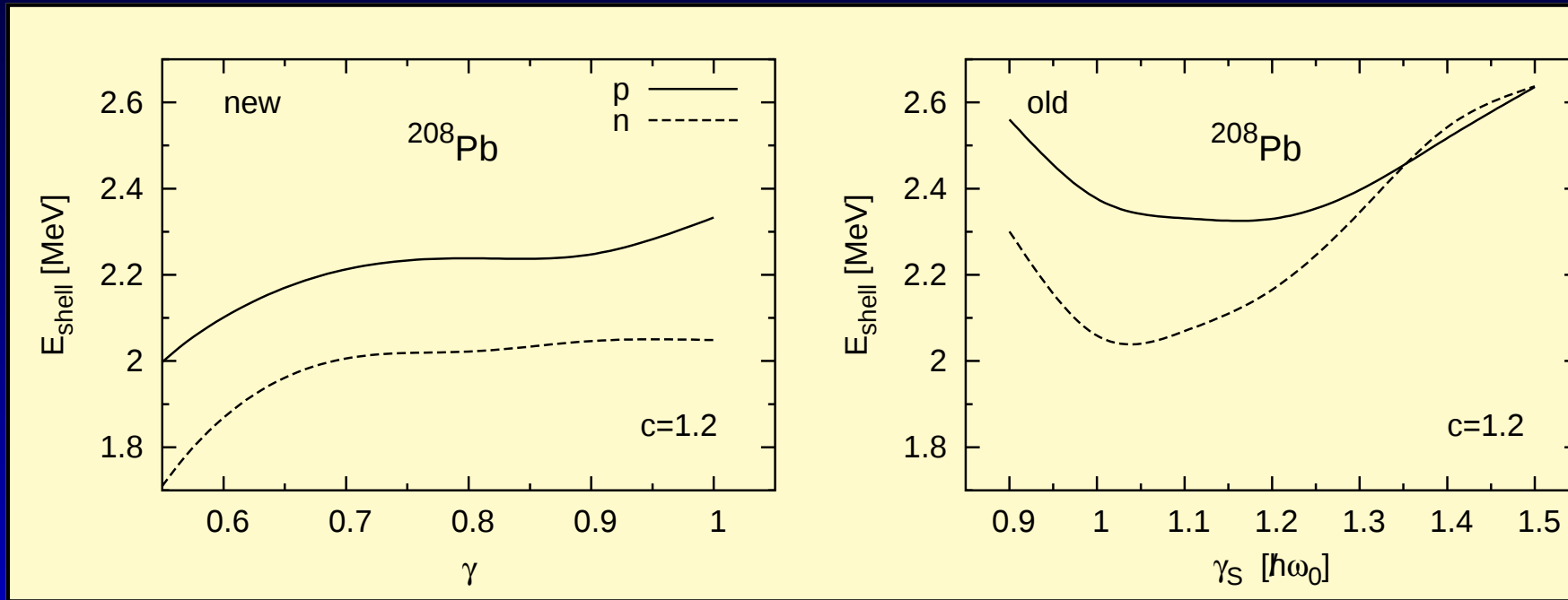
where $\gamma \approx 0.7$ is the smearing width and

$$j(u) = \frac{1}{\sqrt{\pi}} e^{-u^2} \left(\frac{35}{16} - \frac{35}{8} u^2 + \frac{7}{4} u^4 - \frac{1}{6} u^6 \right) .$$

Strutinsky and new shell energy



Plateau condition:



K. Pomorski, Phys. Rev. **C70** (2004) 044306.

$\mathcal{N}$ -folding of the BCS energy:

$$E_{\text{BCS}}(n) = 2 \sum_{\nu} v_{\nu}^2 e_{\nu} - G \sum_{\nu} v_{\nu}^4 - \Delta^2 / G$$

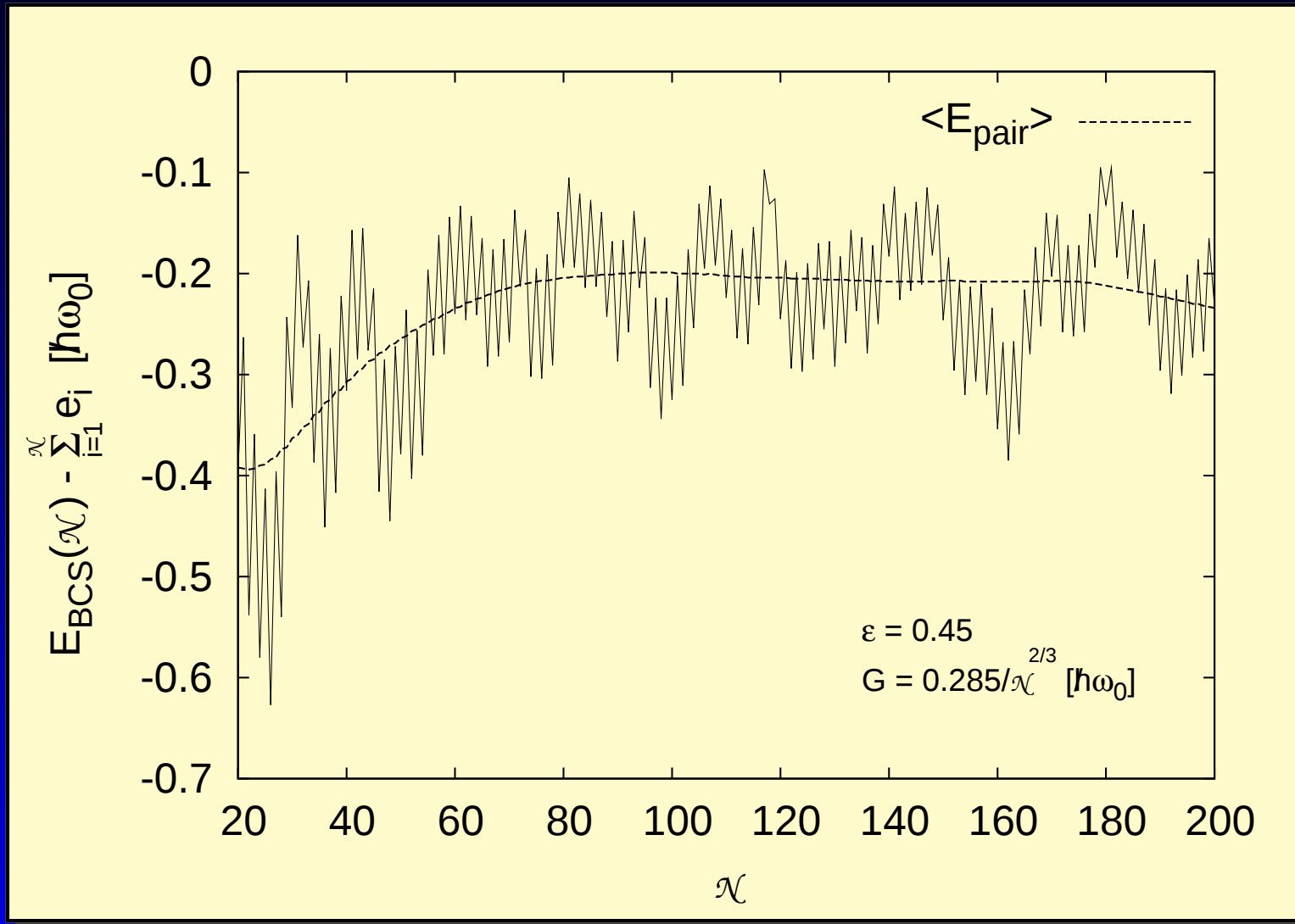
The BCS energy smoothed in $\mathcal{N}^{1/3}$ -space is given by:

$$\langle E_{\text{BCS}}(\mathcal{N}) \rangle = \sum_{n=\mathcal{N}_{\min}}^{\mathcal{N}_{\max}} \frac{2 E_{\text{BCS}}(n)}{3 n^{2/3}} j \left(\frac{\mathcal{N}^{1/3} - n^{1/3}}{\gamma} \right)$$

where

$$j(u) = \frac{1}{\sqrt{\pi}} e^{-u^2} \left(\frac{35}{16} - \frac{35}{8} u^2 + \frac{7}{4} u^4 - \frac{1}{6} u^6 \right) .$$

Smooth pairing energy



K. Pomorski, B. Nerlo-Pomorska, Physica Scripta **T125** (2006) 21.

B. Nerlo-Pomorska, K. Pomorski, Int. Journ. Mod. Phys. **E16** (2007) 328.

Summary and conclusions

- **L**ublin **S**trasburg **D**rop, which apart traditional liquid drop terms contains the curvature energy, describes well the masses of known isotopes.
- The fission barriers as well as the saddle point masses are reproduced well by the LSD model.
- New method of evaluation of the Strutinsky smoothed energy by folding in the particle number space is proposed.
- The $\mathcal{N}$ -folding method can be applied to obtain the average pairing energy.
- Large scale calculation of the nuclear mass table is planned.