

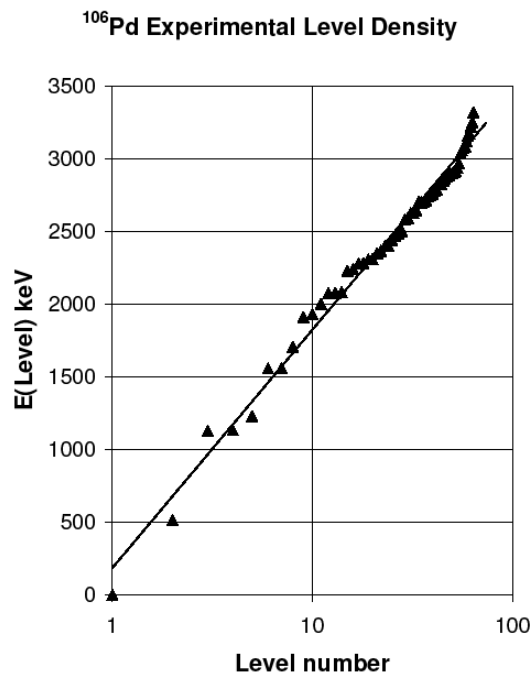


# The Influence of Nuclear Structure on Statistical Decay Properties

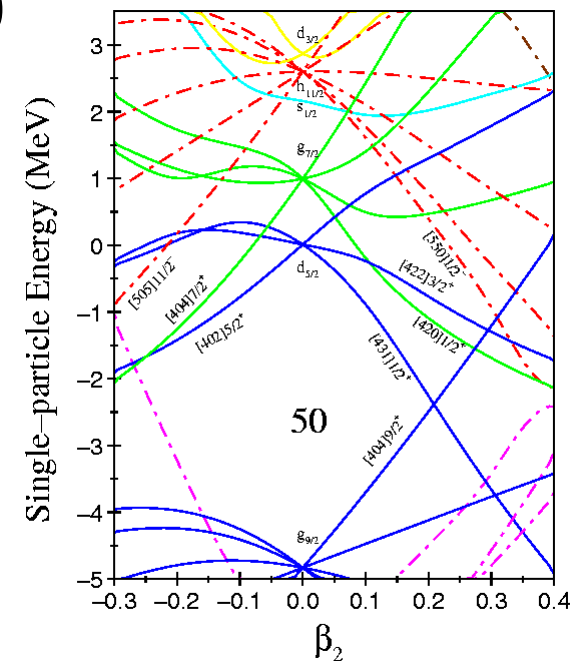
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2<sup>nd</sup> Workshop on Level Density and Gamma Strength



Oslo, May 11 - 15, 2009



# Introduction

New thermal neutron capture  $\gamma$ -ray cross section measurements ( $\sigma_\gamma$ ) have been performed at the Budapest Reactor.

- Evaluated Gamma-ray Activation File (EGAF)
- Statistical model calculations with DICEBOX and COSMO to determine  $\sigma_0$  and  $J^\pi$  values
- Search for nuclear structure (K) dependence.

Previous unpublished nuclear beta decay strength function measurements for  $^{117-124}\text{Cs}$  will be discussed.

- LBNL Total Absorption Spectrometer
- Nuclear structure dependence of the beta decay strength

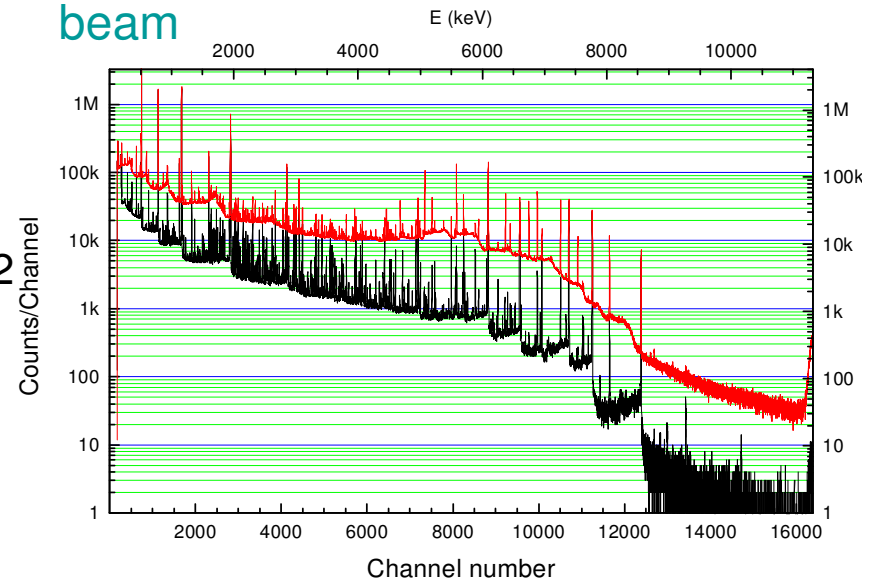


# EGAF (n, $\gamma$ ) Database



Thermal neutron  $\gamma$ -ray cross sections were measured for all elemental targets ( $Z=1,3-60,62-83,92$ ) at the Budapest Reactor\*.

Neutron beam



Thermal beam –  $2 \times 10^6 \text{ n} \cdot \text{s}^{-1} \text{cm}^{-2}$   
 Cold beam –  $5 \times 10^7 \text{ n} \cdot \text{s}^{-1} \text{cm}^{-2}$

\* Collaborators: G.L. Molnar†,  
 Zs. Revay, T. Belgya

† Deceased

Compton suppression lows background by a factor of  $\sim 5$  @ 1332 to  $\sim 40$  at 7 MeV.



# Internal Calibration of (n, $\gamma$ ) data

- Stoichiometric compounds containing elements with well-known cross sections

H, N, Cl, S, Na, Ti, Au

e.g. KCl, (CH<sub>2</sub>)<sub>n</sub>, Pb(NO<sub>3</sub>)<sub>2</sub>, Ti<sub>2</sub>SO<sub>4</sub>

- Homogenous mixtures

Aqueous (H<sub>2</sub>O) or acid (20% HCl) solutions, mixed powders (TiO<sub>2</sub>)

- Cross section of activation products

<sup>19</sup>F, <sup>28</sup>Al, <sup>100</sup>Tc, <sup>235</sup>U

Cross section  $\sigma_{\gamma}$  precision of <1% for strong transitions.  
First reliable database of  $\sigma_{\gamma}$  from thermal neutron capture



# EGAF Database

The Evaluated Gamma-ray Activation File (**EGAF**) was compiled as an IAEA Coordinated Research Project . **EGAF** contains >13,000  $\gamma$ -ray cross sections ( $\sigma_\gamma$ ) from 79 elements.

EGAF thermal (n, $\gamma$ ) Publications:

*Database of Prompt Gamma Rays from Slow Neutron Capture for Elemental Analysis*, R.B. Firestone, H.D. Choi, R.M. Lindstrom, G.L. Molnar, S.F. Mughabghab, R. Paviotti-Corcuera, Zs. Revay, V. Zerkin, and C.M. Zhou, IAEA STI/PUB/1263, 251 pp (2007); on-line at <http://www-pub.iaea.org/MTCD/publications/PubDetails.asp?pubId=7030>.

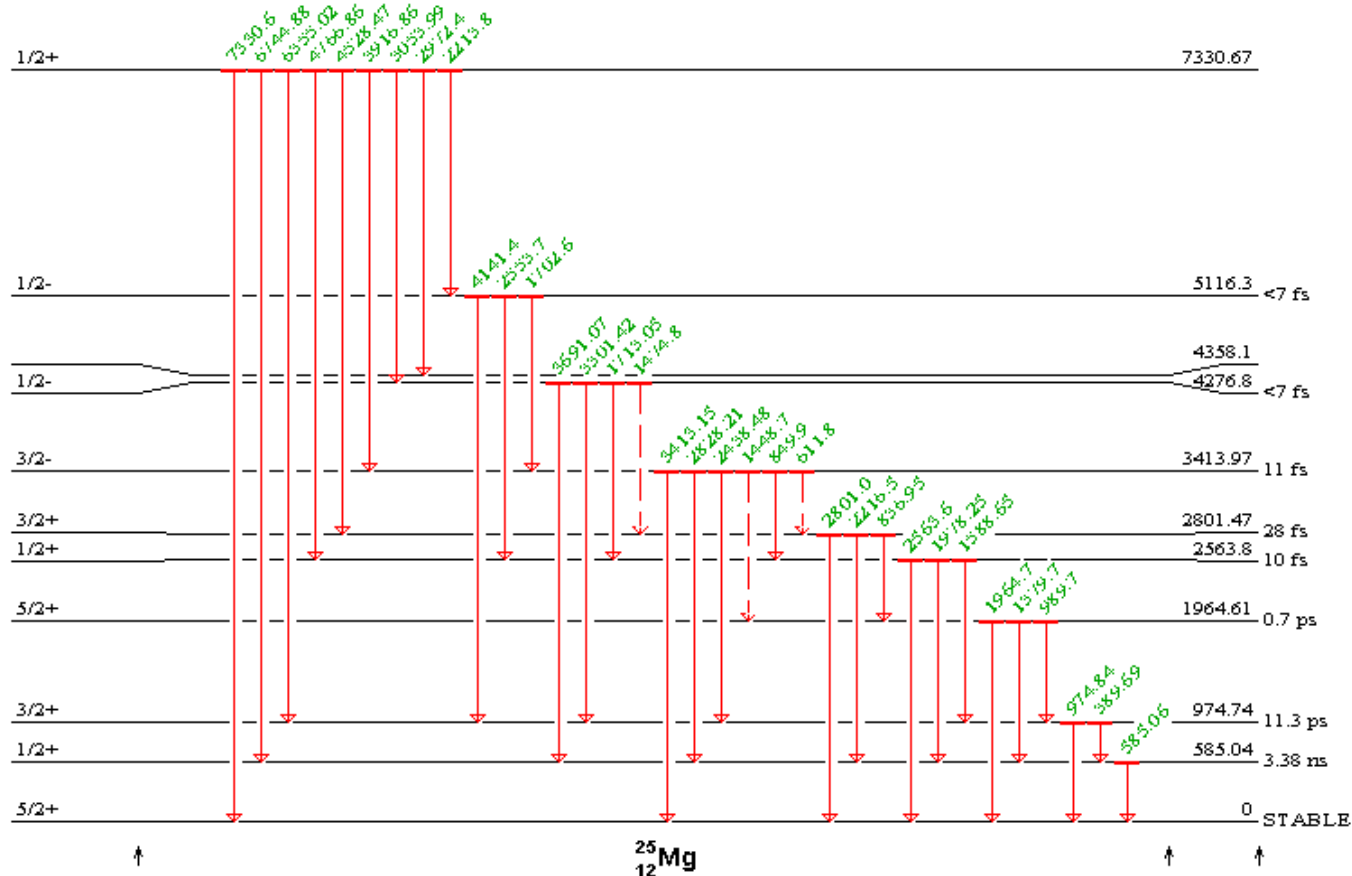
*Handbook of Prompt Gamma Activation Analysis with Neutron Beams*, Zs. Revay, T. Belgya, R.M. Lindstrom, Ch. Yonezawa, D.L. Anderson, Zs. Kasztovsky, and R.B. Firestone, edited by G.L. Molnar (Kluwer Publishers, 2004).



# Low-Z $\sigma_0$ Cross Sections

The decay schemes of low-Z isotopes are complete. Total thermal radiative neutron cross sections are calculated by

$$\sigma_0 = \sum \sigma_{\gamma+e}(\text{GS}) = \sum \sigma_{\gamma+e}(\text{CS})$$





# Cross section balance for $^{24}\text{Mg}(n,\gamma)^{25}\text{Mg}$

Cross section balance for the  $^{25}\text{Mg}$  neutron capture decay scheme

E(Level)	$\sigma(\text{in})$	$\sigma(\text{out})$	$\Delta\sigma$
0	0.0536(14)	0.0	0
585.01(3)	0.0406(11)	0.0398(14)	0.0008(18)
974.68(3)	0.0157(4)	0.0158(4)	0.0001(6)
1964.69(10)	0.00022(2)	0.00026(3)	0.00004(4)
2563.35(4)	0.00202(10)	0.00179(7)	0.00023(12)
2801.54(9)	0.00047(4)	0.00061(5)	0.00013(6)
3413.35(3)	0.0411(14)	0.0416(11)	0.0005(18)
4276.33(4)	0.0105(4)	0.0107(3)	0.0002(5)
4358.2(5)	0.00009(2)	0.0	0.00009(2)
5116.37(15)	0.00038(4)	0.00027(3)	0.00011(5)
7330.53(4)	0.0	0.0539(14)	0.0539(14)
$\sigma(\text{Mughabghab}[23])$		0.0536(15) b	
$\sigma(\text{Measured, average})$		0.0538(14) b	



# High-Z Cross Sections

For  $Z \geq 20$  measured neutron capture  $\gamma$ -ray decay schemes are generally incomplete due to unresolved continuum  $\gamma$ -rays.

$$\sigma_0 = 21.0 \pm 1.5 \text{ b (Mughabghab)}$$

$$\sigma_\gamma(\text{primary } \gamma\text{-rays}) = 0.55 \text{ b}$$

$$\sigma_\gamma(\text{secondary } \gamma\text{-rays}) = 20.26 \text{ b}$$

$^{105}\text{Pd}(n,\gamma)^{106}\text{Pd}$  cross section level feedings calculated from EGAF data.

$^{105}\text{Pd}(n,\gamma)^{106}\text{Pd}$  Level Feedings

E(level)	J <sup>π</sup>	Σσ <sub>γ</sub> (in)	Σσ <sub>γ</sub> (out)	ΔΣσ
0	0+	20.26		
511.844	2+	13.88	17.91	4.03
1128.04	2+	2.371	4.263	1.892
1133.79	0+	0.227	0.565	0.338
1229.2	4+	1.630	3.479	1.849
1557.67	3+	1.183	2.142	0.959
1562.16	2+	0.312	1.869	1.557
1706.44	0+	0.012	0.193	0.181
1909.39	2+	0.063	0.724	0.661
1932.37	4+	0.217	0.590	0.373
2001.56	0+	0.029	0.118	0.089
2077.1	6+	0.001	0.103	0.102
2077.37	(4)+	0.057	0.440	0.383
2084.39	-3	0.123	1.033	0.910
2242.4	2+	0.026	0.499	0.473
2278.47	0+	0	0.056	0.056
2282.89	4+	0.0007	0.275	0.274
2306.01	-3	0.053	0.542	0.489
2308.73	2+	0.000	0.283	0.283
2350.96	4+	0.018	0.304	0.286
2366.09	5+	0.003	0.116	0.114
2397.37	(5)-	0.055	0.263	0.209
2401	(2-,3-)	0.037	0.300	0.263
2439.11	2+	0.065	0.293	0.227
2472.09	0+	0.000	0.055	0.055
2484.76	(1-)	0.043	0.253	0.211
2500.01	-2	0.028	0.296	0.267
2578.64	(4-)	0.00004	0.221	0.221
...	...	...	...	...
...	...	...	...	...
9561.4	2+,3+		0.554	

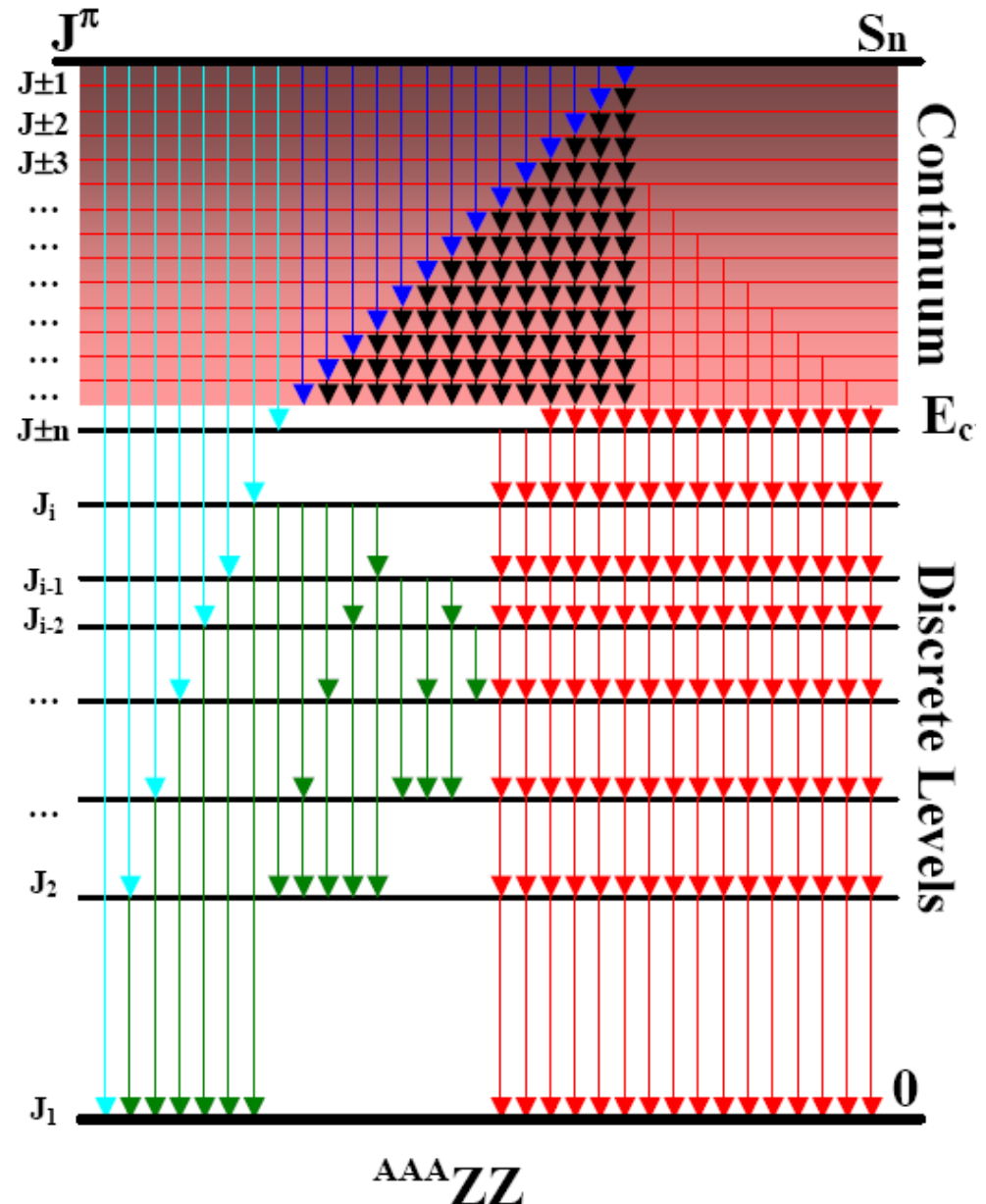




# Statistical Model Calculations

The  $(n,\gamma)$  continuum feeding is completely statistical and can be calculated if

1.  $\sigma_\gamma$  deexciting levels below a cutoff energy  $E_{\text{crit}}$  is complete.
2. Primary  $\sigma_\gamma$  populating the levels below  $E_{\text{crit}}$  from the capture state is complete.
3.  $J^\pi$  of levels below  $E_{\text{crit}}$  are well known.
4. Level density above  $E_{\text{crit}}$  is statistically distributed.
5. Photon strength deexciting levels above  $E_{\text{crit}}$  is statistically distributed.





# DICEBOX Calculations

DICEBOX is Monte Carlo code written by F. Becvar and M. Krlicka that generates complete simulated neutron capture decay schemes constrained by known nuclear properties and statistical models.

- A. Discrete primary and secondary  $\gamma$ -ray data from EGAF
- B.  $J^\pi$  data for  $E < E_{\text{crit}}$  from Reaction Input Parameter Library (RIPL)

C. Level density models

$$\rho(E, J) = \frac{f(J)}{T} \exp\left(\frac{E - E_0}{T}\right)$$

1. Constant temperature (CT)

2. Back-shifted Fermi (BSF) model  $\rho(E, J) = f(J) \frac{\exp\left(2\sqrt{a(E - E_1)}\right)}{12\sqrt{2}\sigma_c a^{1/4}(E - E_1)^{5/4}}$

D. E1 Photon Strength

$$f_{\text{BA}}^{(E1)}(E_\gamma) = \frac{1}{3(\pi\hbar c)^2} \frac{\sigma_G E_\gamma \Gamma_G^2}{(E_\gamma^2 - E_G^2)^2 + E_\gamma^2 \Gamma_G^2}$$

1. Brink-Axel (BA)

2. Kadmenski, Markushev, Furman (KMF) for spherical nuclei

3. Kopecky *et al* generalized Laurentian (GLO), temperature dep.

E. M1 Photon Strength

1. Single Particle (SP),  $f^{(E1)}/f^{(M1)}=5-7$  or  $f(M1)=1.2 \times 10^{-8} \text{ MeV}^{-3}$

2. Spin-Flip (SF), Laurentzian resonance  $\approx 8,5 \text{ MeV}$ , GSF  $\approx 4 \text{ MeV}$



# COSMO Calculations

**COSMO:** (**CO**n tinuous **S**tatistical **MO**del) – by R.B. Firestone (LBNL, Berkeley)

- Similar statistical model approach as DICEBOX
- The statistical decay scheme is binned above  $E_{\text{crit}}$
- Average level densities and photon strengths are assumed for each bin.
- COSMO calculations run faster than DICEBOX on small PCs
- No information on the statistical variation in the calculation is provided



# Palladium Calculations

**EGAF**  $\sigma_\gamma$  data are available for (n, $\gamma$ ) on all stable palladium targets

$^{102,104,105,106,108,110}\text{Pd}$

Calculations of the statistical  $\sigma_\gamma$  populating levels  $<E_{\text{crit}}$  were performed with **DICEBOX** where input model and parameters were selected by

- Comparison of calculated and experimental capture state width
- Comparison of E1 photon strength with photonuclear data from neighboring nuclei
- Dependence on  $E_{\text{crit}}$  of the result
- Comparison of  $\sigma_\gamma$  populating/depopulate levels  $<E_{\text{crit}}$



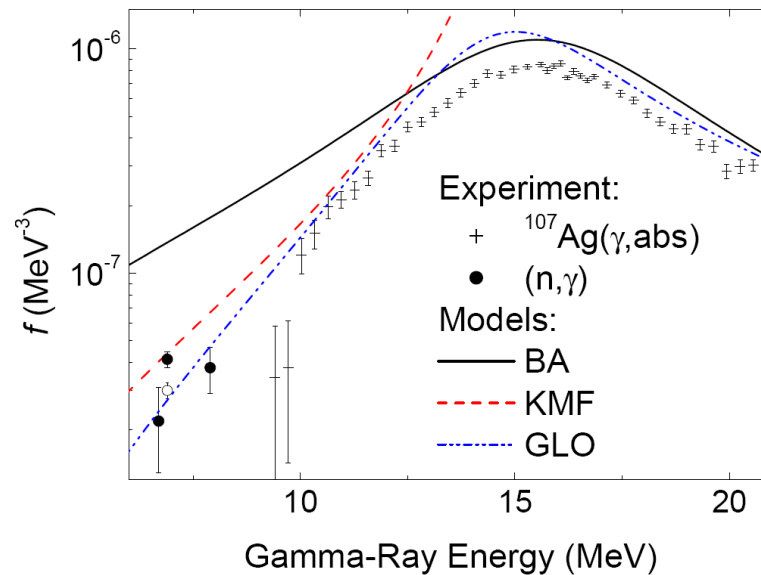
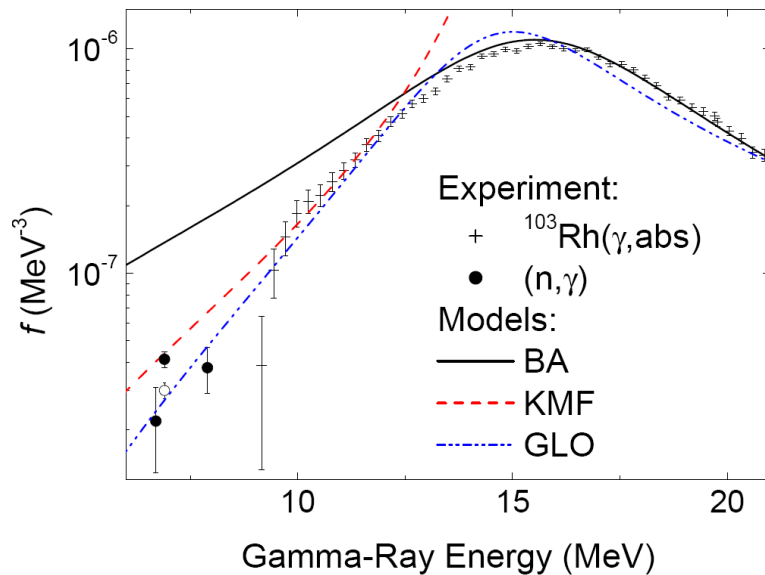
# Model dependence of the $^{106}\text{Pd}$ capture state width

## DICEBOX calculation of the capture state width

E1-PSF	M1-PSF	$\rho(E,J)$	$\Gamma_{\gamma}^{\text{tot}}$
Brink-Axel	Single Particle	Constant Temperature	410±47
Brink-Axel	Spin flip	Constant Temperature	352±42
Kadmenski et al (KMF)	Single Particle	Back-shifted Fermi	201±14
Kadmenski et al (KMF)	Spin flip	Back-shifted Fermi	172±12
Generalized Laurentzian	Single Particle	Back-shifted Fermi	156±8
Generalized Laurentzian	Spin flip	Back-shifted Fermi	126±8
Experiment (Mughabghab)			148±10



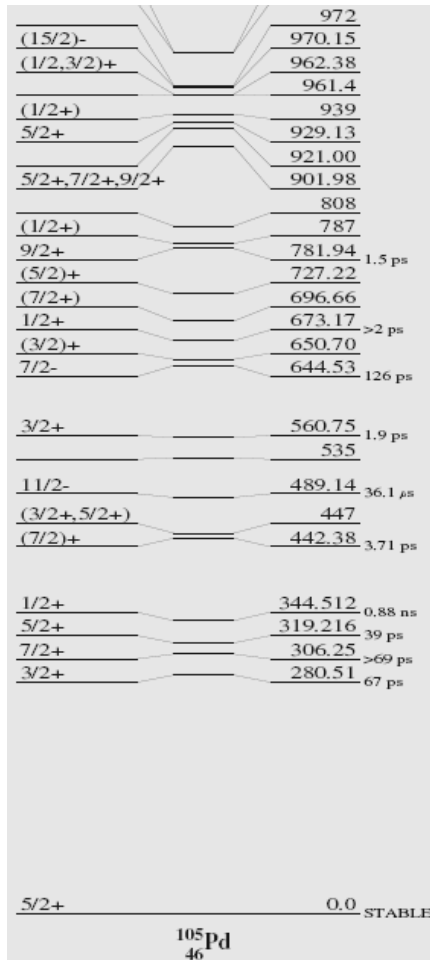
# E1 Photon Strength



Comparison of E1 photon strength models indicate that KMF and GLO models agree better with than Brink-Axel for  $^{103}\text{Rh}$  and  $^{107}\text{Ag}$  photonuclear data.

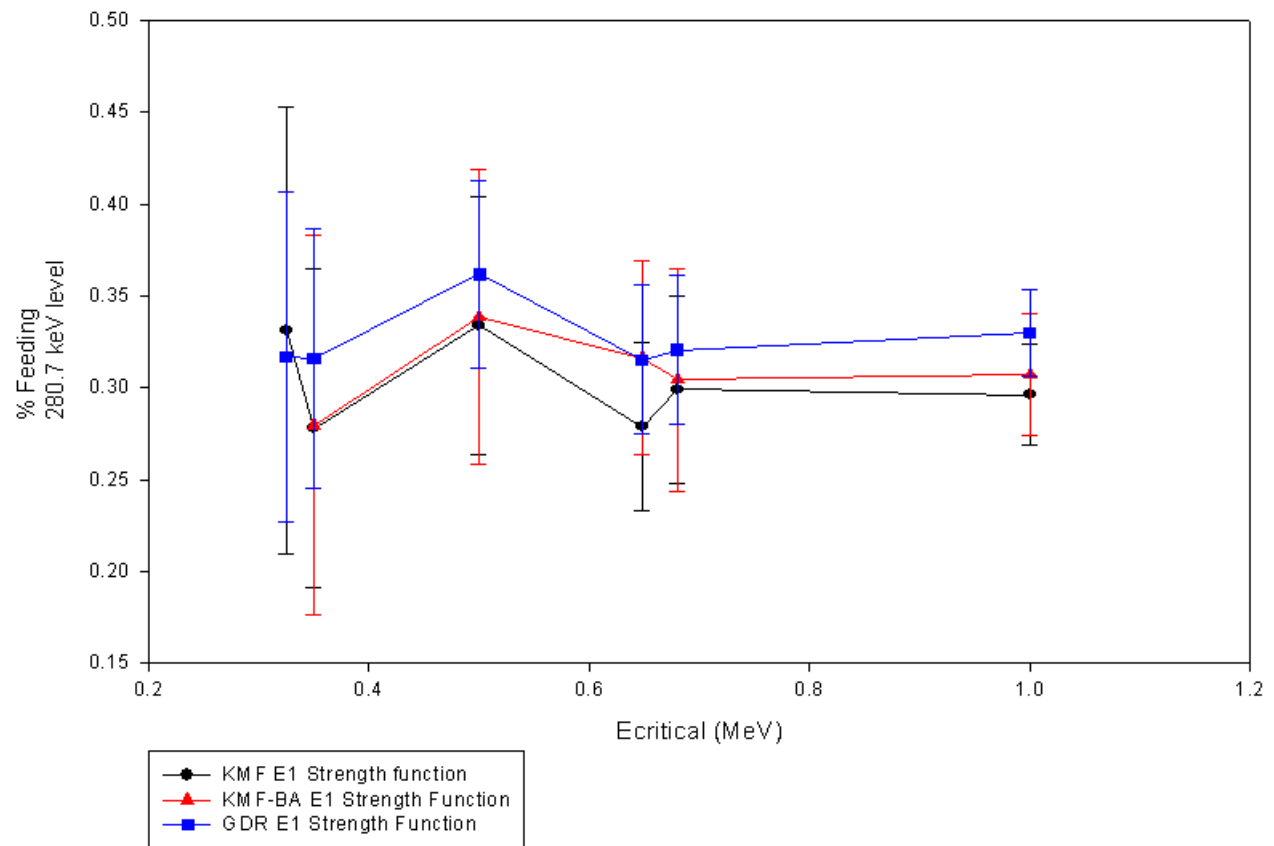


# Dependence of the Statistical Feeding on $E_{crit}$



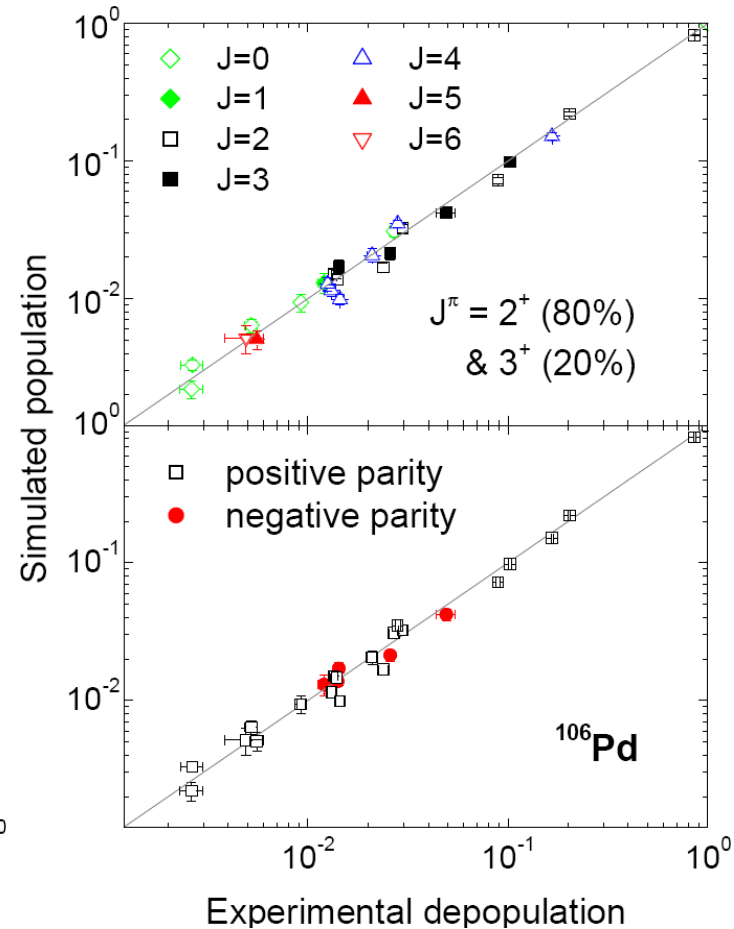
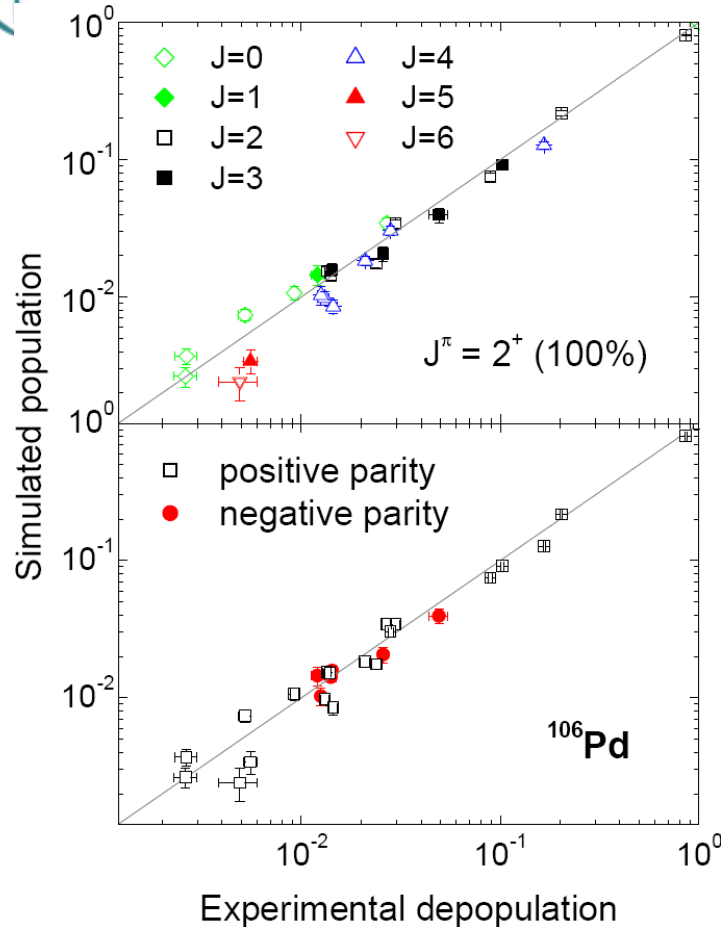
$^{104}\text{Pd}(n,g)$

% Cascade Feeding of 280.7 keV level vs  $E_{critical}$   
KMF, KMF-BA, and GDR E1 Strength functions





# Population/Depopulation Plot



Spin composition of the capture state is determined by a least-squares fit.

$$\sigma_0 = \sigma_\gamma(\text{GS})_{\text{expt}} + \sigma_\gamma(\text{GS})_{\text{Statistical}} = 20.3 \pm 0.3 \text{ b} + 1.4 \pm 0.3 \text{ b} = 21.7 \pm 0.5 \text{ b}$$

$$\sigma_0(\text{Mughabghab}) = 21.0 \pm 1.5 \text{ b}$$





# RIPL Structure Data Library

## RIPL input data for $^{106}\text{Pd}$

```

106Pd
number of levels: 133
number of gamma-rays: 212
number of levels in a complete level scheme: 30
number of levels with assigned spin and parity: 8
neutron separation energy: 9.561510 [MeV]
proton separation energy: 9.345901 [MeV]
  
```

NL	EL[MeV]	S/P	F	T1/2[s]	Ng	s	unc	Nf	Eg[MeV]	s-info	nd	m	p	mode
										Pg			Pe	Icc
1	0.000000	0.0	1	-1.00E+00	0	u				0+	0			
2	0.511851	2.0	1	1.21E-11	1	u				2+	0			
3	1.128010	2.0	1	3.12E-12	2	u	1	0.512	9.946E-01		1.000E+00		5.455E-03	
							2	0.616	6.461E-01		6.482E-01		3.252E-03	
							1	1.128	3.515E-01		3.518E-01		7.525E-04	
4	1.133770	0.0	1	6.80E-12	2	u				0+	0			
							2	0.622	9.968E-01		1.000E+00		3.171E-03	
							1	1.134	0.000E+00		0.000E+00		0.000E+00	
5	1.229250	4.0	1	1.34E-12	1	u				4+	0			
							2	0.717	9.978E-01		1.000E+00		2.183E-03	

$J^\pi$  are taken directly from ENSDF when uniquely assigned.

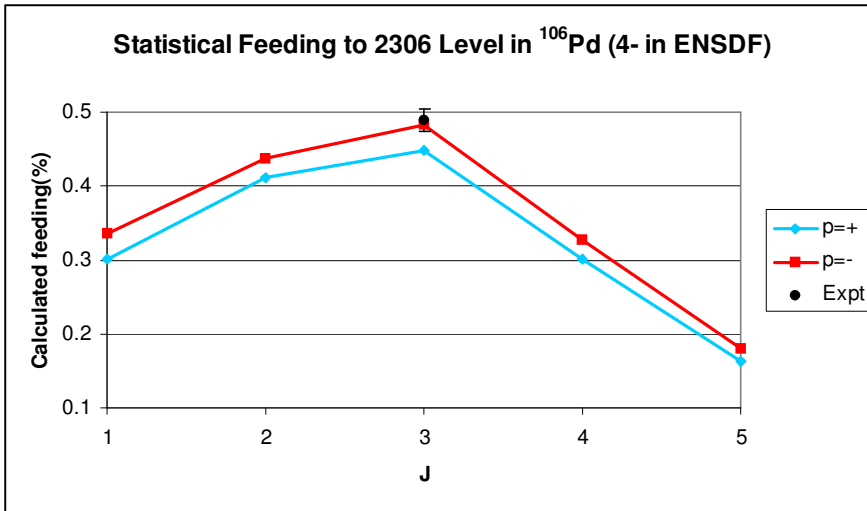
33	2.591200	-1.0	0		1				0.659	(2, 3)+	0		1.000E+00	0.000E+00
34	2.624400	0.0	1		3	u				0+	0			
							7	1.062	3.608E-01		3.611E-01		8.584E-01	
							3	1.496	2.505E-01		2.506E-01		4.201E-01	
							2	2.113	3.883E-01		3.883E-01		0.000E+00	
35	2.626870	-1.0	0		3					(2, 3)+	0			
							7	1.065	7.453E-02		7.453E-02		0.000E+00	
							3	1.499	6.211E-01		6.211E-01		0.000E+00	
							2	2.115	3.043E-01		3.043E-01		0.000E+00	

$J^\pi$  uncertain

$E_{crit}$  is limited to the highest level with a unique, known  $J^\pi$

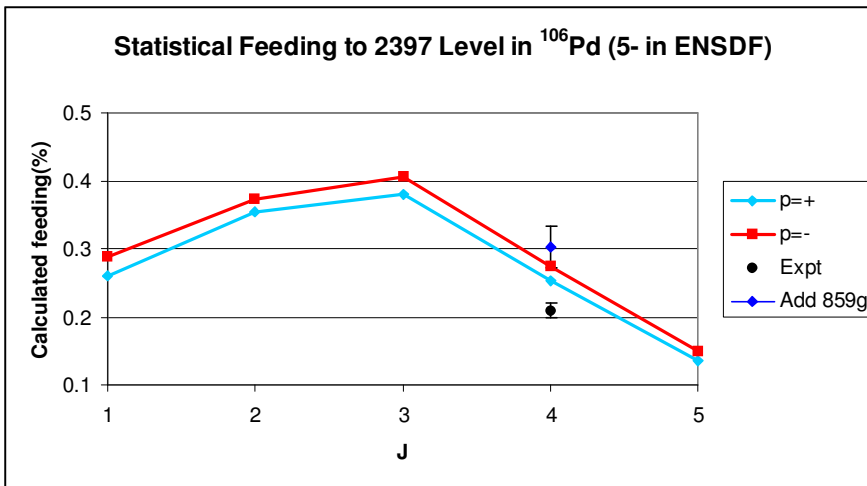


# R IPL Input Data Problems: $J^\pi$ Errors



Statistical model calculations can be used to constrain  $J^\pi$  values.

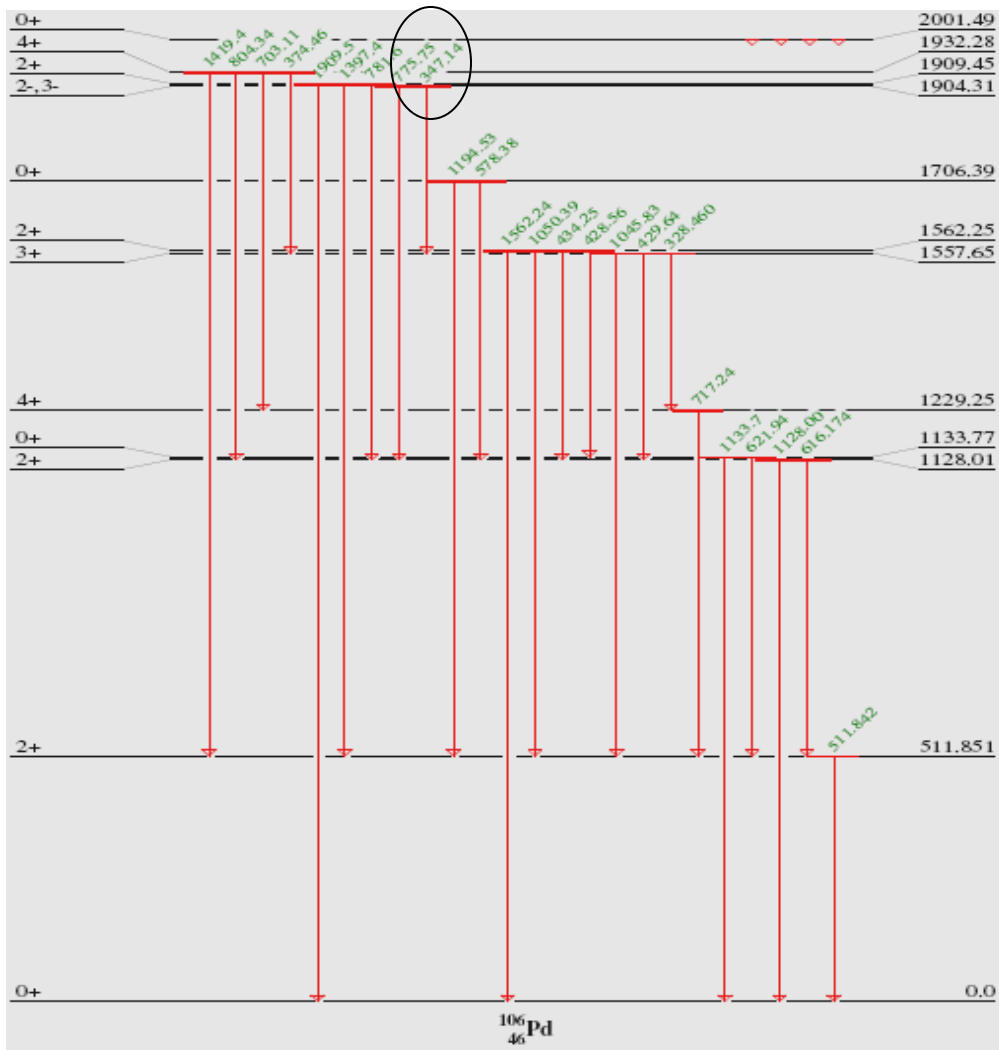
2306-keV level feeding is consistent with  $J^\pi=3^-$  not  $J^\pi=4^-$  adopted in ENSDF on basis of  $\gamma(\theta)$ .



2397-keV level feeding is consistent with  $J^\pi=4^-$  when additional  $\gamma$ -ray is placed. Assignment of  $J^\pi=5^-$  in ENSDF is based on  $L=(5)$  in  $(p,t)$ .



# Level Scheme Errors



**Mistaken level assignment**

1904.3 keV level assigned by the Ritz principal.

$$\sigma_{\gamma}(1904)_{\text{expt}} = 0.12 \text{ b}$$

$$\sigma_{\gamma}(1904)_{\text{DICEBOX}} = 1.13 \text{ b}$$

Reassigning placement of the 347- and 776-keV  $\gamma$ -rays to deexcite the 1909.4-keV level gives

$$\sigma_{\gamma}(1909)_{\text{expt}} = 0.62 \text{ b}$$

$$\sigma_{\gamma}(1909)_{\text{DICEBOX}} = 0.83 \text{ b}$$

**Statistical model calculations can be used to improve the nuclear structure information in RIPL.**



# $^{106}\text{Pd}$ $J^\pi$ Values

New  $J^\pi$  values or level placements were determined for 7 of 64 levels.

The cross section deexciting the 2472 level is inconsistent with  $J > 0$ . Decay to  $0^+$  levels rules out  $J^\pi = 0^\pm$ . Either there is missing  $\gamma$ -ray deexcitation or this level placement is incorrect.

E	dE	JPI(ENSDF)	JPI(This work)	E	dE	JPI(ENSDF)	JPI(This work)
0	0	0+	0+	2626.76	0.1	(2,3)+	(2,3)+
511.844	0.02	2+	2+	2646.3	0.2	(4+)	(4+)
1128.04	0.03	2+	2+	2699.36	0.16	(6)-	(6)-
1133.79	0.06	0+	0+	2705.3	0.08	(1)+	(1)+
1229.2	0.03	4+	4+	2713.78	0.09	2+,3+	2+,3+
1557.67	0.04	3+	3+	2741	0.5	4+	1,2+
1562.16	0.04	2+	2+	2747.72	0.16	2-,3-	2-,3-
1706.44	0.08	0+	0+	2757	0.04	5+	5+
1904.21	0.09	2-,3-	NO LEVEL	2775.88	0.11	(4+)	(4+)
1909.39	0.07	2+	2+	2783.71	0.13	2+	2+
1932.37	0.05	4+	4+	2820.94	0.11	2+	2+
2001.56	0.11	0+	0+	2828	1.7	0+	0+
2076.69	0.04	4+	4+	2861.41	0.17	(+)	(+)
2077.01	0.06	6+	6+	2850.79	0.15	2+,3+	2+,3+
2084.39	0.07	3-	3-	2878.27	0.19	0+	0+
2229.2	0.21		NO LEVEL	2886.16	0.23	(-)	(-)
2242.4	0.06	2+	2+	2897.42	0.16	(1-,4-)	4-
2278.47	0.14	0+	0+	2902.31	0.2	2+	2+
2282.89	0.09	4+	4+	2908.53	0.17	(1-)	(1-)
2306.01	0.06	4-	3-	2918.56	0.13	2+	2+
2308.73	0.11	2+	2+	2935.8	0.2	(2-,3-)	(2-,3-)
2350.96	0.09	4+	4+	2968.5	0.25	3-	3-
2366.09	0.1	5+	5+	3037.45	0.11	1,2	1,2
2397.37	0.08	(5)-	4-	3056.38	0.12	1+	1+
2400.84	0.25	2-,3-	2-,3-	3071.03	0.23	(2,3)-	(2,3)-
2439.11	0.1	2+	2+	3083.52	0.13	0	0+
2472.09	0.17	1+,2+	1+,2+???	3118.45	0.18	(6+)	(6+)
2484.76	0.25	(1-)	(1-)	3161.1	0.5	2+	2+
2500.01	0.12	2-	2-	3173.8	0.7	(2+,3+)	(2+,3+)
2578.64	0.1	(5-)	(4-)	3221.64	0.15	0+	0+
2591.2	0.4	(2,3)+	(2,3)+	3252	0.4	2+	2+
2624.21	0.13	0+	0+	3319.52	0.25	0+	0+



# Pd total radiative cross section $\sigma_0$ results

Reaction E(n)=thermal	# levels below $E_{\text{crit}}$	$\sigma_0$ (literature) (barns)	$\sigma_0$ (this work) (barns)
$^{102}\text{Pd}(n,\gamma)^{103}\text{Pd}$	2	$1.6\pm 0.2$	$1.1\pm 0.4$
$^{104}\text{Pd}(n,\gamma)^{105}\text{Pd}$	5	$0.65\pm 0.30$	$0.77\pm 0.17$
$^{105}\text{Pd}(n,\gamma)^{106}\text{Pd}$	28	$21.0\pm 1.5$	$21.7\pm 0.5$
$^{106}\text{Pd}(n,\gamma)^{107}\text{Pd}$	5	$0.30\pm 0.03$	$0.36\pm 0.10$
$^{108}\text{Pd}(n,\gamma)^{109}\text{Pd}$	11	$7.6\pm 0.5$	$7.2\pm 0.5$
$^{108}\text{Pd}(n,\gamma)^{109}\text{Pd}^m$	2	$0.185\pm 0.010$	$0.185\pm 0.011$
$^{110}\text{Pd}(n,\gamma)^{111}\text{Pd}$	5	$0.70\pm 0.17$	<b><math>0.34\pm 0.10</math></b>

Total radiative thermal neutron cross sections can be determined even when only one  $\gamma$ -ray is observed.



# Treatment of Spin and Parity

The spin distribution of the level density  $f(J)$  has been defined as (Bethe, 1937)

$$f(J) = \frac{2J+1}{2\sigma_c^2} \exp\left(-\frac{(J+1/2)^2}{2\sigma_c^2}\right)$$

Where  $\sigma_c^2$  is the spin cutoff parameter. This should be multiplied by a parity distribution parameter  $f(\pi)$ .

One parameterization of  $f(\pi)$  is suggested by Al-Quraishi [1] and used in COSMO calculations is

$$f(\pi=+) = 0.5 \times [1 + (1 + \exp(C_\pi(E - \Delta_\pi))^{-1})]$$

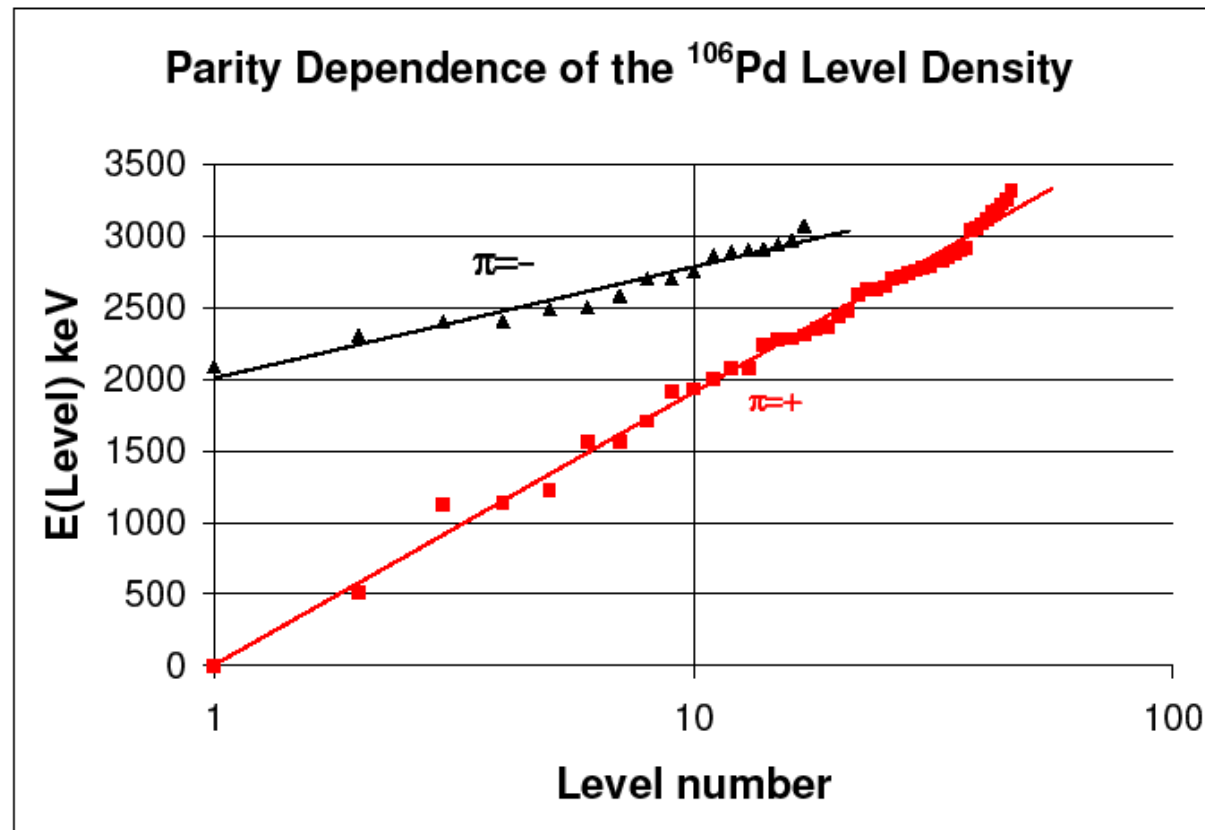
$$f(\pi=-) = 1 - F(\pi=+)$$

where  $C_\pi$ ,  $\Delta_\pi$  are given by the authors.

[1] S.I. Al-Quraishi, S.M. Grimes, T.N. Massey, and D.A. Resler, Phys. Rev. **C67**, 015803.



# Parity Dependence of the Experimental Level Density

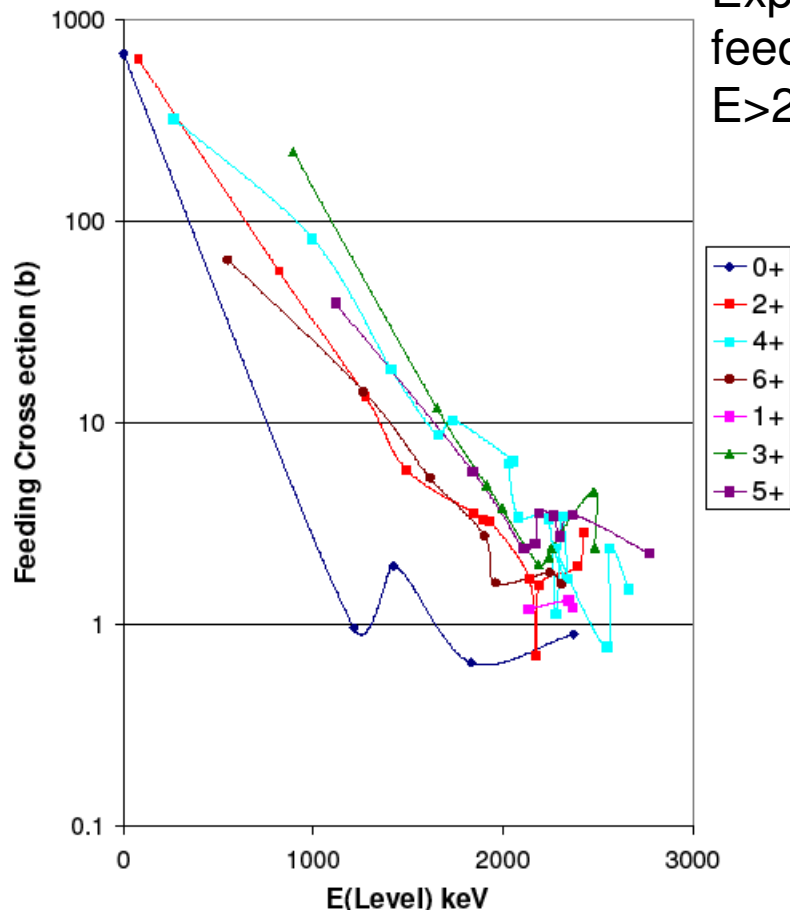


Positive and negative parities may have different temperatures



# Experimental and Statistical $J^{\pi=+}$ Feedings

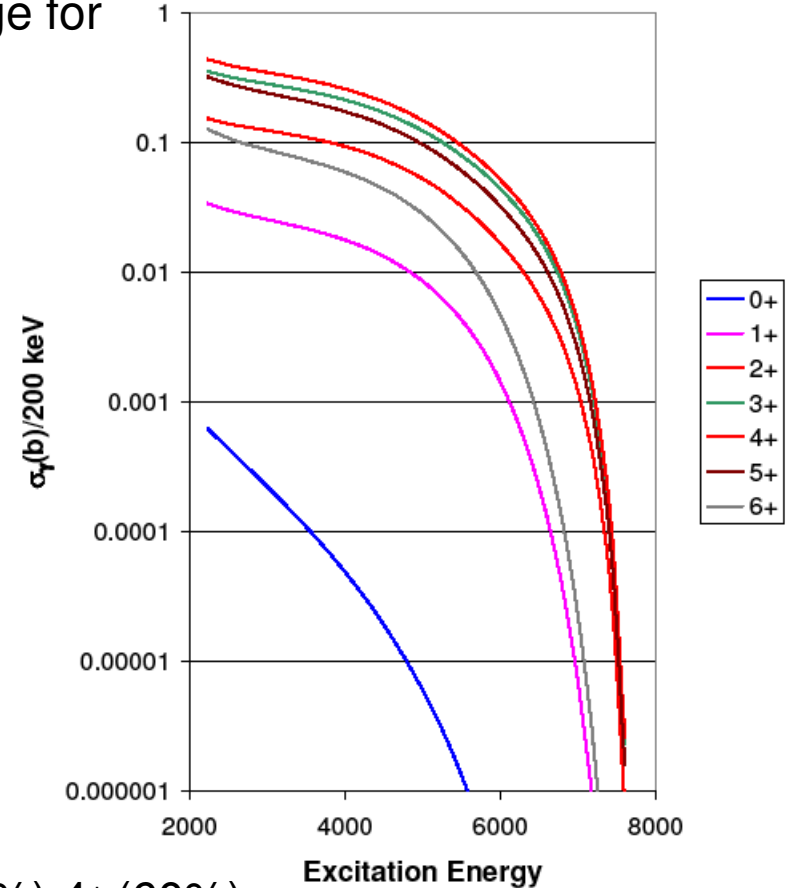
$^{167}\text{Er}(n,\gamma)^{168}\text{Er}$



Experimental level feedings converge for  $E > 2$  MeV.

Capture state is  $3^+$  (40%),  $4^+$  (60%)

$^{168}\text{Er}$  Statistical  $\pi=+$  Feedings



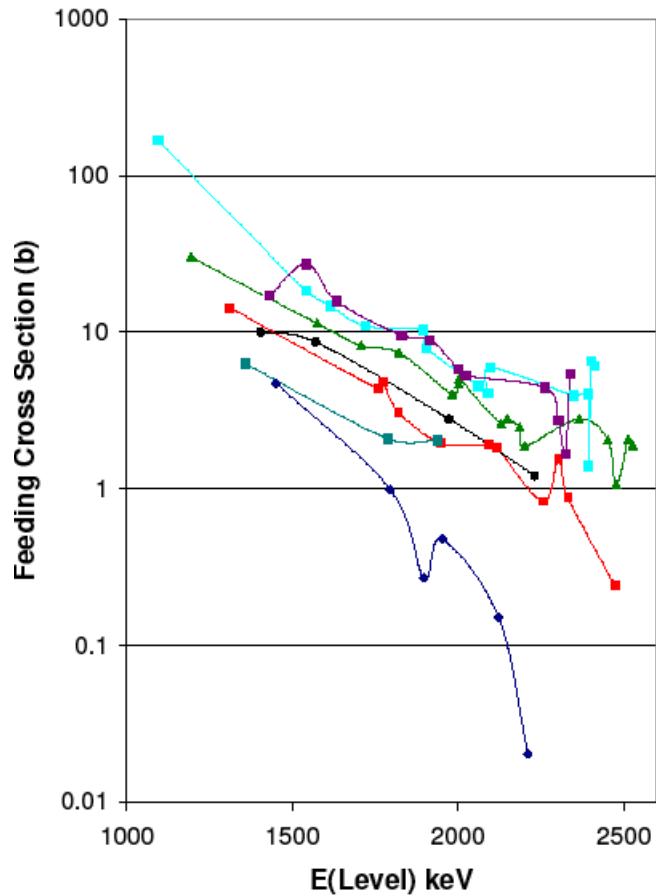




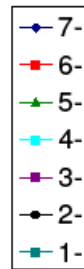
# Experimental and Statistical $J^{\pi=-}$ Feedings

$^{167}\text{Er}(n,\gamma)^{168}\text{Er}$

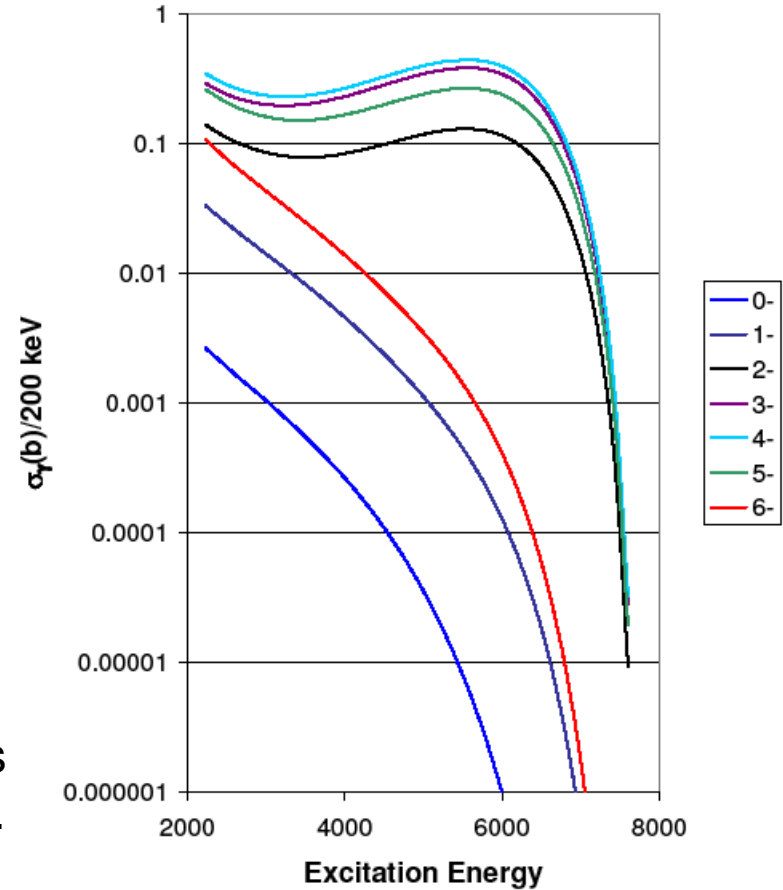
$^{168}\text{Er}$  Statistical  $\pi=-$  Feedings



Note that no  $J^{\pi}=0^-$  levels are observed.



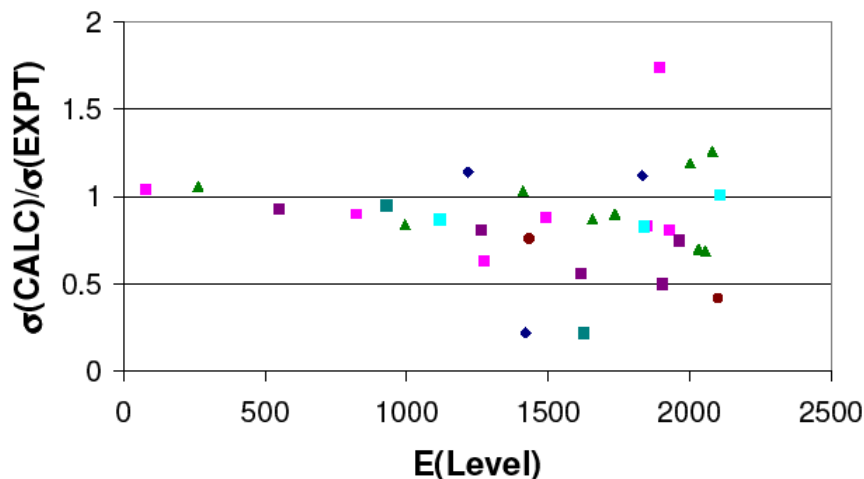
Input  $J^{\pi}$  values are from RIPL.



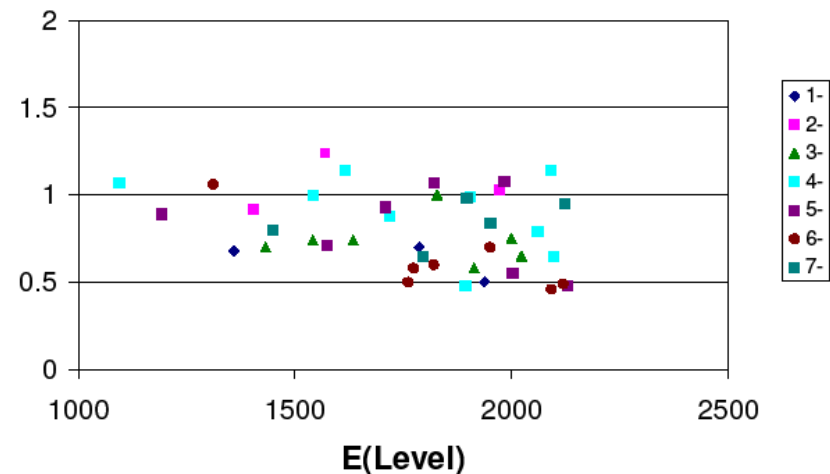


# Comparison of $^{168}\text{Er}$ Experimental and Statistical Cross Sections

$^{168}\text{Er}$ : Comparison of Experimental and Calculated Level Feedings  $\pi=+$



$^{168}\text{Er}$ : Comparison of Experimental and Calculated Level Feedings  $\pi=-$



Level densities - CT model, E1 dipole strength Brink-Axel.  $E_{\text{crit}}=2135$  keV (80 levels)

For  $E < 1300$  keV,  $\sigma(\text{CALC})/\sigma(\text{EXPT})=0.96$ .

For  $E > 1300$  keV,  $\sigma(\text{CALC})/\sigma(\text{EXPT})= 0.80$

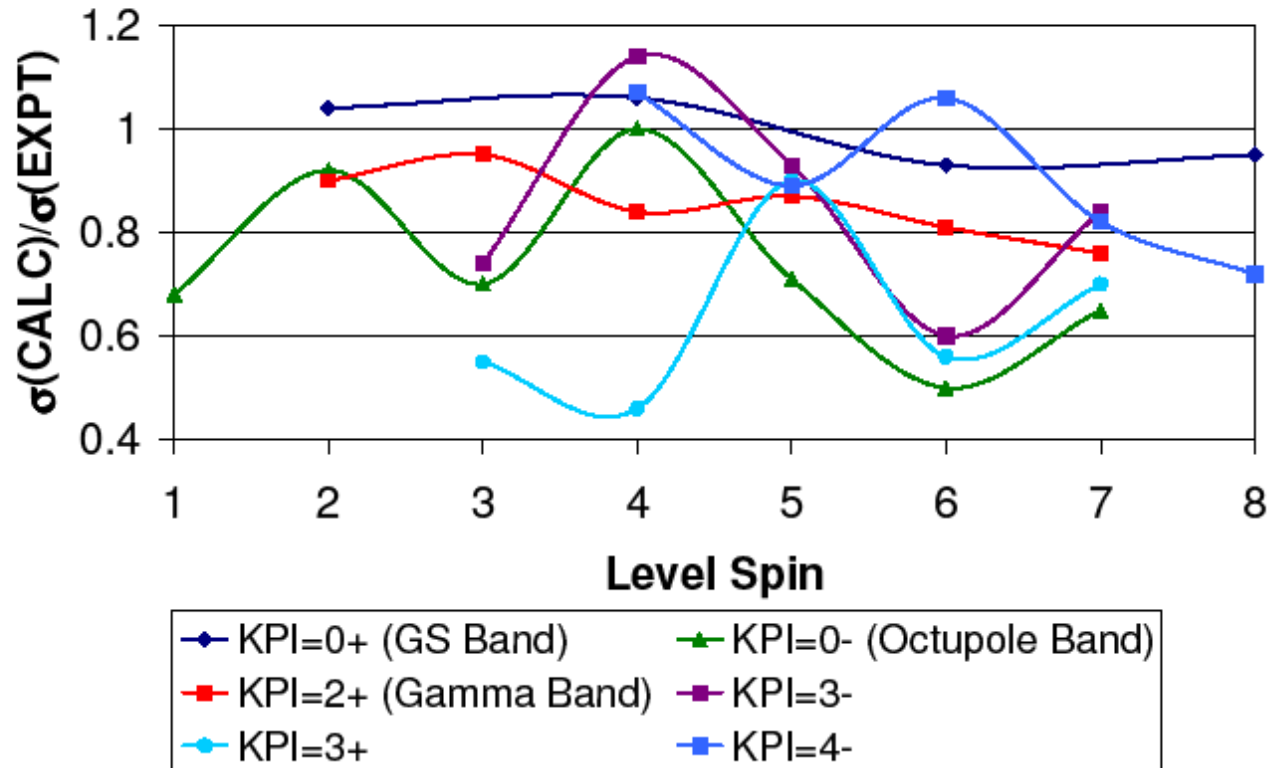
- Incomplete experimental level scheme
- Nuclear structure contributions

Outliers may be due to low statistics, missing transitions, and incorrect  $J^\pi$  assignments.



# K-Dependence

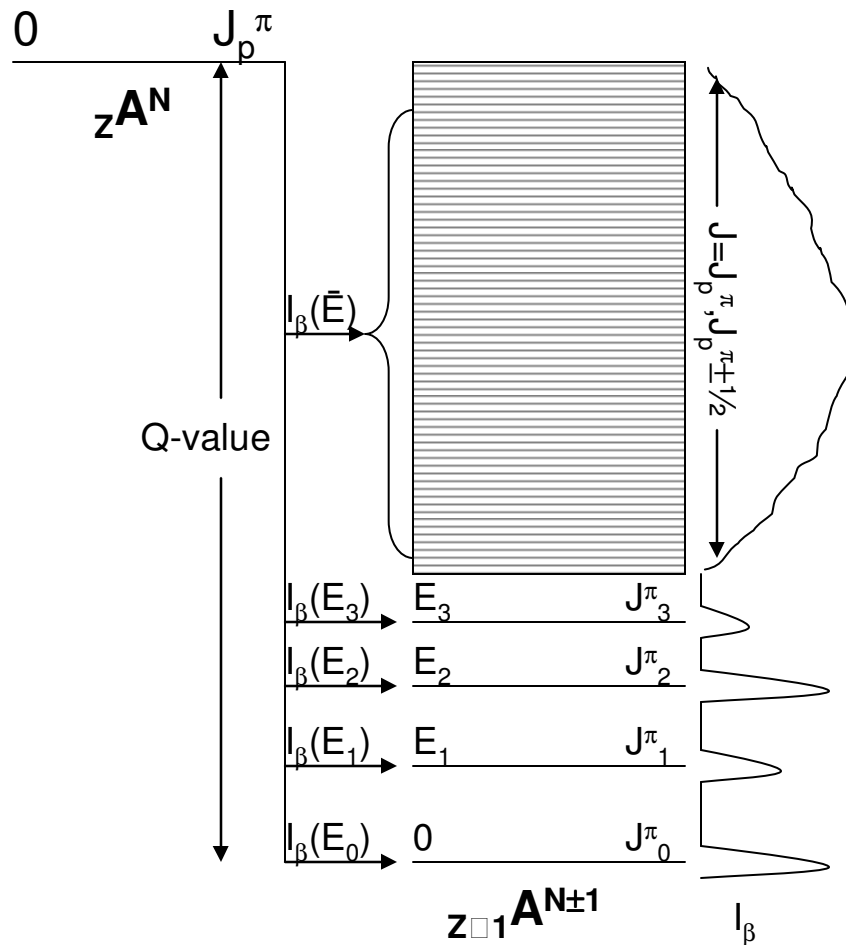
$^{168}\text{Er}$  K-Dependence of the Statistical Model



Significant dependence on K is observed. Note that oscillations in the fit appear to be correlated, possibly due to incorrect spin distribution.



# $\beta$ -Strength Function



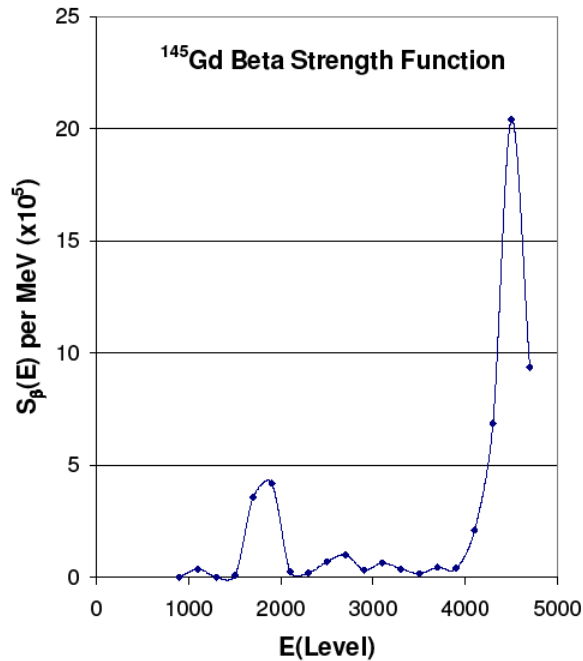
$$S_\beta(E) = \frac{I_\beta(E)}{f(Q_\beta - E)t_{1/2}} \text{ (individual levels)}$$

$$S_\beta(E) = \frac{\rho(E)}{f(Q_\beta - E)t_{Ave}} \text{ (continuum levels)}$$

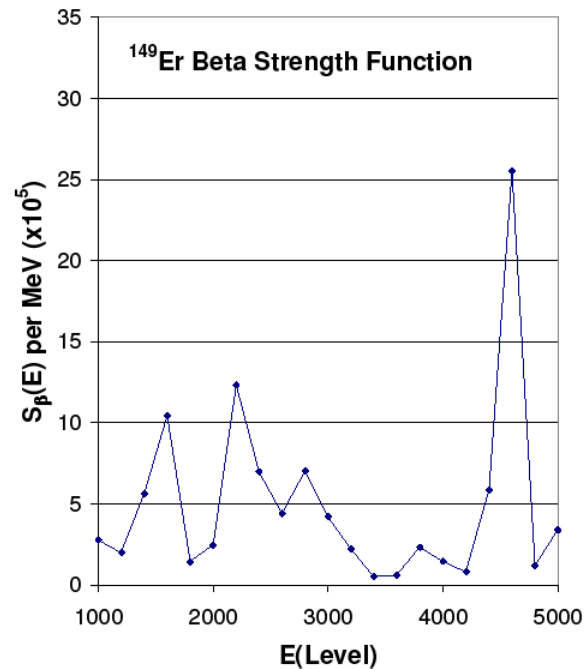
In the continuum  $S_\beta(E)$  is calculated from the product of the level density  $\rho(E)$  and an average beta strength  $ft_{Ave}$  analogous to the statistical model for thermal neutron capture  $\gamma$ -ray decay analysis.



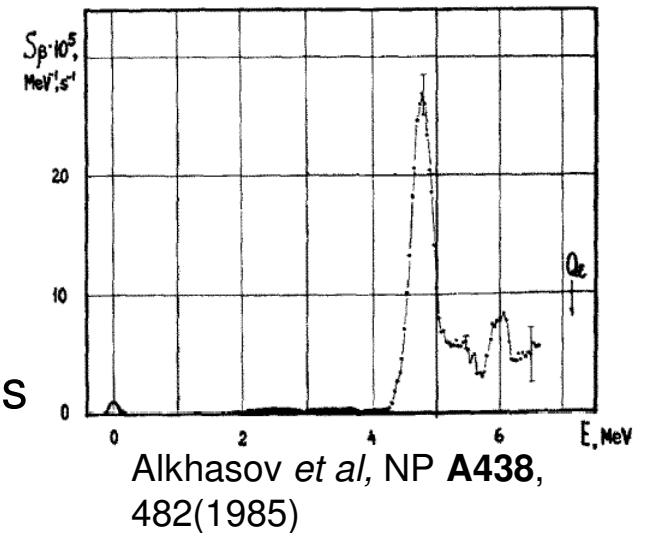
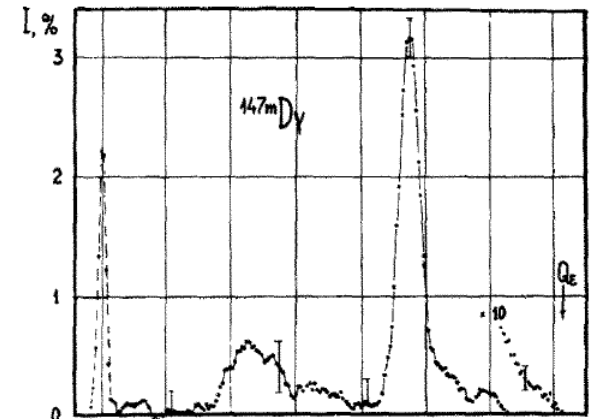
# Nuclear Structure in the $\beta$ -Strength Function



R.B. Firestone *et al*,  
PRC**25**, 527(1982)



R.B. Firestone *et al*,  
PRC**39**, 219(1988)



It was shown in detailed decay scheme studies and Total Absorption Spectrometer (TAS) measurements that the beta strength to the continuum in  $^{145}\text{Gd}$ ,  $^{147}\text{Dy}$ , and  $^{149}\text{Er}$  is not statistical.



# Nuclear Structure and the Decay of N=81 Isotopes

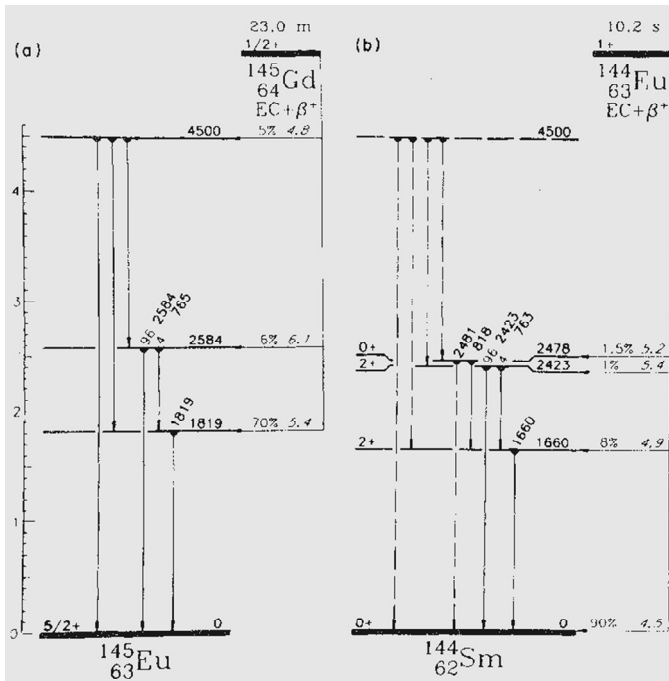
The  $\beta$ -strength function is dominated by the fast  $\pi h_{11/2} \rightarrow \nu h_{9/2}$  spin-flip transition from virtual  $\pi h_{11/2}$  pairs.

- $\pi s_{1/2}$  (XX)
- $\pi d_{3/2}$  (XX)
- $\pi h_{11/2}$  (XX)
- $\pi d_{5/2}$  XX XX XX
- $\pi g_{7/2}$  XX XX XX XX

Shell + weak coupling model calculation.

$(\pi h_{11/2})^1 (\nu h_{9/2})^1 (vs_{1/2} + vd_{3/2})^{-1}$	5300
$(\pi h_{11/2})^1 (\nu i_{13/2})^1 (vs_{1/2} + vd_{3/2})^{-1}$	4900
$(\pi h_{11/2})^1 (\nu f_{7/2})^1 (vs_{1/2} + vd_{3/2})^{-1}$	3900

$^{145}\text{Gd}$  proton valence states



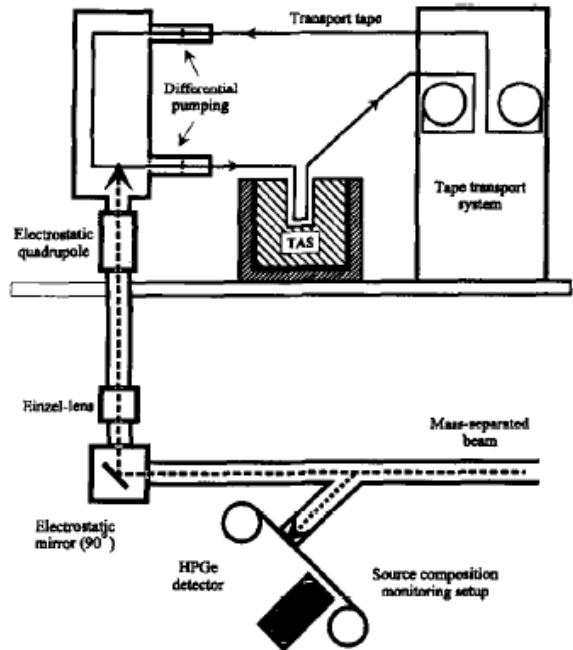
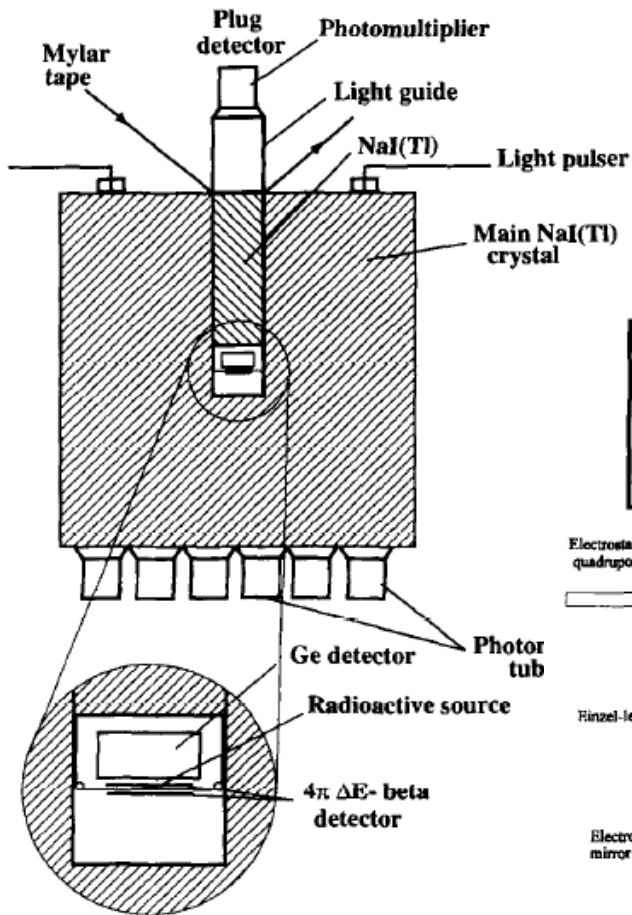
Lower level feeding dominated by core decay to  $\pi d_{5/2} \otimes 2^+$  configurations.

$(\pi d_{3/2})^1$	1042
$(\pi s_{1/2})^1$	808
$(\pi h_{11/2})^1$	716
$(\pi g_{7/2})^{-1} (\pi d_{5/2})^2$	330
$(\pi d_{5/2})^1$	0



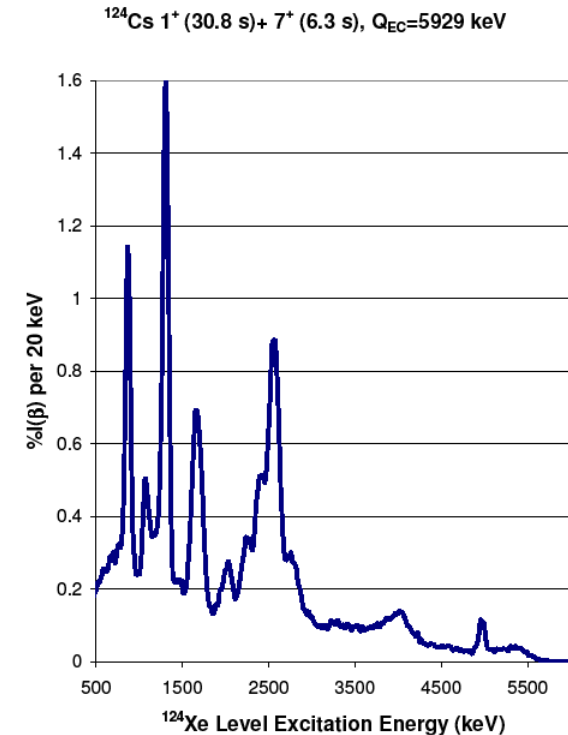
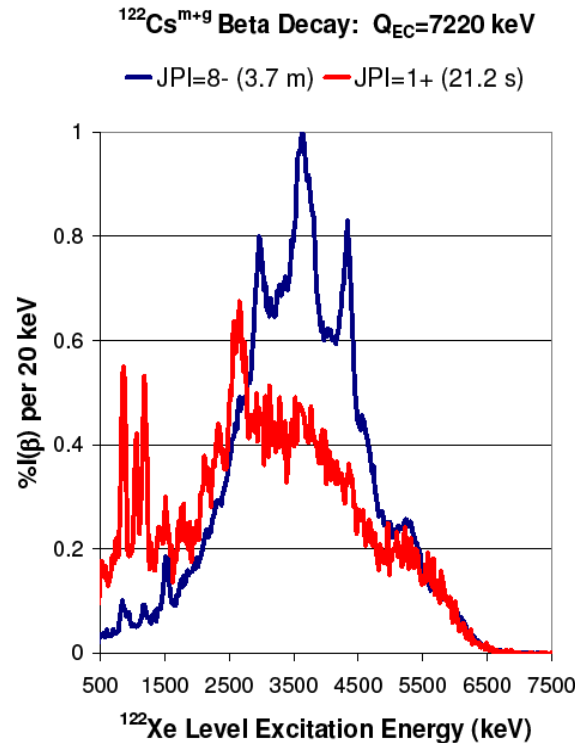
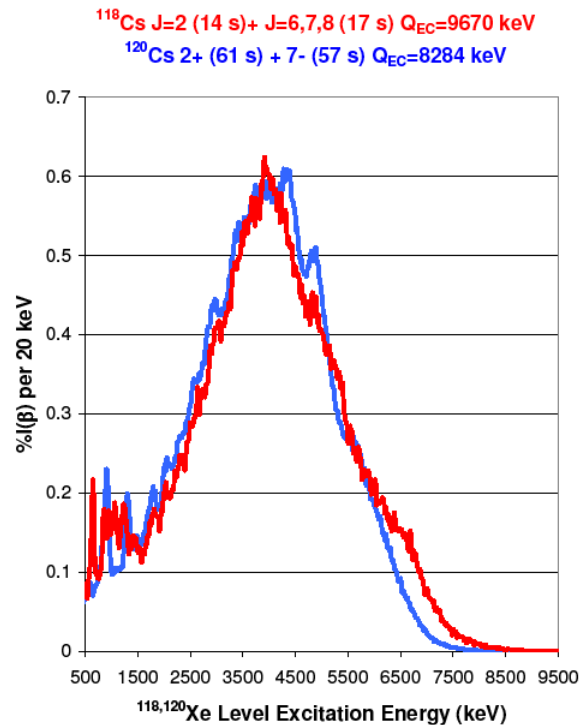
# LBL Total Absorption Spectrometer (TAS)

Mass separated sources produced with the LBNL HILAC (1995) and transferred by tape to inside a 36×36 cm NaI(Tl) detector. TAS is ~95% efficient for 10 MeV  $\gamma$ -rays.





# 118-124Cs Odd-Odd Isotopes



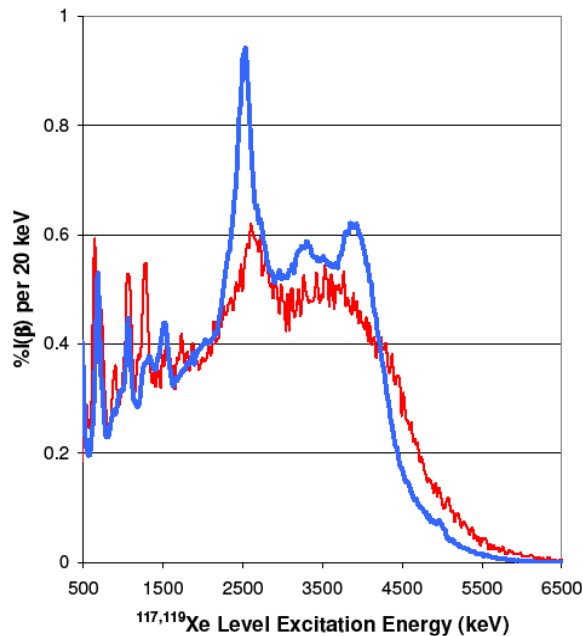
- $^{118}\text{Cs}$  and  $^{120}\text{Cs}$  decays have nearly the same  $\beta$ -strengths with little structure
- Onset of high excitation nuclear structure appears at  $^{122}\text{Cs}$  with N=65
- Isomeric and GS decays can have very different nuclear structure features



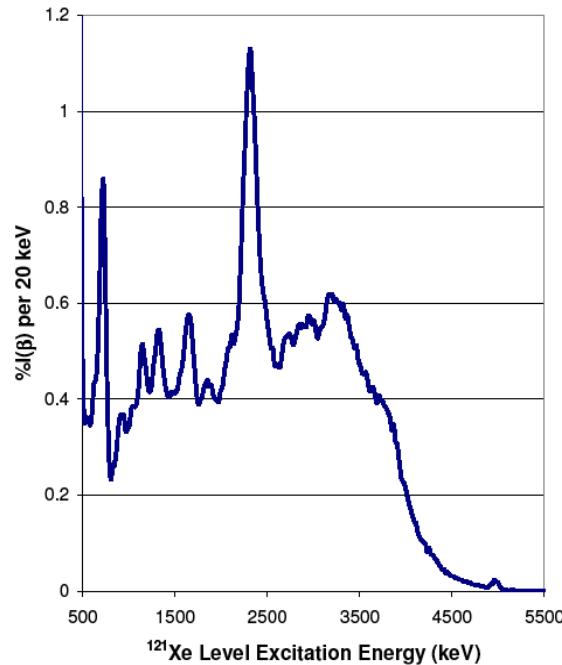


# 117-123Cs Odd-Even Isotopes

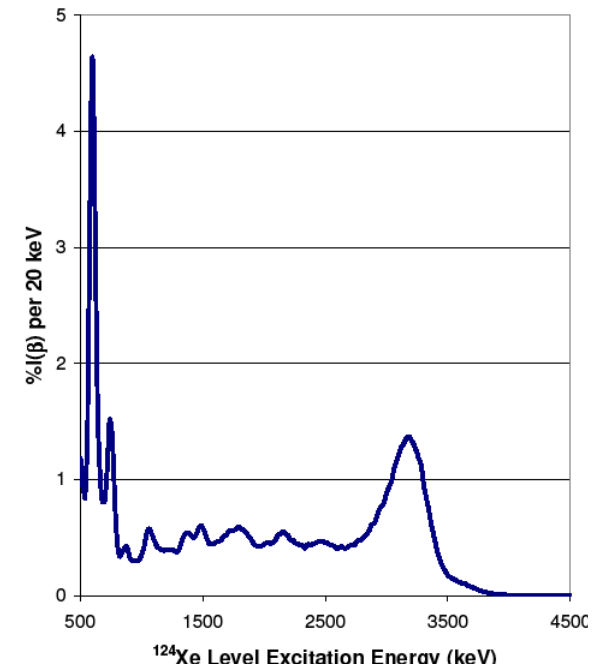
$^{117}\text{Cs}$  3/2+ (6.5 s)+ 9/2+ (8.4 s)  $Q_{\text{EC}}=7740$  keV  
 $^{119}\text{Cs}$  3/2+ (30.4 s)+ 9/2+ (43 s)  $Q_{\text{EC}}=6489$  keV



$^{121}\text{Cs}$  3/2+ (155 s)+ 9/2+ (122 s)  $Q_{\text{EC}}=5372$  keV



$^{123}\text{Cs}$  1/2+ (5.87 m)+ 11/2- (1.64 s),  $Q_{\text{EC}}=4265$  keV



- Little nuclear structure >3 MeV is seen in  $^{117}\text{Cs}$  Decay
- Comparable  $\beta$ -strengths with resonances >3 MeV are seen in  $^{119-123}\text{Cs}$  decays
- Transition to high excitation resonance structure appears to occur at N=64.



# Nuclear Structure in the $\beta$ -strength of the Cs isotopes

Strong  $\beta$ -strength is expected for  $\pi g_{9/2} \rightarrow \nu g_{7/2}$  spin-flip transition.

The Cs isotopes have  $\beta_2=0.2-0.3$  bringing the  $\pi g_{9/2+}[404]$  configuration near the GS. A fast spin-flip  $\beta$ -decay transition is expected to the  $\nu g_{7/2+}[404]$  configuration at  $\approx 4$  MeV.

Disappearance of the high excitation structure in the  $\beta$ -strength for  $N < 64$  has not been explained.

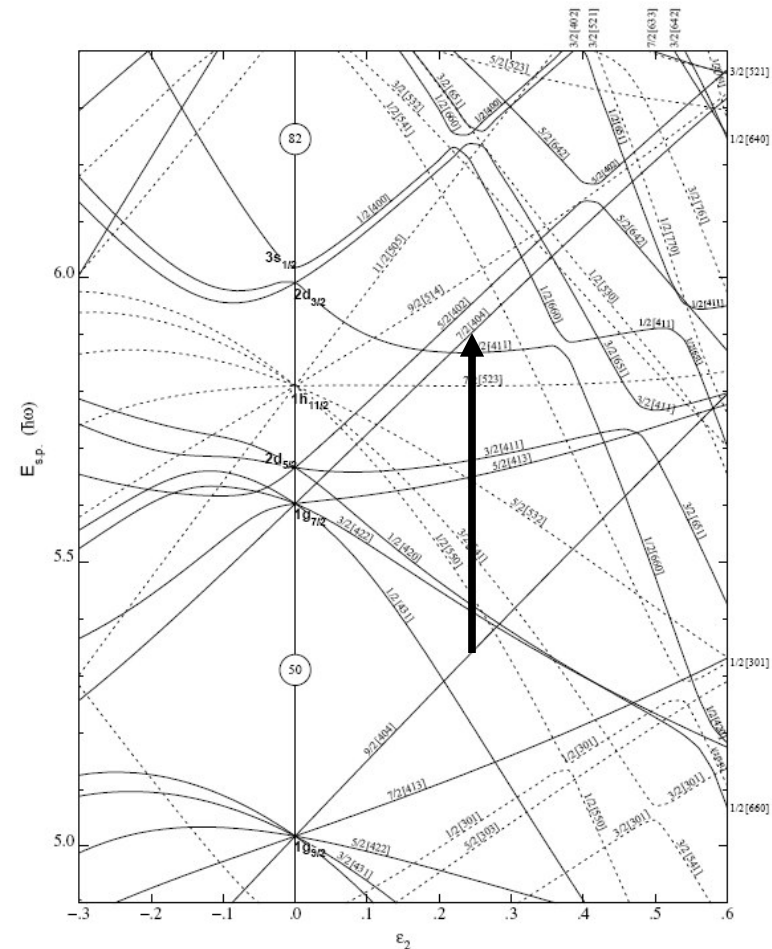


Figure 12. Nilsson diagram for protons,  $50 \leq Z \leq 82$  ( $\epsilon_x = -\epsilon_y^2/6$ ).



# Conclusions

- Statistical model analysis of thermal neutron capture  $\sigma_\gamma$  data can be used to accurately determine total radiative cross sections  $\sigma_0$  and improve level  $J^\pi$  values.
- Preliminary evidence of nuclear structure effects, possible K dependence, are seen in the neutron capture data.
- Strong evidence of nuclear structure effects is seen in the  $\beta$ -strength function for the decay of nuclei far from stability.