





#### Why ?...(very brief)

#### Pygmy Resonance may have an very important impact on the r-process <u>nucleosynthesis</u>

Giant resonances are of paramount important for nuclear astrophysics. Often, relevant reaction rates under astrophysical conditions are dominated by giant-resonance contributions, frequently in unstable nuclei. For instance, neutron-rich nuclei with loosely bound valence neutrons may exhibit very strong ( $\gamma$ ,n) strength components near particle threshold and thus, in turn, enhanced neutron-capture rates. Re-

Nupecc long range plan 2004



S.Goriely, Phys. Lett. B436 10 (1998) S.Goriely and E. Khan, Nucl. Phys. A706 (2002) 217

### <u>AND</u> • how collective properties change with neutron number







### <u>Stable nuclei</u> ⇒

photon scattering,Photoabsorption  $(\gamma,\gamma'),(\gamma,n)...$ 



<u>Erotic nuclei</u>⇒ <u>Virtual photon breakup</u> <u>or Virtual photon scattering</u>





#### Soft E1 Excitation with Coulomb Breakup of <sup>11</sup>Be,...

T.Nakamura *et al.*, PLB 331,296(1994) N.Fukuda et al., PRC70, 054606 (2004)



#### Virtual photon scattering technique (1)

- Peripheral heavy-ion collision on a high Z target at relativistic energies
- Virtual photon excitation and decay





### High resolution $\gamma$ -spectroscopy at the FRS

<sup>68</sup>Ni beam produced by fragmentation of <sup>86</sup>Kr @ 900 MeV/u on thick Be target (4g/cm<sup>2</sup>):



### Coulomb excitation of <sup>68</sup>Ni @ 600 AMeV



### <u>FRS+RISING ARRAY</u>



#### Euroball 15 Clusters

Located at 16.5°, 33°, 36° Energetic threshold ~ 100 keV

Hector 8 BaF<sub>2</sub>

Located at 142° and 88° Energetic threshold ~ 2 MeV

Miniball 7 HPGe segmented detectors Located at 46°, 60°, 80°, 90° Energetic threshold ~ 100 keV

Beam identification and tracking detectors Before and after the target

### Coulomb excitation of <sup>68</sup>Ni @ 600 AMeV







~ 6 Days of effective beam time ~ 400 GB of data recorded

~ 1.10<sup>8</sup> ' good <sup>68</sup>Ni events ' Incoming+Outgoing <sup>68</sup>Ni





### Spectrum of the ejectile measured in the zero degree calorimeter (CATE)



#### **RESULTS** Coulomb excitation : <sup>197</sup>Au(<sup>68</sup>Ni@600AMeV,<sup>68</sup>Ni<sup>\*</sup>)<sup>197</sup>Au



## γ-rays spectrum of BaF<sub>2</sub> detectors



## **Data Analysis**

Coulomb excitation Yield is product of <u>3 terms</u>: Virtual photo number, photoabsorption cross section, Branching

$$\frac{d^2 \sigma_{C\gamma}}{d\Omega dE_{\gamma}} (E_{\gamma}) = \frac{1}{E_{\gamma}} \frac{dn_{\gamma}}{d\Omega} (E_{\gamma}) \sigma_{\gamma} (E_{\gamma}) R_{\gamma} (E_{\gamma}).$$

[... Beene, Bortignon, Bertulani ...]

#### Calculate the ground state γ-ray decay from a GR state following a Coulomb excitation

# **! Coulomb excitation probability is directly proportional to the Photonuclear cross section**

[Eisenberg, Greiner, Bertulani, Alder, Winther,...]

$$\frac{d^2\sigma_{C\gamma}}{d\Omega dE_{\gamma}} (E_{\gamma}) = \left( \frac{1}{E_{\gamma}} \frac{dn_{\gamma}}{d\Omega} (E_{\gamma}) \sigma_{\gamma} (E_{\gamma}) R_{\gamma} (E_{\gamma}) \right).$$

The functions  $n_{\pi\lambda}(\varepsilon)$  are called the *virtual photon numbers*, and are given by

$$n_{E1}(b,\varepsilon) = \frac{Z_1^2 \alpha}{\pi^2} \frac{\xi^2}{b^2} \left(\frac{c}{v}\right)^2 \left\{ K_1^2 + \frac{1}{\gamma^2} K_0^2 \right\}$$
  
= number of equivalent photons  
Does **NOT** depend on the nuclear structure.

**Equivalent(virtual)-photon method** *Flux of virtual photons per unit area impinging on collision partners.* 



## Gamma decay – Branching ratio and level density

$$\frac{d^2\sigma_{C\gamma}}{d\Omega dE_{\gamma}} (E_{\gamma}) = \frac{1}{E_{\gamma}} \frac{dn_{\gamma}}{d\Omega} (E_{\gamma}) \sigma_{\gamma} (E_{\gamma}) R_{\gamma} (E_{\gamma}),$$



\*[Beene, et al PLB (1985) and ref therein]

C.N. Gilbreth and Y. Alhassid, private communication Shell model Monte Carlo (SMMC) Y. Alhassid et al. PRL 99, 162504 (2007)

#### (GDR-PDR) Coulomb excitation of <sup>26</sup>Ne or <sup>68</sup>Ni at relativistic energies

The extraction of the B(E1) strength requires the estimation of the direct and compound  $\gamma$ -decay of the dipole state to the ground state [Beene, et al PLB (1985)]

$$R_{decay}(E1) = \left[\frac{\Gamma_{\gamma 0}^{\uparrow}}{\Gamma} + C\frac{\Gamma^{\downarrow}}{\Gamma}\frac{\left\langle\Gamma_{\gamma 0}^{CN}\right\rangle}{\left\langle\Gamma^{CN}\right\rangle}\right]$$

J.Beene et al PRC 41(1990)920

$$\left\langle \Gamma_{\gamma 0}^{CN} \right\rangle = \frac{2}{\pi} \frac{\Gamma_{\gamma 0}^{\uparrow}}{\Gamma^{\downarrow}} \frac{1}{\rho}$$

J.Beene et al PLB 164(1985)1

Direct 
$$\Gamma^{\dagger}_{\gamma_0}=rac{1}{3\pi^2(\hbar c)^2}E^2\sigma_{int}$$

 $\sigma_{int}$  = photo-absorption cross-section

The compound term depends on the <u>ratio</u> between the <u>gamma and total</u> <u>decay width</u>

The gamma decay width depends on The value of the <u>level density</u> <u>at the resonance energy</u>

Spreading width  $\frac{\Gamma^{\downarrow}}{\Gamma} \simeq 1$   $\frac{\Gamma^{\downarrow}}{\Gamma} < 1$ 

Heavier nuclei Lighter nuclei



#### IMPORTANT Cross CHECK !



#### Euroball Clusters of HPGe IMPORTANT Cross CHECK !











Euroball Clusters of HPGe



O. Sorlin et al., Phys. Rev. Lett. 88, 092501 (2002).

$$\sigma_{E2}^{(app)} = \frac{8\pi^2}{75} \frac{Z_T^2 \alpha}{\left(\hbar c\right)^3} E_x^3 B\left(E2\right) \left(\frac{c}{v}\right)^4 \left[\frac{2}{\gamma^2} K_1^2 + \xi \left(1 + \frac{1}{\gamma^2}\right)^2 K_0 K_1 - \frac{v^4 \xi^2}{2c^4} \left(K_1^2 - K_0^2\right)\right]$$

#### → With B(E2)=255e^2fm^4 → 18mBarn

 $b_{min} = 1.34 (ap^{(1/3)}+at^{(1/3)})-(ap^{(-1/3)}+at^{(-1/3)})=12.85 fm$ 



[10] Y. Alhassid et al., Phys. Rev. Lett. 99, 162504 (2007) and C.N.Gilbreth and Y.Alhassid, private communication

- [11] S. I. Al-Quraishi et al. Phys. Rev. C 63, 065803 (2001) and S. I. Al-Quraishi et al. Phys. Rev. C 67, 015803 (2003)
- [12] P. Demetriou, S. Goriely, Nucl. Phys. A695, 95 (2001)
- [13] W. Dilg, W. Schantl, H. Vonach, and M. Uhl, Nucl. Phys. A217, 269 (1973)

#### [10] Shell Model Monte Carlo Method

- [11] Extrapolation from known levels to exotic nuclei
- [12] HF-BCS MSk7 Skyrme http://www-nds.ipen.br/RIPL-2/densities.html
- [13] Back-shifted Fermi gas model
- (\*)*Microscopic LD* HF-Bogolyubov BSK 14 Skyrme *http://www-astro.ulb.ac.be/Nucdata/Nld\_comb\_ph/z028.tab* [S. Goriely, S. Hilaire and A.J. Koning, Phys. Rev. C78 (2008) 064307]

<LD> (-1) around E\*=11 MeV: Al-Quraishi-ca.50 BRUSLIB\*-ca.180 Alhassid-ca.500 Empire.-ca.600 Cascade-ca.2000

#### Coulomb excitation of <sup>67</sup>Ni (600 MeV A) . . .





Predictions are available only for <sup>68</sup>Ni for <sup>67</sup>Ni we see a strength at an energy 2 MeV lower. As a simple rule the energy of the PDR should be correlated to the neutron binding energy

 $E_{b}$  (<sup>68</sup>Ni) ~ 7.8 MeV  $E_{b}$  (<sup>67</sup>Ni) ~ 5.8 MeV

#### ..... In progress ... to extract information on the skin radius







A. Klimkiewicz et al. PHYSICAL REVIEW C 76, 051603 ®(2007)

FIG. 4. (Color online) Fraction of the energy-weighted sum rule contained in the low-energy region (5-10 MeV) relative to that in the high-energy region (10-25 MeV) as a function of the neutron skin of the various Sn-isotopes.

## **<u>Planed next Measurements</u>** Predictions in <sup>XX</sup>Ne

System Neon Massh Number by QRRPA calculations



Cao L.-G. and Ma Z.-Y. Phys. Rev. C 71, 034305 (2005)



Figure 2: Dipole strength distributions calculations for <sup>26</sup>Ne in function of the gamma ray energy. In the left panel predictions by QRRPA calculations [17] are shown, in the right panel the results of HFB+QRPA calculations are shown [18]. The two different calculated strength distributions for <sup>26</sup>Ne in function of the gamma ray energy gives for the PDR state, 4,5% (of TRK sum rule) at 8 MeV and 1,5% at 10.5 MeV respectively. PRL 101, 212503 (2008)



<u>Planed next Measurements</u>

Pygmy in light nuclei :

 e.g. Ne neutron rich nuclei (measured at RIKEN with virtual photon break-up)
 improve the count rate of the set up because it has a smaller cross section than <sup>68</sup>Ni

Pygmy in neutron rich nuclei in mass 100-130

□ Improve the energy resolution of the particle detector after target

We <u>need</u> level density calculations (or measurements) for E around the threshold for proposals, and analysis

Even-odd nuclei?

#### CONCLUSION

> Measurement of high energy  $\gamma$ -rays from Coulex of <sup>68</sup>Ni at 600 MeV/u.

- > Strength at 11 MeV has been observed in all three kind of detectors
- > We found an extra strenght at 11 MeV with around 5% (+/- 1.5%) of the EWSR. The error is related by the assumption of the Level density.
- > The theory (RMF and RRPA calculations) predicts 4-8%.
- > The results opens new perspectives for other experiments and are very promising for Future measurements especially with high resolution

#### <u>Measurements with stable beams</u>



1.Restoration of the <u>Isopin mixing</u> in N=Z  $^{80}$ Zr at high temperature studied with the measurement of the  $\gamma$ -decay of GDR

Method and ("Theoretical") backgroundMeasurementOutlook





### Introduction (very brief and fast)

**Isospin**  $(I_z = \frac{1}{2}(N-Z))$  is a **good** quantum number for the **strong** interaction and therefore the Hamiltonian is symmetric against exchange of n and p

→ Coulomb interaction is isospin violating/breaking → MIXING (Isospin is not conserved)

• in the g.s and in low lying states: isospin mixing  $\alpha$ , around <5%

#### **OPEN QUESTIONS**

• how does mixing <u>change</u> (with Temperature  $\rightarrow$  E\*) for Z and A ?? •How to measure this in un/stable N=Z nuclei ??

### <u>Isospin mixing</u> and <u>forbidden</u> E1 decay in N=Z

N=Z Nuclei : The electric <u>dipole</u> transitions (without mixing) in long- wavelength limit are strictly (ISOSPIN ) forbidden in states with the same isospin

Temperature= 0



isospin mixing !

#### Temperature > C



### Temperature > 0

how does mixing <u>change</u> (with Temperature) for Z and A ??
How to <u>measure</u> this in hot un/stable nuclei ??







•up to mass  $\approx$  40 the isospin mixing  $\alpha^2$ , increases similarly to its behaviour at T=0 MeV,

•while at A = 60 it does not increase as one would expect?

•Measurement at A=80 will help to clarify !



### **GARFIELD+HECTOR** *@* Laboratori Nazionali di Legnaro



## Preliminary Results Measurement

Stat. Model.



→Thermalized emission from "identical" both Compound systems at same E\* and T

## **Preliminary analysis:**



Isopin mixing in N=Z <sup>80</sup>Zr at high temperature

It was possible to measure symmetric <sup>40</sup>Ca+<sup>40</sup>Ca=<sup>80</sup>Zr reaction at high temperature, not accessible with other methods

Statistical model simulations together with experimental data and teoretical analysis will give **insight** into restoration of Isospin mixing at high temperature

#### >FUTURE:

<b>Beams to use:</b>	<sup>44</sup> Ti, <sup>56</sup> Ni, <sup>72</sup> Kr …?				
Nuclei to study:	<sup>84</sup> Mo, <sup>96</sup> Cd, <sup>112</sup> Ba,	Proj.	Target	CN	
		<sup>44</sup> Ti	<sup>24</sup> Mg	<sup>68</sup> Se	
<b>Method:</b> excite GDR in $N = Z$ nucleus by		<sup>56</sup> Ni	<sup>12</sup> C	<sup>68</sup> Se	
complete fusion reaction		<sup>56</sup> Ni	<sup>24</sup> Mg	<sup>80</sup> Zr	
measure gamma-ray (and light particle) spectra		<sup>44</sup> Ti	<sup>40</sup> Ca	<sup>84</sup> Mo	
analyze GDR statistical gamma-decay		<sup>56</sup> Ni	<sup>28</sup> Si	<sup>84</sup> Mo	
		<sup>72</sup> Kr	<sup>12</sup> C	<sup>84</sup> Mo	
		<sup>56</sup> Ni	<sup>40</sup> Ca	<sup>96</sup> Cd	
		<sup>72</sup> Kr	<sup>24</sup> Mg	<sup>96</sup> Cd	
		<sup>72</sup> Kr	<sup>28</sup> Si	<sup>100</sup> .5n	

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