

# *Structure of hot nuclear states*

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Argonne National Laboratory

2nd Workshop on  
Level Density and Gamma Strength  
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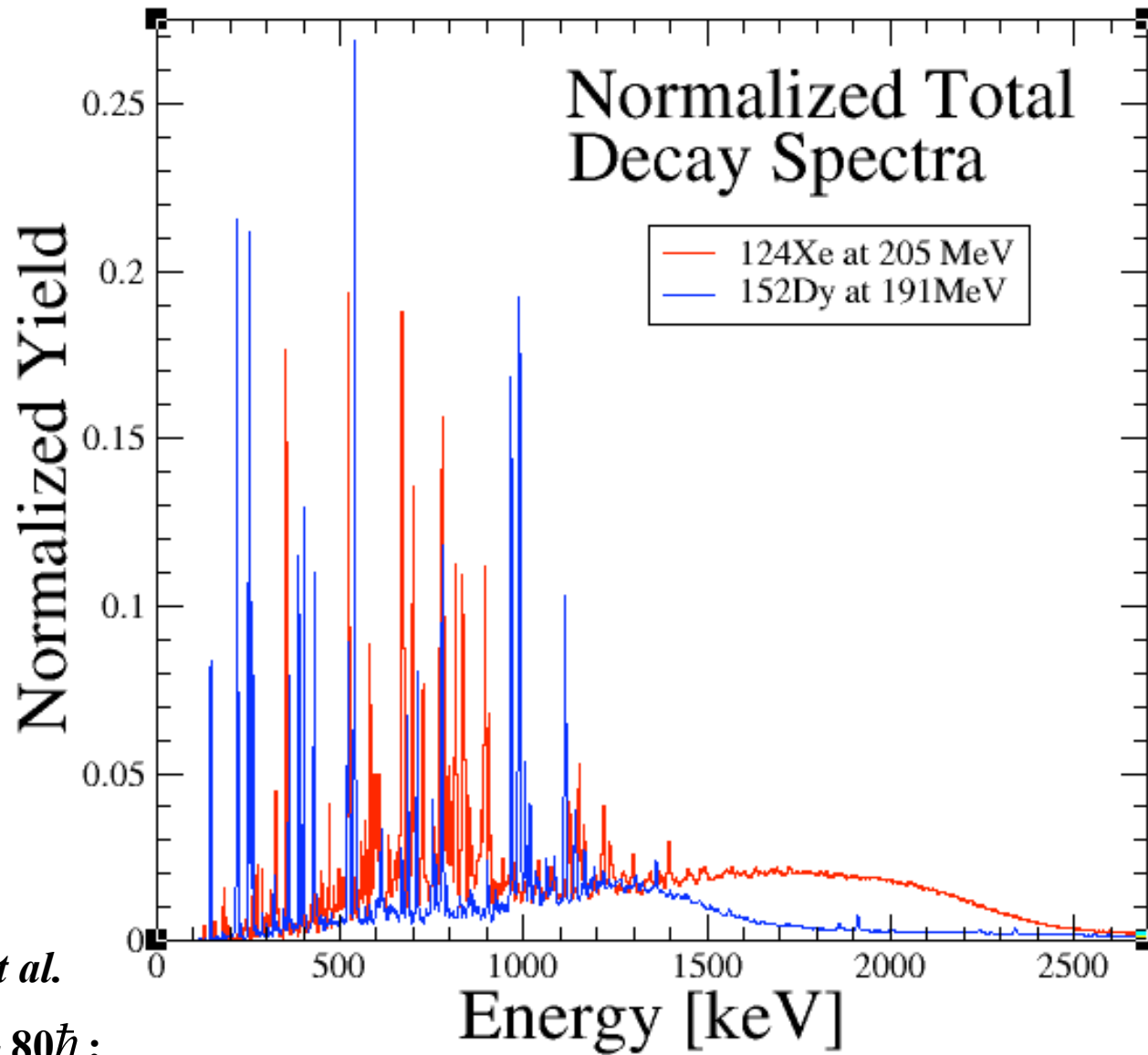
# Samuel Taylor Coleridge

## The Rime of the Ancient Mariner

*And the Albatross begins to be avenged.*

Water, water, every where,  
And all the boards did shrink ;  
Water, water, every where,  
Nor any drop to drink.

$\gamma$ s,  $\gamma$ s every where,  
And all the spectra did shrink ;  
 $\gamma$ s,  $\gamma$ s every where,  
Nor any physicist to drink.



Lauritsen *et al.*

$^{124}\text{Xe}$ ,  $l_{\text{max}} \sim 80\hbar$ ;

$\sim 1/2$  of  $\gamma$ s unresolved, from hot states.

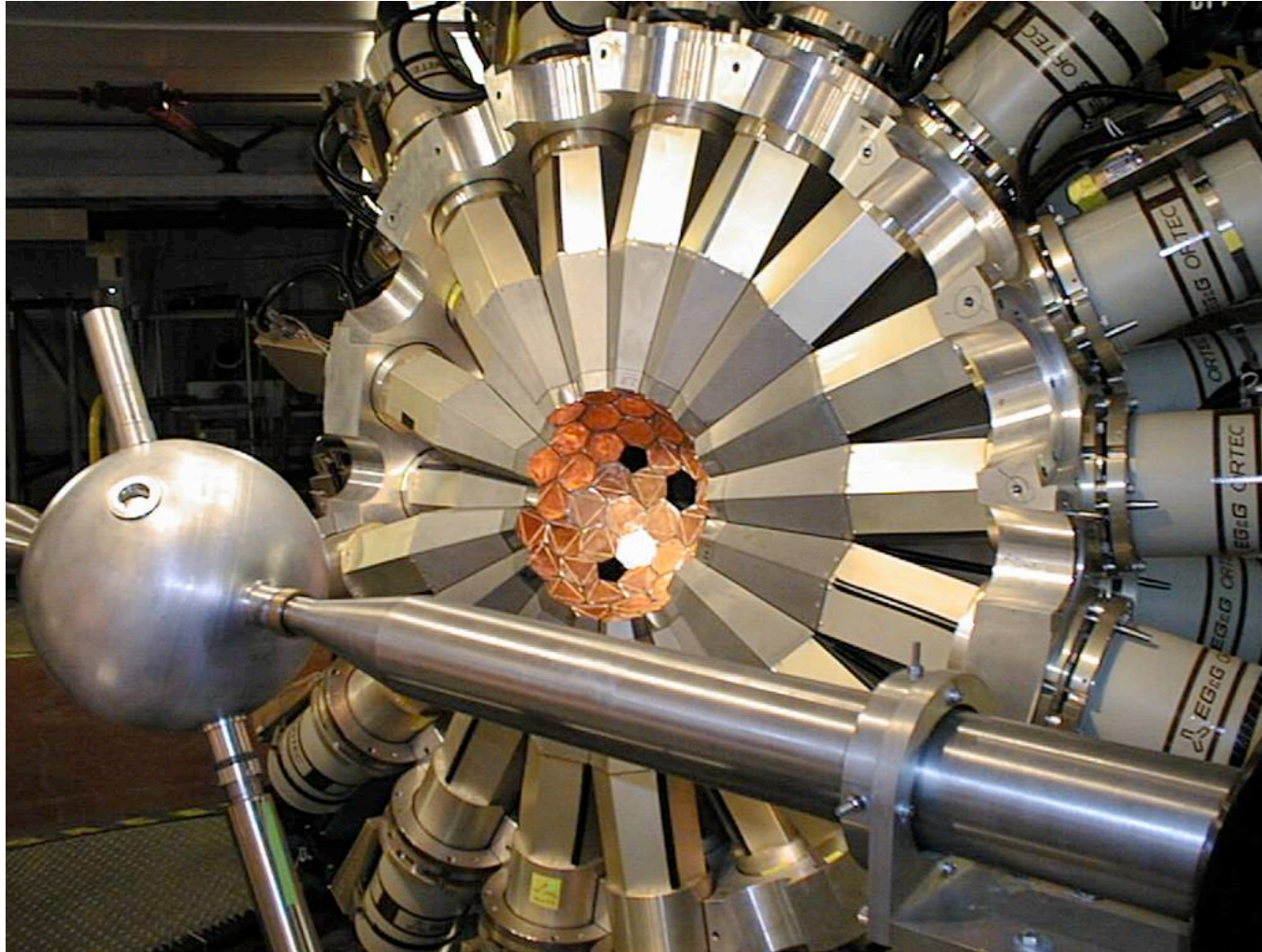
# Topics

- Phase transitions
- Fission barriers and  $l_{\max}$  in Superheavy Nuclei
- Order to chaos transition in superdeformed nuclei
- Ergodic superdeformed bands

Information extracted from  $\gamma$  spectra.

Theoretical description of  $\gamma$  spectrum requires knowledge of  $\rho$  &  $S_\gamma$ .

$\gamma$ s detected with Gammasphere  
@Argonne & Berkeley





# Adventures in the $\beta\gamma$ plane & above

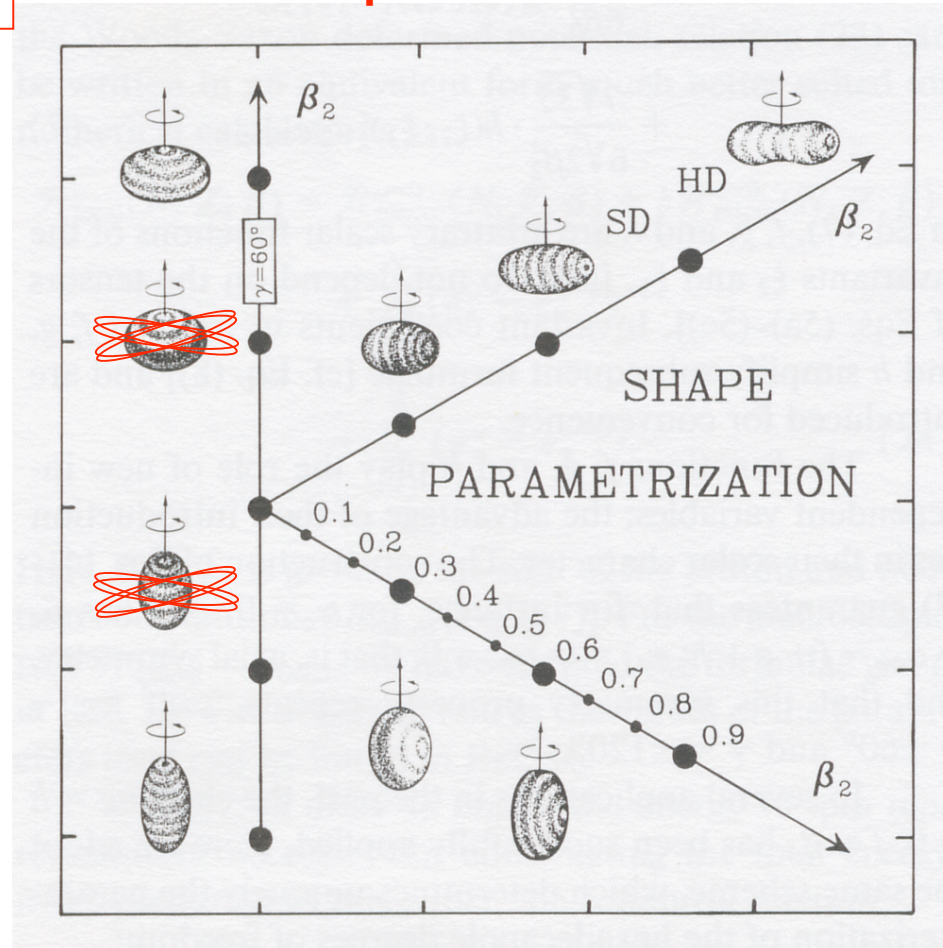
“Rotation”  
around symmetry  
axis

$J \sim J_0(1 + 0.6\beta)$

$J_0 = J_{\text{rigid sphere}}$

$J \sim J_0(1 - 0.6\beta)$

Oblate s.p.



prolate

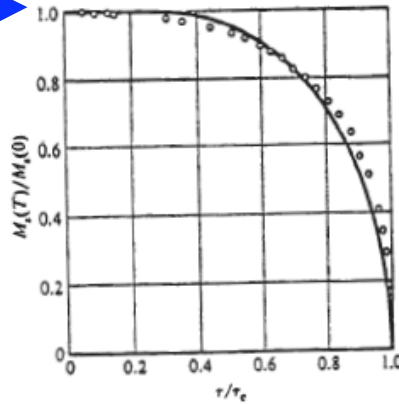
Collective oblate

prolate s.p.

Collective & aligned-particle rotation of nuclei

**Magnetization vs.  $T/T_c$**

Figure 10.16 Saturation magnetization of nickel as a function of temperature, together with the theoretical curve for spin  $\frac{1}{2}$  on the mean field theory.



**symmetry broken  
ordered phase**

**$\beta$  vs.  $T, I=0$   
 $^{170}\text{Er}$ , Goodman**

**deformed**

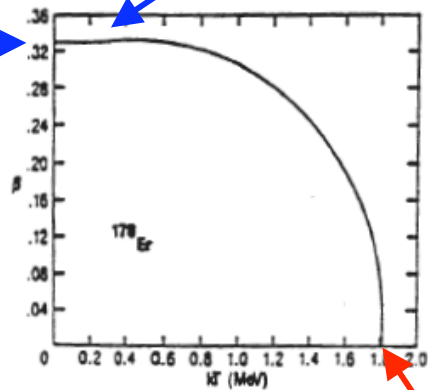
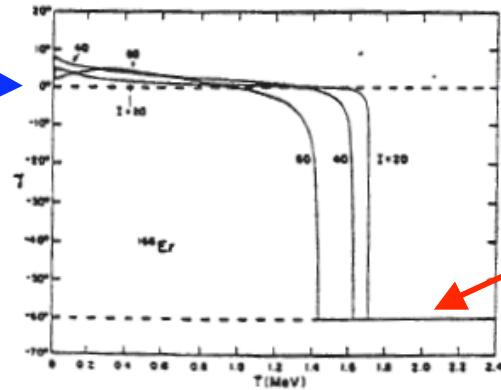


Fig. 2. The quadrupole deformation vs the temperature for  $^{170}\text{Er}$ .

**symmetry restored  
disordered phase**

**spherical**

**$\gamma$  vs.  $T, I=60, 40, 20 \hbar$   
 $^{164}\text{Er}$ , Goodman**



**$\gamma = 0^\circ$ , prolate**

**$\gamma = 60^\circ$ ,  
oblate**

**Universality class  
Critical exponent  $\beta$   
Order param  $\propto (1 - T/T_c)^\beta$**

**Phase Transitions**

## Phase Transitions above the Yrast Line in $^{154}\text{Dy}$

W. C. Ma,<sup>1</sup> V. Martin,<sup>2</sup> T. L. Khoo,<sup>3</sup> T. Lauritsen,<sup>3</sup> J. L. Egidio,<sup>4</sup> I. Ahmad,<sup>3</sup> P. Bhattacharyya,<sup>5</sup>  
M. P. Carpenter,<sup>3</sup> P. J. Daly,<sup>5</sup> Z. W. Grabowski,<sup>5</sup> J. H. Hamilton,<sup>6</sup> R. V. F. Janssens,<sup>3</sup> D. Nisius,<sup>3</sup> A. V. Ramayya,<sup>6</sup>  
P. G. Varmette,<sup>1</sup> and C. T. Zhang<sup>5</sup>

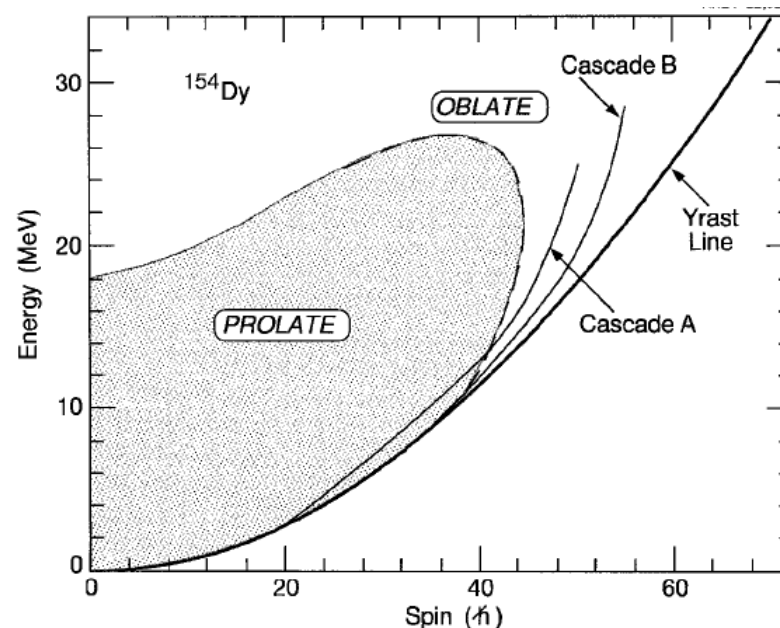
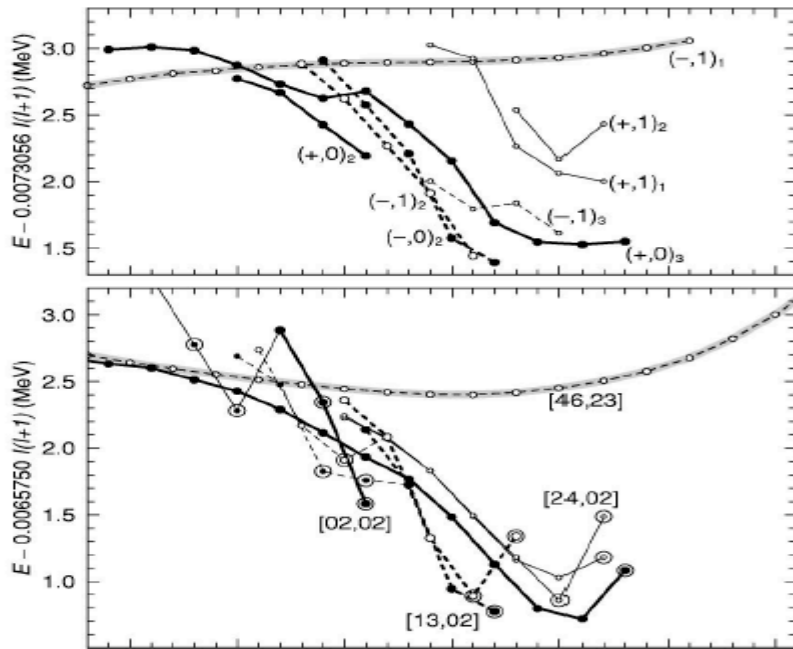


FIG. 1. Theoretical yrast line and regions of prolate (dotted) and oblate phases in  $^{154}\text{Dy}$ . The phase boundary (dashed line) corresponds to the  $\gamma = -60^\circ$  line in finite-temperature Hartree-Fock-Bogoliubov calculations without fluctuations [4]. Sketches of two cascade paths (A, B) are shown, which connect the experimental entry and exit points for cascades feeding into two selected regions of the yrast line,  $I = (16-22)\hbar$  and  $I = (34-36)\hbar$ .

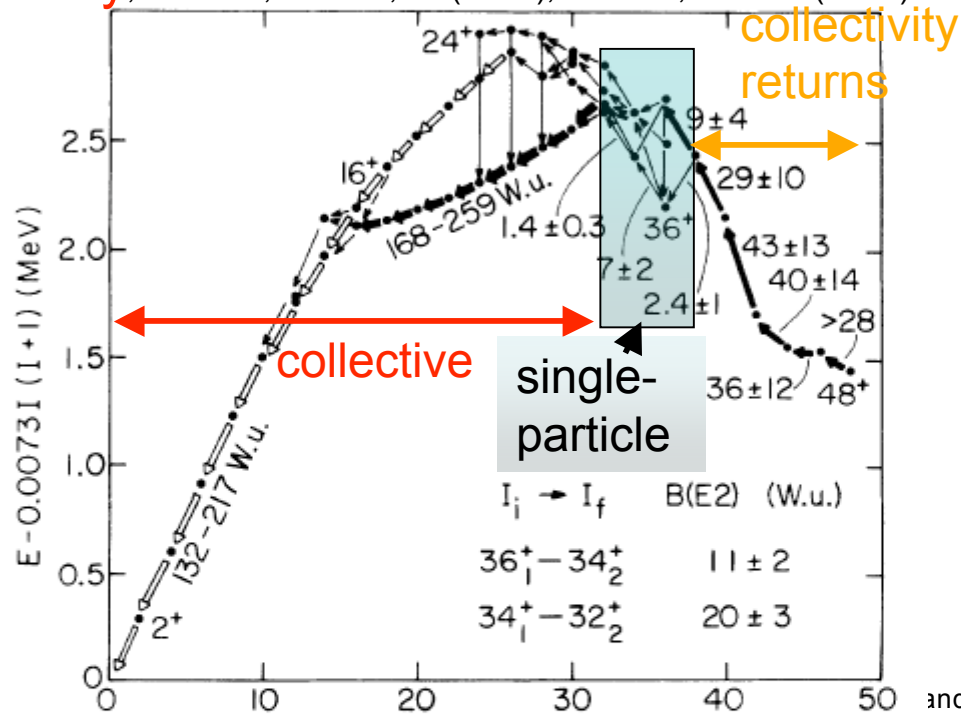


# Prolate ( $\gamma=0^0$ ) & oblate ( $\gamma=60^0$ ) states in N=88, 90 nuclei

- $N \leq 86$ , e.g.  $^{152}\text{Dy}$ . Particle alignment ( $\gamma=60^0$ , oblate), yrast isomers.
- **$N = 88, 90$ , e.g.  $^{154,156}\text{Dy}$ . Both prolate & oblate  $\rightarrow$  rotational & terminating bands.**
- $N \geq 92$ , e.g.  $^{158}\text{Dy}$ , prolate rotors ( $\gamma=0^0$ ).



$^{154}\text{Dy}$ , Ma et al., PRL **61**, 46 (1988), PRC **70**, 034315 (2004)



## Phase Transitions above the Yrast Line in $^{154}\text{Dy}$

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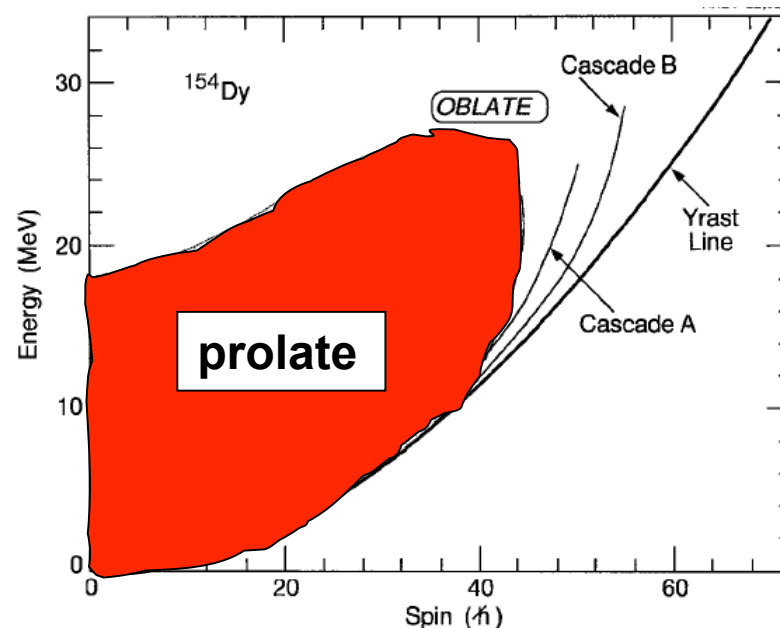
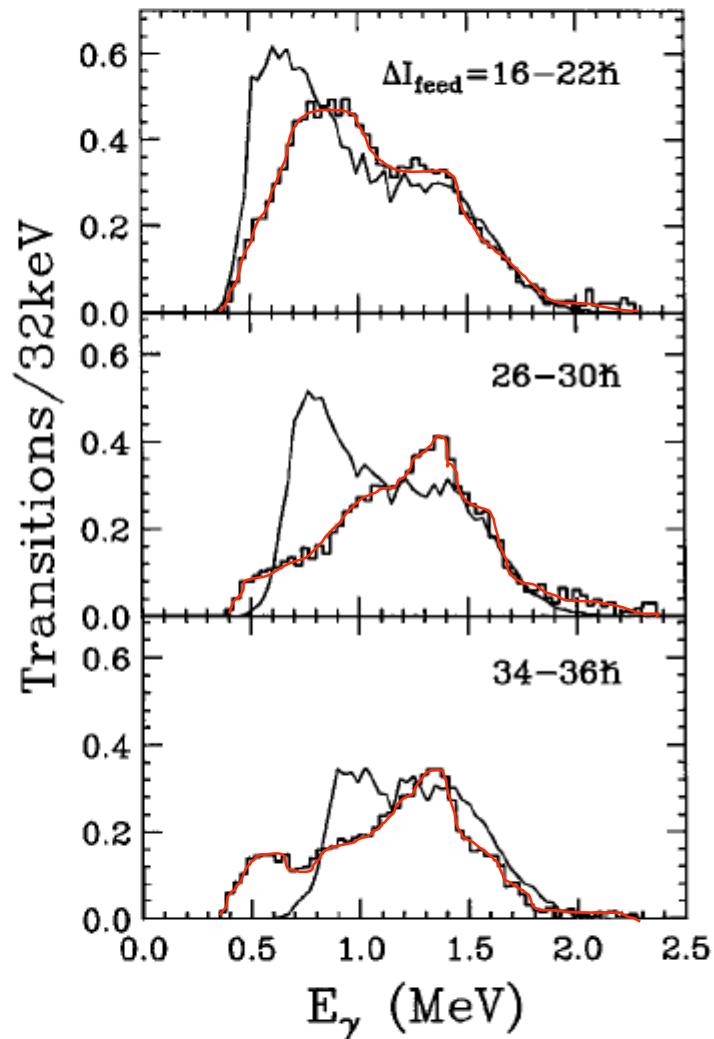


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Cascade A  
 oblate & prolate  
 zones  $\rightarrow$  **2 peaks**

Theory THFB ———  
 Experiment ———

Cascade B  
 oblate zone  $\rightarrow$  **1 peak**



Ma et al., PRL **84**, 5967 (2000)

**Remanent signatures  
 of phase transition at  
 finite temperature in  
 a mesoscopic system**

FIG. 3. Differential  $E2$  spectra feeding into the yrast line *only* in the indicated spin region  $\Delta I_{\text{feed}}$ . Histograms and solid lines correspond to experiment and theory. The approximate decay pathways corresponding to the top and the bottom spectra are shown as cascades A and B, respectively, in Fig. 1.

# Superheavy nuclei: at the limits of $Z, I, E^*$

## What are the limits?

### Physics questions.

- $B_f(I, E^*, Z, N)$  &  $E_{\text{shell}}(I, E^*, Z, N)$ .

- Variation with  $I, E^*$  (as well as  $Z, N$ )  $\rightarrow$  incisive tests of shell structure.

- Spectroscopy  $\rightarrow$  detailed tests of  $E_{\text{sp}}, e(\omega), E_{\text{band}}(I), J^{(1,2)}$  (i.e.  $\partial E / \partial I, \partial^2 E / \partial^2 I$ ).

# $\rho, S_\gamma$ in SHN

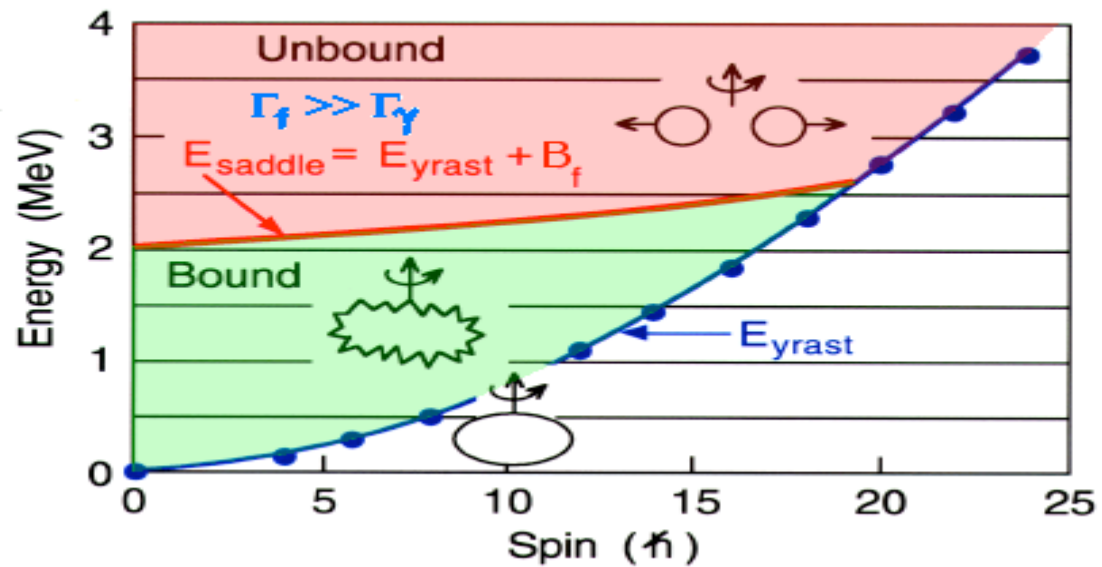
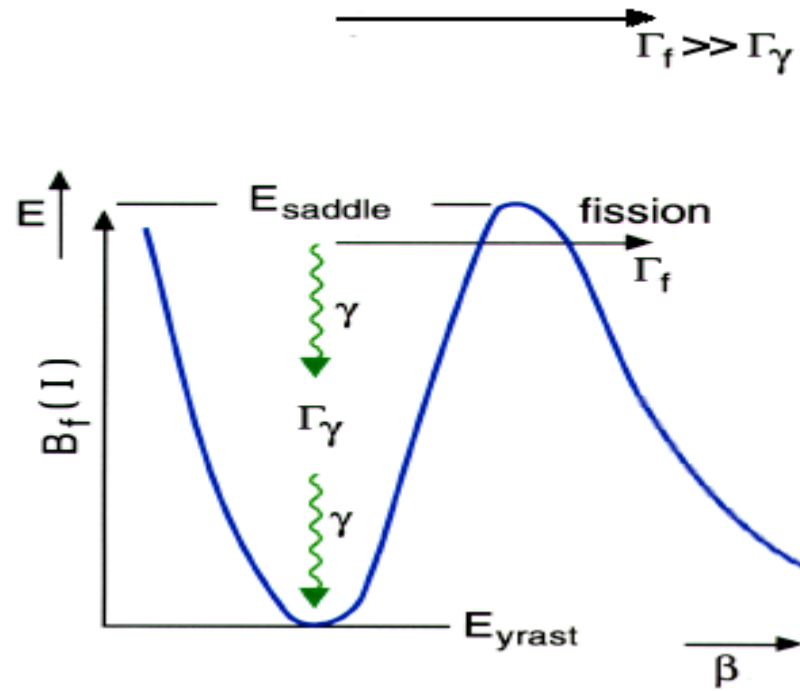
- $\rho, S_\gamma$  in SHN with transfer reactions,  
e.g.  $^{249}\text{Cf}(d,p), (\alpha, ^3\text{He}), (d,t), (^3\text{He}, \alpha)$
- Any new aspects, e.g. near top of fission barrier?



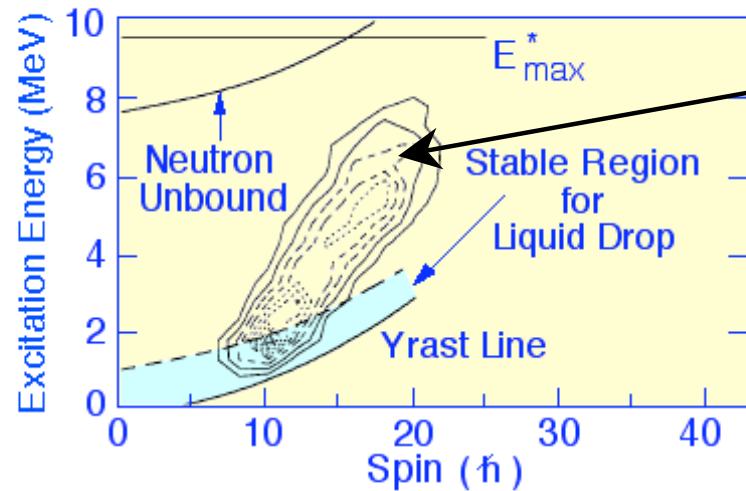
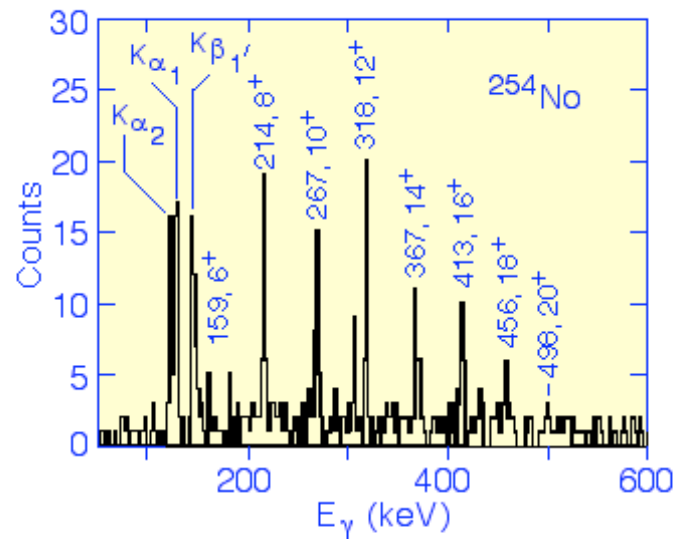
# How delicate are superheavy nuclei?

- How fragile are these loosely-bound nuclei, barely held together shell-created barriers?
- If you tickle them, will they “laugh” and fall apart?
- If you spin them, will they fission immediately?





## Limits of Stability in Spin and Energy



Initial states for  $\gamma$  decay to g.s.

$^{254}\text{No}$  is stable up to at least spin  $22\hbar$  and 8 MeV energy -- much more robust than expected.

$$B_f(\ell) > 5 \text{ MeV} \quad \ell \geq 10 \hbar$$

GammaSphere at ANL

Reiter et al.

PRL 84, 3542 (2000)

## Entry Distribution, Fission Barrier, and Formation Mechanism of $^{254}_{102}\text{No}$

P. Reiter,<sup>1,2</sup> T. L. Khoo,<sup>1</sup> T. Lauritsen,<sup>1</sup> C. J. Lister,<sup>1</sup> D. Seweryniak,<sup>1</sup> A. A. Sonzogni,<sup>1</sup> I. Ahmad,<sup>1</sup> N. Amzal,<sup>3</sup> P. Bhattacharyya,<sup>4</sup> P. A. Butler,<sup>3</sup> M. P. Carpenter,<sup>1</sup> A. J. Chewter,<sup>3</sup> J. A. Cizewski,<sup>1,5</sup> C. N. Davids,<sup>1</sup> K. Y. Ding,<sup>5</sup> N. Fotiades,<sup>5</sup> J. P. Greene,<sup>1</sup> P. T. Greenlees,<sup>3</sup> A. Heinz,<sup>1</sup> W. F. Henning,<sup>1</sup> R.-D. Herzberg,<sup>3</sup> R. V. F. Janssens,<sup>1</sup> G. D. Jones,<sup>3</sup> H. Kankaanpää,<sup>7</sup> F. G. Kondev,<sup>1</sup> W. Korten,<sup>6</sup> M. Leino,<sup>7</sup> S. Siem,<sup>1,8</sup> J. Uusitalo,<sup>1</sup> K. Vetter,<sup>9</sup> and I. Wiedenhöver<sup>1</sup> M. Asai<sup>1,10</sup>, D. Grayson<sup>1</sup>

<sup>1</sup>Argonne National Laboratory, Argonne, Illinois 60439

<sup>2</sup>Ludwig-Maximilians-Universität, Am Coulombwall 1, D-85748 Garching, Germany

<sup>3</sup>University of Liverpool, Liverpool L69 7ZE, England

<sup>4</sup>Purdue University, West Lafayette, Indiana 47097

<sup>5</sup>Rutgers University, New Brunswick, New Jersey 08903

<sup>6</sup>DAPNIA/SPhN, CEA Saclay, F-91191 Gif-sur-Yvette Cedex, France

<sup>7</sup>University of Jyväskylä, Jyväskylä, Finland

<sup>8</sup>University of Oslo, Oslo, Norway

<sup>9</sup>Lawrence Berkeley National Laboratory, Berkeley, California 94720

(Received 3 January 2000)

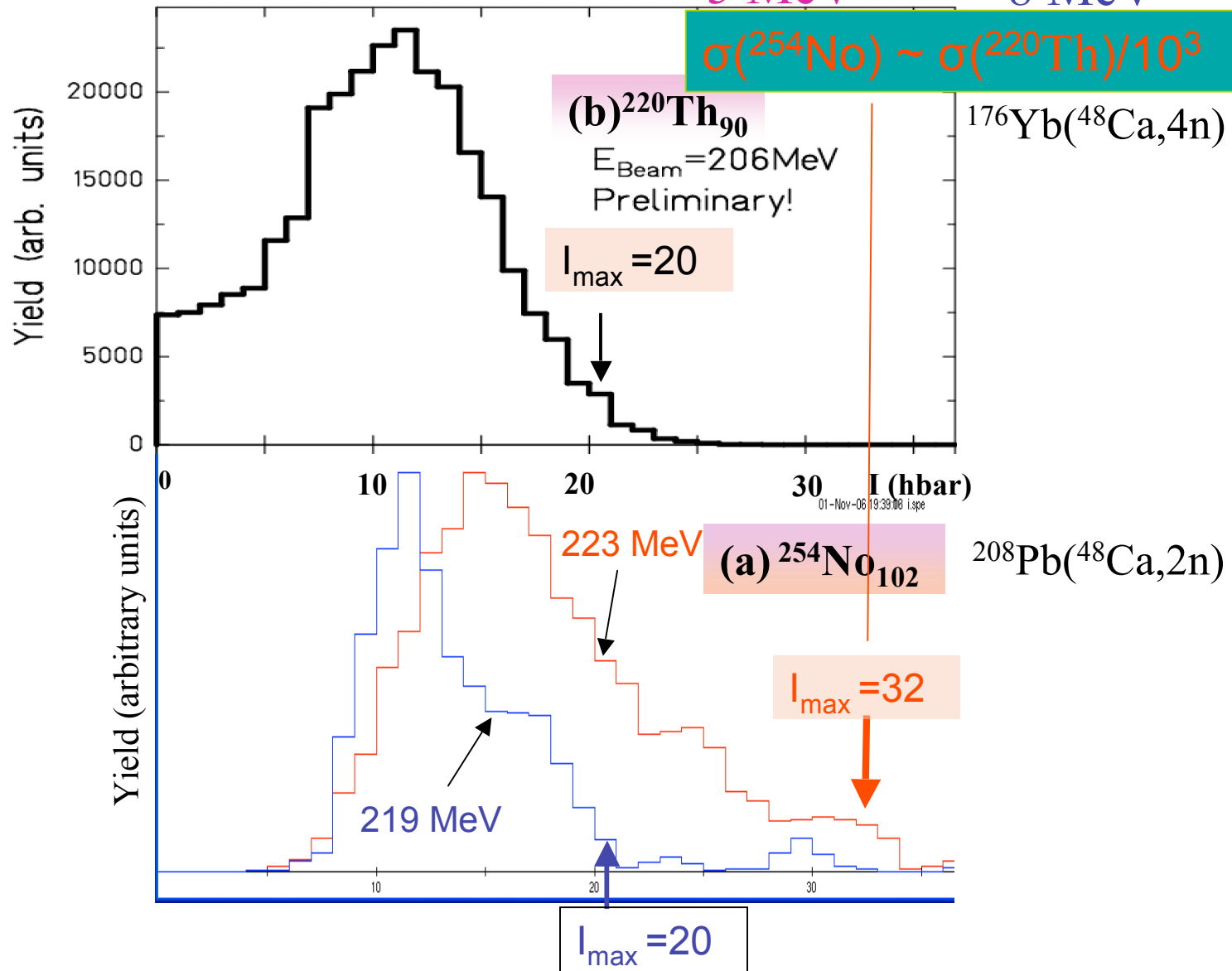
The entry distribution in angular momentum and excitation energy for the formation of  $^{254}\text{No}$  has been measured after the  $^{208}\text{Pb}(^{48}\text{Ca}, 2n)$  reaction at 215 and 219 MeV. This nucleus is populated up to spin  $22\hbar$  and excitation energy  $\geq 6$  MeV above the yrast line, with the half-maximum points of the energy distributions at  $\sim 5$  MeV for spins between  $12\hbar$  and  $22\hbar$ . This suggests that the fission barrier is  $\geq 5$  MeV and that the shell-correction energy persists to high spin.

$$I_{\max}(^{254}\text{No}) > I_{\max}(^{220}\text{Th})!$$

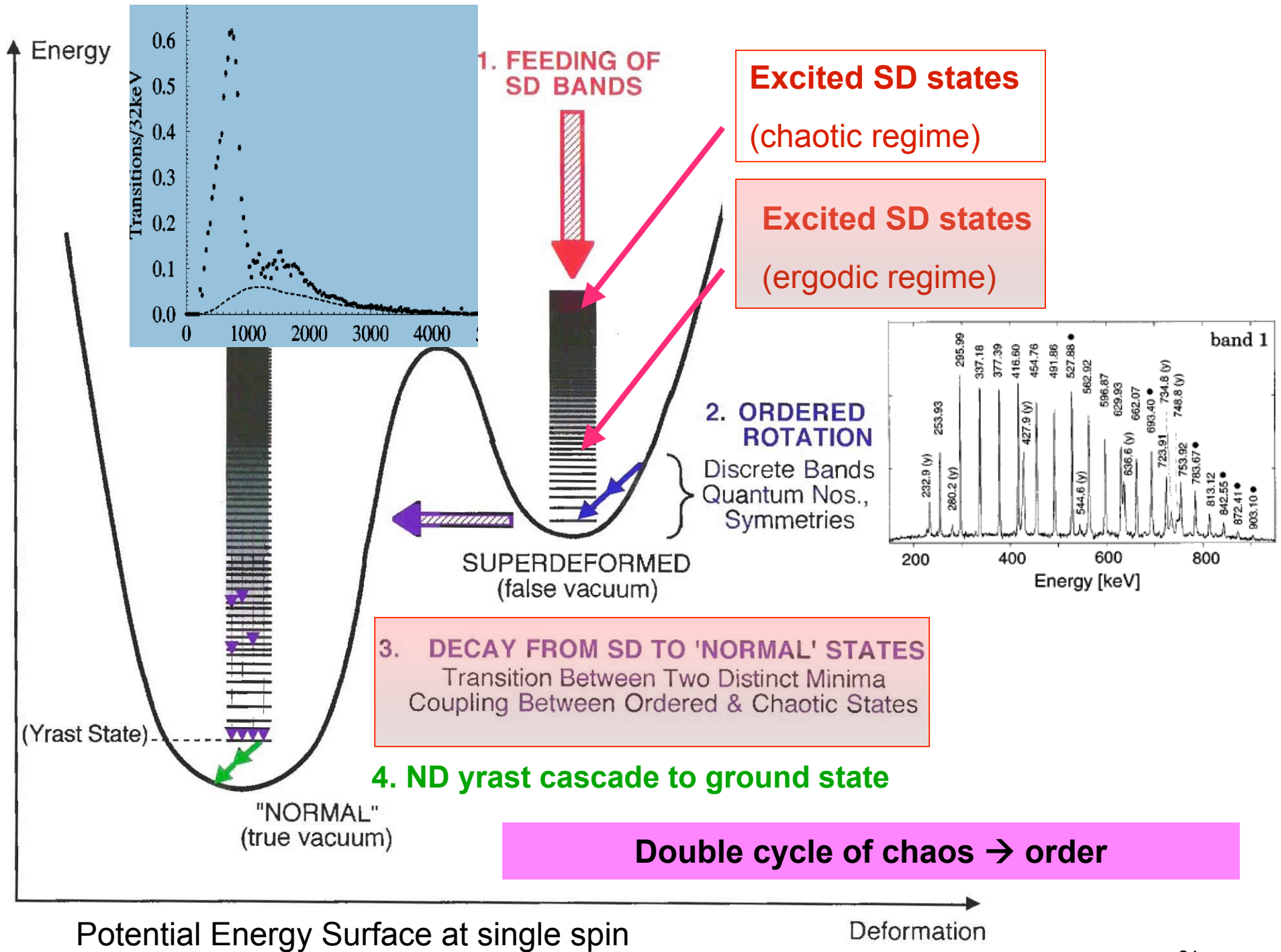
$$B_f(^{254}\text{No}) > B_f(^{220}\text{Th})$$

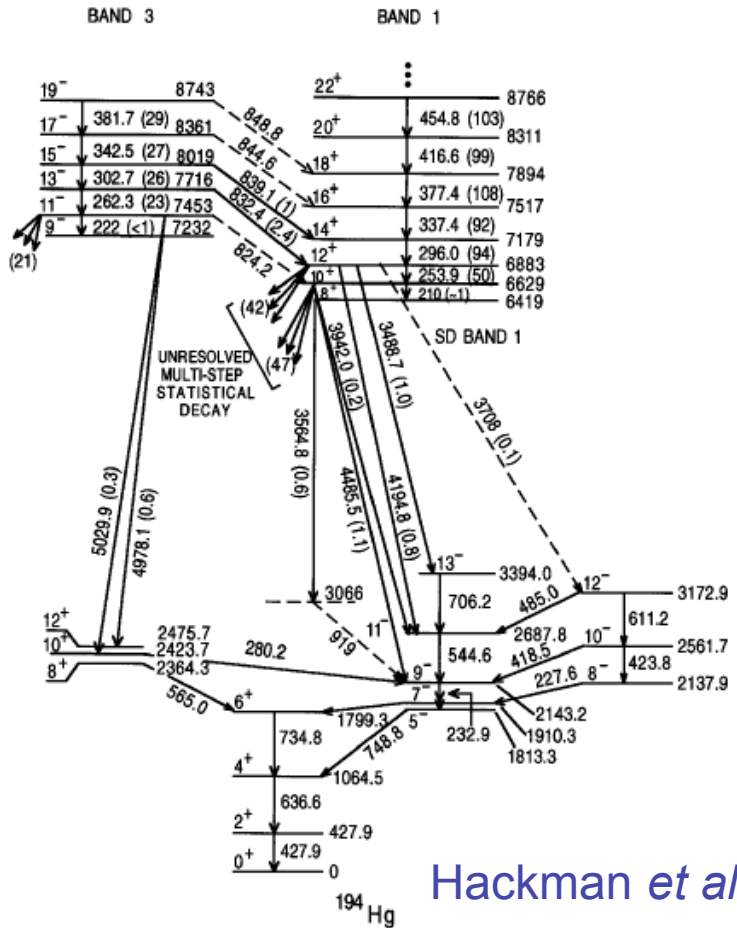
$$> 5 \text{ MeV} \quad \sim 8 \text{ MeV}$$

$$\sigma(^{254}\text{No}) \sim \sigma(^{220}\text{Th})/10^3$$

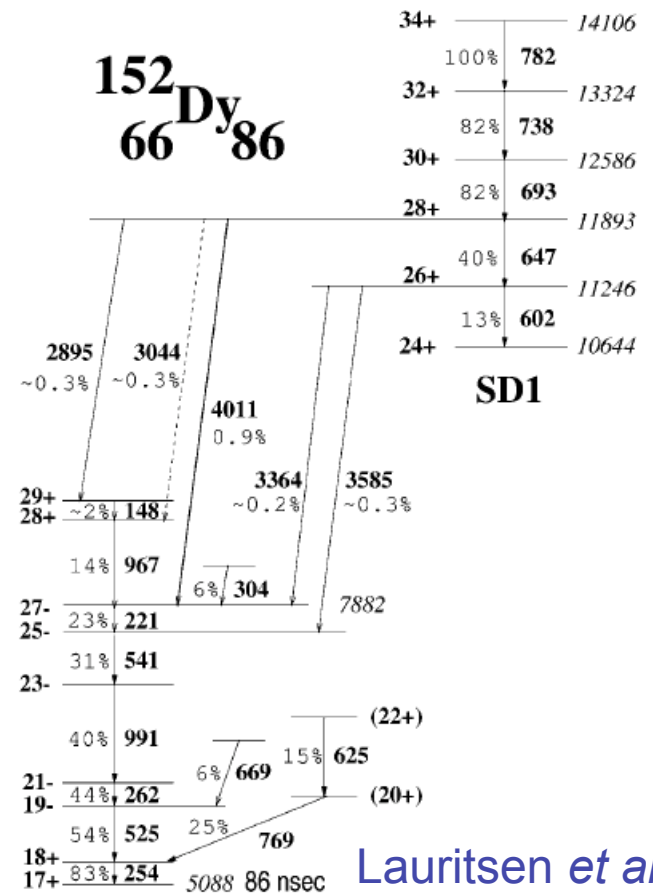
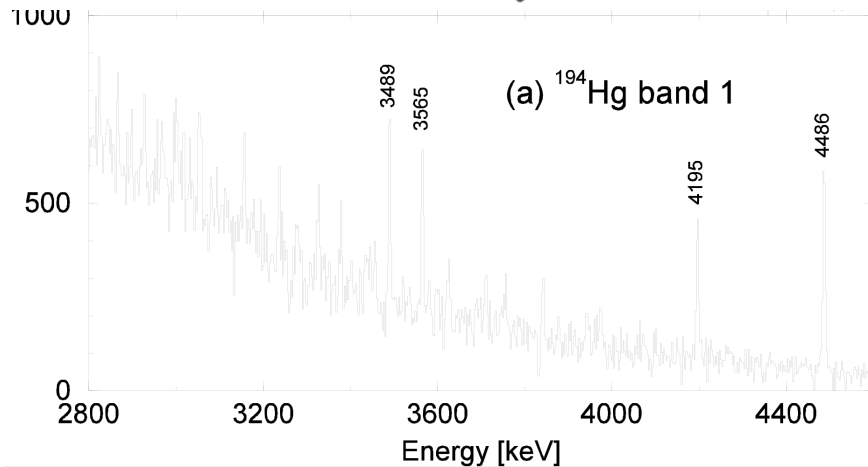




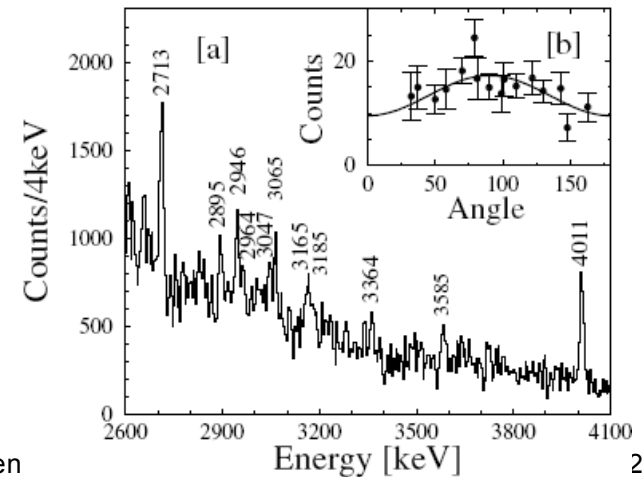




Hackman *et al.*



Lauritsen *et al.*



**Motional Narrowing and Ergodic Bands in Excited Superdeformed States of  $^{194}\text{Hg}$** 

Lopez-Martens,<sup>1</sup> T. Døssing,<sup>2</sup> T. L. Khoo,<sup>3</sup> M. Matsuo,<sup>4</sup> B. Herskind,<sup>2</sup> T. Lauritsen,<sup>3</sup> M. P. Carpenter,<sup>3</sup> R. V. F. Janssens,<sup>3</sup> G. Hackman,<sup>3,\*</sup> I-Y. Lee,<sup>5</sup> A. O. Macchiavelli,<sup>5</sup> E. Vigezzi,<sup>6</sup> and K. Yoshida<sup>7</sup>

<sup>1</sup>*C.S.N.S.M, IN2P3-CNRS, Batiment 104-108, 91405 Orsay, France*

<sup>2</sup>*Niels Bohr Institute, DK-2100, Copenhagen, Denmark*

<sup>3</sup>*Argonne National Laboratory, Argonne, Illinois 60439, USA*

<sup>4</sup>*Graduate School of Science and Technology, Niigata University, Niigata 950-2181, Japan*

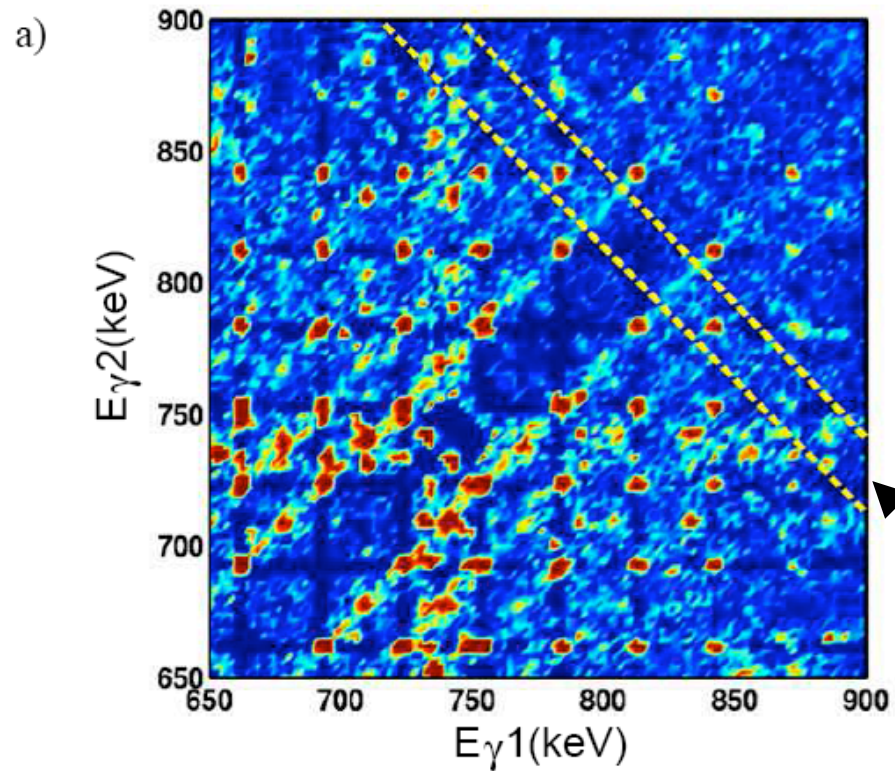
<sup>5</sup>*Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA*

<sup>6</sup>*INFN Sezione di Milano and Dipartimento di Fisica, Università di Milano, Via Celoria 16, 20133 Milano, Italy*

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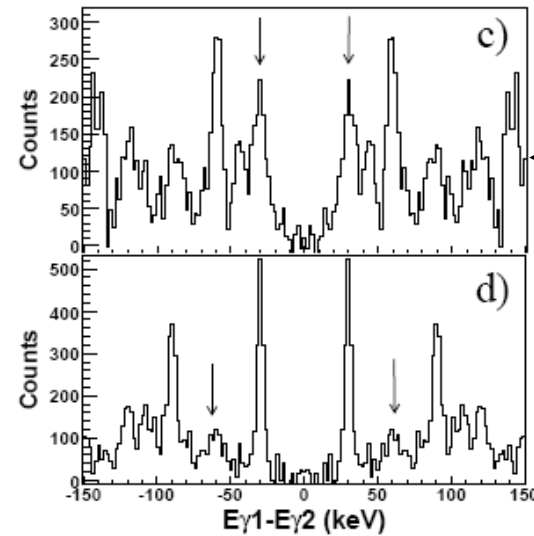
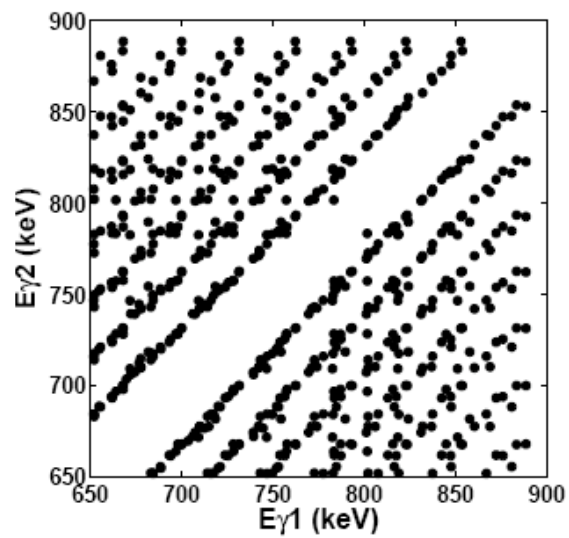
The  $E_\gamma$ - $E_\gamma$  coincidence spectra from the electromagnetic decay of excited superdeformed states in  $^{194}\text{Hg}$  reveal surprisingly narrow ridges, parallel to the diagonal. A total of 100–150 excited bands are found to contribute to these ridges, which account for nearly all the unresolved  $E2$  decay strength. Comparison with theory suggests that these excited bands have many components in their wave functions, yet they display remarkable rotational coherence. This phenomenon can be explained in terms of the combination of shell effects and motional narrowing.



$^{194}\text{Hg}$  SD gated  
Lopez-Martens *et al.*

$E_{\gamma} E_{\gamma}$  correlations

Slice c)



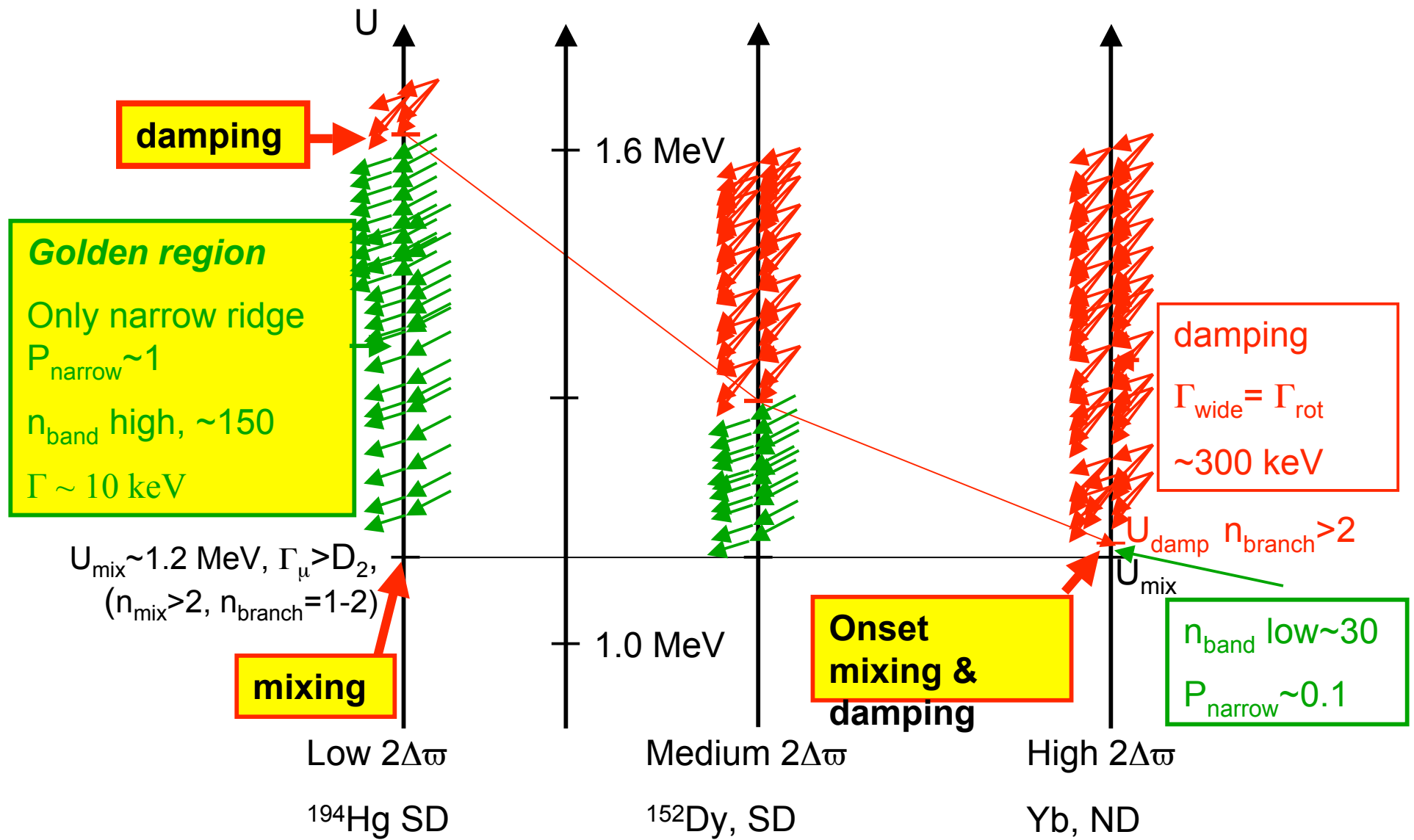
Discrete SD bands,  $A \sim 190$

Cuts perpendicular to diagonal

# Observations: *new phenomenon*

- Exceptionally narrow SD ridges (~~~10 keV vs. ~50 keV for ND~~).
- Ridge exhausts ~100% of E2 total strength (vs. ~10% for ND).
- No detectable broad component with rotational damping  $\Gamma_{\text{rot}} \sim 300$  keV.
- *Number of bands contributing to the ridge ~150 (vs. ~30 for ND).*
- *Cf.  $\text{FWHM}_{\text{SD}} \sim 10$  keV vs.  $\text{FWHM}_{\text{ND}} \sim 350$  keV*
- Narrow ridge implies E2 transitions flow within parallel rotational bands with nearly identical  $J^{(2)}$ .
- $J^{(2)}$  identical to that of SD band 1.
- Theory
  - predicted narrow ridges (Matsuo and Yoshida).
  - suggests that ~2-8 (4, average) basis configurations in excited SD states; from  $U=1.2-1.6$  MeV.
- Predicted by Mottelson







$$\Gamma_{E2} < D < \Gamma_{\mu}$$

$$\sim 10 \quad \sim 20 \quad \sim 40 \text{ keV}$$

Ergodic condition

$$2\Delta\omega < \Gamma_{\mu}$$

$$\sim 25 \quad \sim 40 \text{ keV}$$

Motional narrowing

### chaotic intrinsic-motion

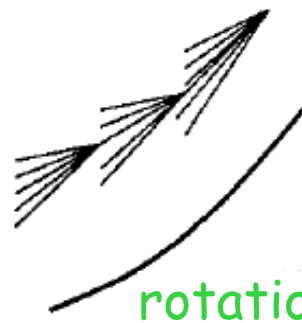
Chaotic, orderly  
rotational decay



ergodic  
bands

$n_{\text{branch}} \sim 1$

chaotic  
rotational decay



rotational  
damping

$> 2$

conditions for ergodic bands

small

$$\Delta\omega_0 \approx \frac{\Delta i}{\mathcal{I}}$$

large  $\varepsilon$ ,  $K$

Large  $\varepsilon$ ,  $A$

SD states in A  $\sim 190$   
small Coriolis intrn.

- Small  $\Delta\omega \rightarrow \sim$ identical  $J^{(2)} \rightarrow$ “identical” bands.
- Ergodic and identical bands intimately connected.

# From Order to Chaos

Excitation Energy above yrast line,  $U$  [MeV]

## Chaotic

Quantum numbers lost (except  $I, \pi$ ) → ***no selection rules.***

Statistical spectrum.

Transitions strengths unpredictable, governed by Porter-Thomas fluctuations.

## Ergodic

Mostly chaotic, with:

- (a) complicated wavefunctions and
- (b) Porter-Thomas fluctuations in all transitions except collective E2 transitions, with rotational band structure preserved. Unique in  $^{194}\text{Hg}$ .

## Ordered

Good quantum numbers → ***selection rules.***

Well-defined spectrum; equi-spaced (picket-fence) sharp lines from rotation of a cold SD object.

# Summary

- 1. Phase transitions along and above the yrast line.**
- 2. Superheavy nuclei:**  
**survive to  $I = 32$  hbar;**  
**shell structure robust to high spin.**
- 3. Superdeformed bands:**  
**double cycle of chaos-to-order transition;**  
**new ergodic regime with “orderly” E2 flow in chaotic states.**