

# Mass and Energy Dependence of the Spin Cutoff Parameter

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## Bethe(1936)

$$\rho(U, J) = \frac{n_s(U)}{\sqrt{2\pi\sigma}} \exp\left(\frac{-J^2}{2\sigma^2}\right)$$

$\sigma = \langle J_z^2 \rangle^{1/2}$  spin cutoff parameter

$$\begin{aligned}\rho_L(U, J) &= \rho_s(U, J) - \rho_s(U, J+1) \\ &= \left(\frac{d\rho_s}{dJ}\right)_{J=J+1/2}\end{aligned}$$

- Used to convert

$$\rho(U, 1/2) \text{ to } \sum_J \rho(U, J)$$

- at the binding energy

## Bethe(1936) cont.

$$\rho_{\text{total}}(U) = 2\sigma^2 \rho(U, 1/2)$$

$$\sigma^2 = \frac{I\theta}{\hbar^2}$$

$$I = \frac{2}{5}mR^2$$

$$\theta = \sqrt{\frac{U}{a}}$$

$$\sigma^2 = \frac{2mR^2}{5\hbar^2} \sqrt{\frac{U}{a}} \quad \text{Rigid Body Model (RBM)}$$

## Ericson (1960)

$$\sigma^2 = n \langle m^2 \rangle = \frac{6}{\pi^2} a \sqrt{\frac{U}{a}} \langle m^2 \rangle$$

- ▶ where  $n$  is the particle-hole number

$$= \frac{6}{\pi^2} \sqrt{U a} \langle m^2 \rangle \quad \text{Microscopic Model (MM)}$$

- ▶ Note:

$$\sigma^2(RBM) \propto a^{-1/2}$$

$$\sigma^2(MM) \propto a^{1/2}$$

- ▶ At closed shells,  $a$  drops, RBM  $\sigma^2$  increases at magic number, while MM  $\sigma^2$  is lower at magic number.

$$\sigma^2(RBM) \propto \sqrt{(U/a)} A^{5/3}$$

$$\sigma^2(MM) \propto \sqrt{U a} \langle m^2 \rangle$$

## Reconcile RBM and MM

- ▶ If  $a = \alpha A$

$$\sigma^2(RBM) \propto \sqrt{(U/\alpha)} A^{7/6}$$
$$\sigma^2(MM) \propto \sqrt{U\alpha} A^{1/2} < m^2 >$$

If  $< m^2 > \propto A^{\frac{2}{3}}$  the A dependence agrees.

- ▶ If we assume  $a = A/8$  ( $\alpha = 1/8$ ) then

$$< m^2 > = 0.1985 A^{\frac{2}{3}} \text{ makes MM equivalent to RBM}$$

- ▶ p shell  $< m^2 > / A^{2/3} = 0.144$
- ▶ sd shell  $< m^2 > / A^{2/3} = 0.164$
- ▶ fp shell  $< m^2 > / A^{2/3} = 0.175$
- ▶ sdg shell  $< m^2 > / A^{2/3} = 0.1824$
- ▶ This is close to the estimate above.

# Experimental Compilations

Fits to spins of known low-energy levels:

- ▶ Von Egidy-Bucurescu (2009)

$$\sigma^2 > \sigma^2(RBM) A < 50$$

$$\sigma^2 \leq 0.6\sigma^2(RBM) \quad A > 120$$

Does not approach  $\sigma^2(RBM)$  as  $U \rightarrow 20$  MeV.

- ▶ Al-Quraishi et al. (2003)  $20 \leq A \leq 110$

$$\sigma^2 \approx 0.6\sigma^2(RBM) \quad A \sim 20$$

$$\sigma^2 \approx 0.4\sigma^2(RBM) \quad A \sim 110$$

$$\sigma_{even-even}^2 < \sigma_{odd-A}^2 < \sigma_{odd-odd}^2$$

Does not approach  $\sigma^2(RBM)$  as  $U \rightarrow 20$  MeV.

- ▶ Capote et al (2009)

$$\sigma^2 = 0.5\sigma^2(RBM) A < 50$$

Does not approach  $\sigma^2(RBM)$  as  $U \rightarrow 20$  MeV.

## Experimental Summary

- ▶ Three compilations indicate  $\sigma^2 \approx 0.5\sigma^2(RBM)$ .
- ▶ Poor consistency.
- ▶ No clear indication of shell effects in  $A$  or  $U$ .

# Angular distributions of Hauser-Feshbach Processes

- ▶  $(p,p'), (n,n'), (p,n)$  nearly isotropic.
- ▶  $(\alpha, \alpha'), (\alpha, p), (\alpha, n), (p, \alpha)$  can yield  $\sigma^2$  values.
- ▶ Measurements from 1972-1978
  - ▶ Grimes et al. 1978:  $28 \leq A \leq 30$
  - ▶ Hille et al. 1974, Lu et al. 1972, Grimes et al. 1974:  $48 \leq A \leq 65$
  - ▶ Hille et al. 1974:  $A = 93, 95, 96$ , and 118
  - ▶  $\sigma^2$  at  $A = 30 \sim 0.6\sigma^2(RBM)$
  - ▶  $\sigma^2$  at  $A = 50 \sim 1.2\sigma^2(RBM)$
  - ▶  $\sigma^2$  at  $A = 60 \sim 0.8\sigma^2(RBM)$
- ▶ For  $A \sim 50$   $\sigma^2$  at  $U = 4$  MeV nearly equal to  $\sigma^2$  at  $U = 8$  MeV.
- ▶  $\sigma^2$  at  $A = 90 \sim 0.7 \sigma^2(RBM)$
- ▶ Data not comprehensive, but show oscillations of  $\sigma^2$  against  $\sigma^2(RBM)$  as a function of  $U$  and  $A$ .



# Calculations

- ▶ BCS Hamiltonian used with Statistical Mechanics by Sano and Yamasaki (1963) and L. Moretto (1972).
- ▶ Single particle energies Nilsson (1955), Seeger Perisho (1967), Seeger Howard (1975).
- ▶ Results show average  $\sigma^2$  Values for  $U$  and  $A$  averaged over the single particle basis.

## Calculations cont.

- ▶ Find  $\sigma_{even-even}^2 < \sigma_{odd-A}^2 < \sigma_{odd-odd}^2$  at a given  $U$  in a narrow  $A$  range.
- ▶ Generally have  $\sigma^2$  between 0.2 and 0.5 of  $\sigma^2(RBM)$  at 3-4 MeV.
- ▶  $\sigma^2$  as  $U \rightarrow 20$  MeV is usually within 10 % of  $\sigma^2(RPM)$ .
- ▶ In range of binding energy significant shell effects.
  - ▶ From  $A = 20$  to  $A = 30$   $\sigma^2$  drops by 20% although  $\sigma^2(RBM)$  increases 60%. Apparent reason is moving from filling the  $d_{5/2}$  orbital to filling the  $s_{1/2}$  and  $d_{3/2}$  orbitals.
  - ▶ A maximum is found at  $A = 50$  with  $\sigma^2 = 1.3\sigma^2(RBM)$
  - ▶ A minimum at  $A = 60$  with  $\sigma^2 = 0.8\sigma^2(RBM)$
  - ▶ A maximum at  $A = 80$  with  $\sigma^2 = 1.2\sigma^2(RBM)$
  - ▶ A minimum at  $A = 90, 140, 200$
  - ▶ A maximum at  $A = 120, 180$
  - ▶ Note: Magic numbers are close to minima  $A = 90, 190, 200$ , but also one maximum at  $A=120$ .

# Calculation Compilation

Table: Calculated spin cutoff  $\sigma^2$  parameters

A	U (MeV)	$\sigma^2$	$\sigma^2/\sigma^2(RBM)$	A	U (MeV)	$\sigma^2$	$\sigma^2/\sigma^2(RBM)$
20	2	1.2	0.6	60	2	2.1	0.3
	4	2.4	0.85		4	7.43	0.7
	6	5.9	1.7		6	11.14	0.9
	8	7.8	1.96		8	12.86	0.9
	10	8.1	1.81		10	13.65	0.85
	20	6.96	1.11		20	23.02	1.02
30	2	0.97	0.3	70	2	2.96	0.33
	4.0	2.6	0.58		4	9.71	0.8
	6	4.2	0.76		6	19.32	1.3
	8	5.43	0.85		8	24.32	1.42
	10	7.52	1.05		10	25.41	0.32
	20	11.0	1.1		20	30.11	1.11
40	2	1.36	0.31	80	2	5.0	0.5
	4	5.49	0.87		4	12.47	0.88
	6	11.29	1.46		6	22.58	1.3
	8	14.1	1.6		8	30.55	1.52
	10	16.1	1.7		10	32.6	1.45
	20	15.9	1.13		20	36.3	1.14
50	2	2.44	0.42	90	2	4.04	0.35
	4	8.04	0.98		4	10.58	0.65
	6	16.0	1.6		6	18.9	0.94
	8	16.1	1.38		8	20.7	0.9
	10	16.25	1.25		10	24.43	0.95
	20	17.4	0.95		20	38.2	1.05

# Calculation Compilation cont.

Table: Calculated spin cutoff  $\sigma^2$  parameters cont.

A	U (MeV)	$\sigma^2$	$\sigma^2/\sigma^2(RBM)$	A	U (MeV)	$\sigma^2$	$\sigma^2/\sigma^2(RBM)$
100	2	3.9	0.3	180	2	13.43	0.52
	4	15.1	1.0		4	33.14	0.91
	6	31.9	1.42		6	58.61	1.31
	8	23.66	1.29		8	78.3	1.52
	10	35.1	1.21		10	83.8	1.45
	20	47.3	1.15		20	94.75	1.16
120	2	10.5	0.65	200	2	8.8	0.3
	4	19.57	0.86		4	20.9	0.5
	6	28.98	1.04		6	27.2	0.54
	8	34.4	1.07		8	37.4	0.64
	10	39.22	1.09		10	46.7	0.72
	20	53.94	1.06		20	79.4	0.86
140	2	5.9	0.31	220	2	13.3	0.41
	4	19.1	0.7		4	34.3	0.74
	6	28.03	0.84		6	61.1	1.08
	8	34.68	0.85		8	73.0	1.12
	10	40.93	0.95		10	84.8	1.16
	20	66.8	1.1		20	119.0	1.15
160	2	9.0	0.40	240	2	17.0	0.53
	4	27.9	0.88		4	43.8	0.86
	6	44.8	1.15		6	80.3	1.29
	8	58.97	1.31		8	88.9	1.23
	10	64.94	1.29		10	93.2	1.15
	20	77.6	1.09		20	126.16	1.09

# Rigid Body Model

- ▶ Rigid Body Model has  $\sigma^2 \propto \frac{1}{\sqrt{a}}$  gives the maximum value of  $\sigma^2$  at closed shell.
- ▶ Microscopic Model
  - ▶  $\sigma^2 \propto \sqrt{a} < m^2 >$
  - ▶  $\sqrt{a}$  causes drop in  $\sigma^2$
  - ▶  $< m^2 >$  could enhance drop, not change it, or could cancel it.

## General Conclusions

- ▶ Microscopic Model goes to the rigid body model if  $\langle m^2 \rangle \propto A^{\frac{2}{3}}$  and  $a \propto A$
- ▶ Convergence  $\sim 20$  MeV to Rigid Body Model.
- ▶ Not yet converged at 6-8 MeV.

## Other Issues

### ► Parity Ratio

- Is  $\rho^+(U) = \rho^-(U)$  at the binding energy?
- Al-Quarishi et al. (2003) produced empirical fit based on known levels.
- Found deviations  $> 10\%$  at the binding energy for  $A \leq 90$ .

### ► Additional Questions

- Is  $\sigma^2(\text{Positive parity}) = \sigma^2(\text{Negative parity})$  at the binding energy?
- Gorieli et al. (2008) find  $\sigma_+^2$  and  $\sigma_-^2$  can differ by 20-25% as the binding energy is approached.

## Deformed Nuclei

- ▶ Bohr and Mottelson Proposed

$$\rho_{\text{deformed}}(U, J) = \sigma_{\perp}^2 \rho_{\text{spherical}}(U, J)$$

- ▶ Acknowledge not valid for all J.
- ▶ If this is used at the binding energy then,

$$\rho_{\text{total}}(U) = 2\sigma^2 \rho(U, \frac{1}{2})$$

.

- ▶ This can be shown to be incorrect.

Spherical Basis	Deformed Basis
J = 0	J = 0, K = 0
J = 2	J = 2, K = 0 J = 2, K = 1 J = 2, K = 2



## Deformed Nuclei cont.

- ▶ Get multiplier of  $J + 1$ .
- ▶ Bands also have higher multipliers for large  $J$ .
- ▶ New Rotational Enhancement Factor Grimes (2013).

$$R(J, K) = \frac{(J + 1)^2 - K^2}{(2J + 1)}$$

- ▶ This factor has dependence on  $K$  as well as  $J$ .
- ▶ When summed over  $J$  and  $K$  the deformed level density is about  $\sigma_{\perp}^2/2$  larger than the sum of the spherical level density.
- ▶ Since the enhancement factor is small for small  $J$ , the corrected level density for resonance counting is *increased*

# Conclusions

- ▶ Have compared Rigid Body and Microscopic Model predictions for  $\sigma^2$
- ▶ Find  $A$  and  $U$  dependence similar if  $a \propto A$  and  $\langle m^2 \rangle \propto A^{2/3}$
- ▶ Data base limited: low  $U \rightarrow \sigma^2 \sim \frac{1}{2}\sigma^2(RBM)$
- ▶ Some evidence for shell modulations in calculations with microscopic model for  $U \leq 10$  MeV.
- ▶ Goes to  $\sigma^2(RBM)$  as  $U \rightarrow 20$  MeV
- ▶ Two body model calculations agree with data and generally with Microscopic Model.
- ▶ Need reanalysis of level densities for deformed nuclei.
- ▶ Need more measurements of  $\sigma^2$  in a variety of  $A$  ranges.