

# *Measuring Neutron Capture Cross Sections on s-Process Radioactive Nuclei*

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 Lawrence Livermore  
National Laboratory

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Density and Gamma Strength  
Oslo, May 18 - 22, 2015**

**LLNL-PRES-670315**

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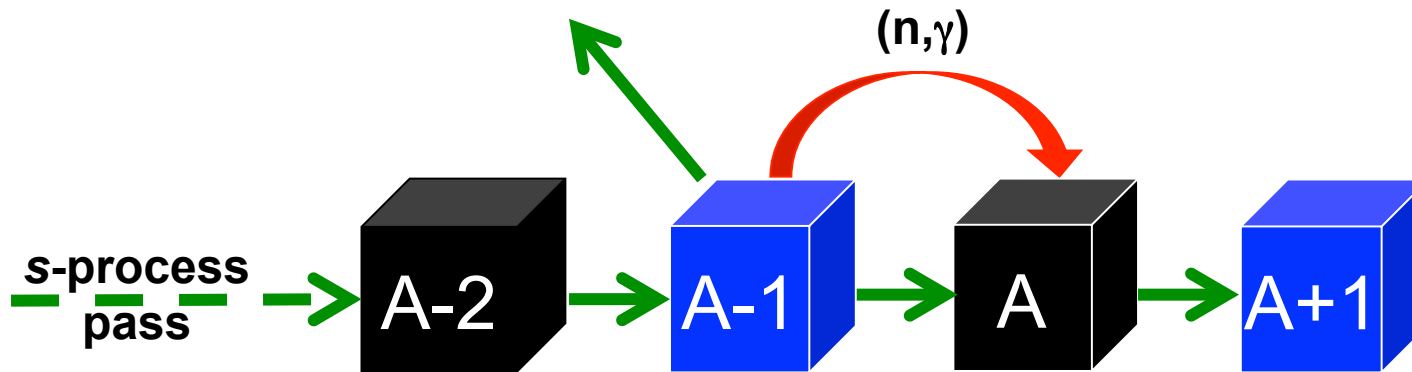


# Outline

- ❑ Low-energy dipole strength: a key to understand decay mechanism for capture reactions
- ❑ Experimental techniques
- ❑ Experimental results for capture cross-section measurements on *s*-process branching point nucleus:  $^{85}\text{Kr}$ ,  $^{75}\text{Ge}$ , and  $^{205}\text{Pb}$
- ❑ Summary



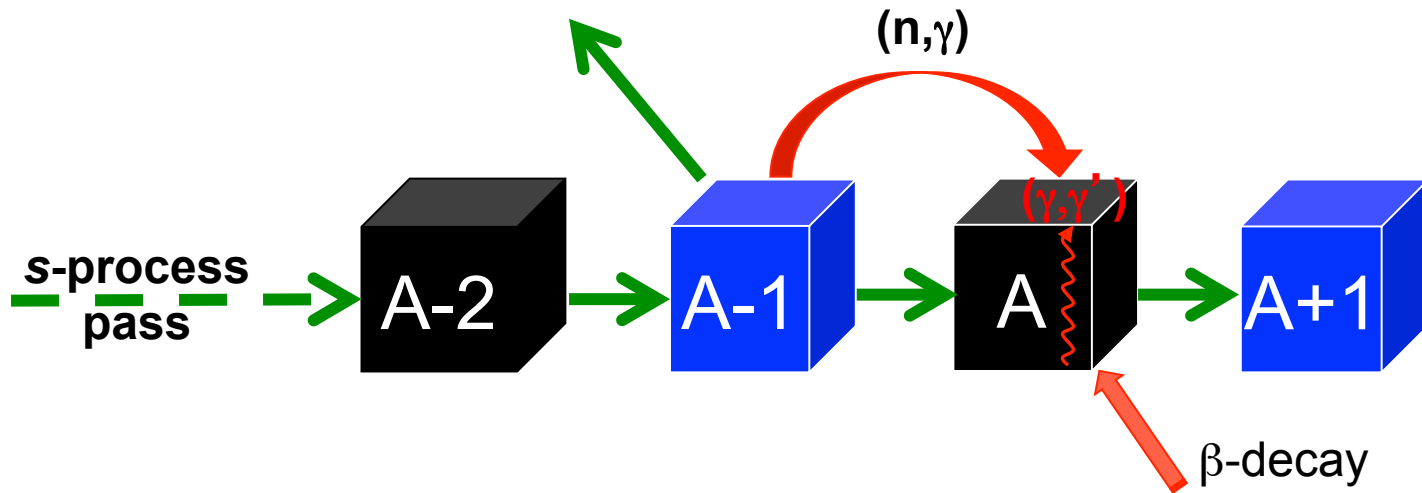
## Neutron Capture on s-Process Nuclei



**A is stable and A-1 is a s-process branching point nucleus**



# Neutron Capture on s-Process Nuclei



Leading indirect techniques:

- Nuclear Resonance Fluorescence (NRF)
- $\beta$ -delayed  $\gamma$  and neutron spectra
- Surrogate method

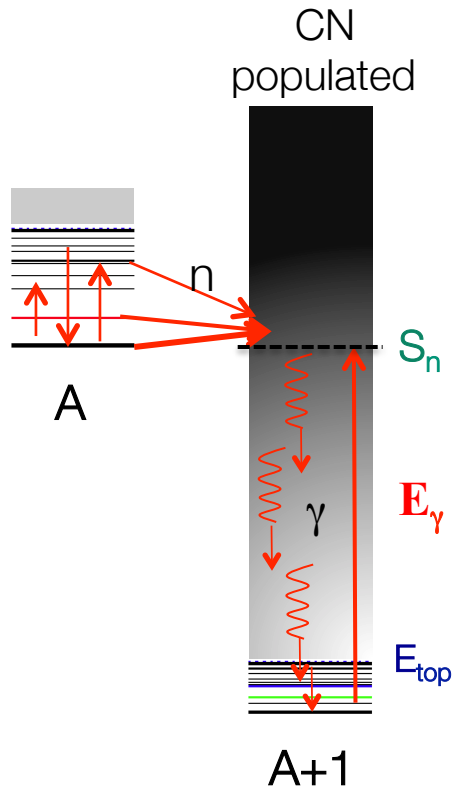
**A is stable and A-1 is a s-process branching point nucleus**



# Calculating (n,γ) Cross Sections

## Capture reaction

n+target → population of compound nucleus (CN)  
Subsequent decay by competition of γ emission and neutron evaporation



## Theoretical description

Hauser-Feshbach (HF) formalism:

$$\sigma_{n\gamma} = \frac{\pi}{k_n^2} \sum_{J,\pi} g_J \frac{T_\gamma(E, J, \pi) T_n(E, J, \pi)}{T_{tot}}$$

where

$T_\gamma$  is the gamma transmission coefficient

$$T_\gamma(E, J, \pi) = \sum_{X,\lambda} \int T_{X\lambda}(\epsilon_\gamma) \rho(E - \epsilon_\gamma) d\epsilon_\gamma$$

$T_n$  is the neutron transmission coefficient

$$\text{and } T_{tot} = T_n + T_p + T_d + T_t + T_\alpha + T_\gamma$$

## Challenge for calculations

Need accurate description of:

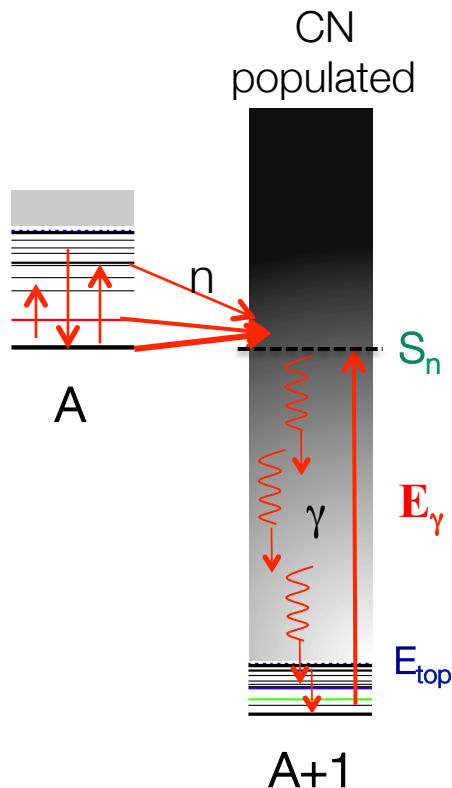
- $\gamma$ -ray strength function ( $\gamma$ SF)
- level densities (LDs)
- discrete low-lying levels with  $J^\pi$ , branching ratios – esp. important for isomers!
- optical model for n+target: (in decent shape ~10% below 1-2 MeV, 3-5% above)



# Calculating $(n,\gamma)$ Cross Sections from Photon Scattering

## Capture reaction

$n$ -target  $\rightarrow$  population of compound nucleus (CN)  
 Subsequent decay by competition of  $\gamma$  emission and neutron evaporation



## Theoretical description

Hauser-Feshbach (HF) formalism:

$$\sigma_{n\gamma} = \frac{\pi}{k_n^2} \sum_{J,\pi} g_J \frac{T_\gamma(E, J, \pi) T_n(E, J, \pi)}{T_{tot}}$$

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$T_n$  is the neutron transmission coefficient

$$\text{and } T_{tot} = T_n + T_p + T_d + T_t + T_\alpha + T_\gamma$$

## Constrain on $\gamma$ SF

$$T_{X\lambda}(\epsilon_\gamma) = 2\pi \epsilon_\gamma^{(2\lambda+1)} f_{X\lambda}(\epsilon_\gamma)$$

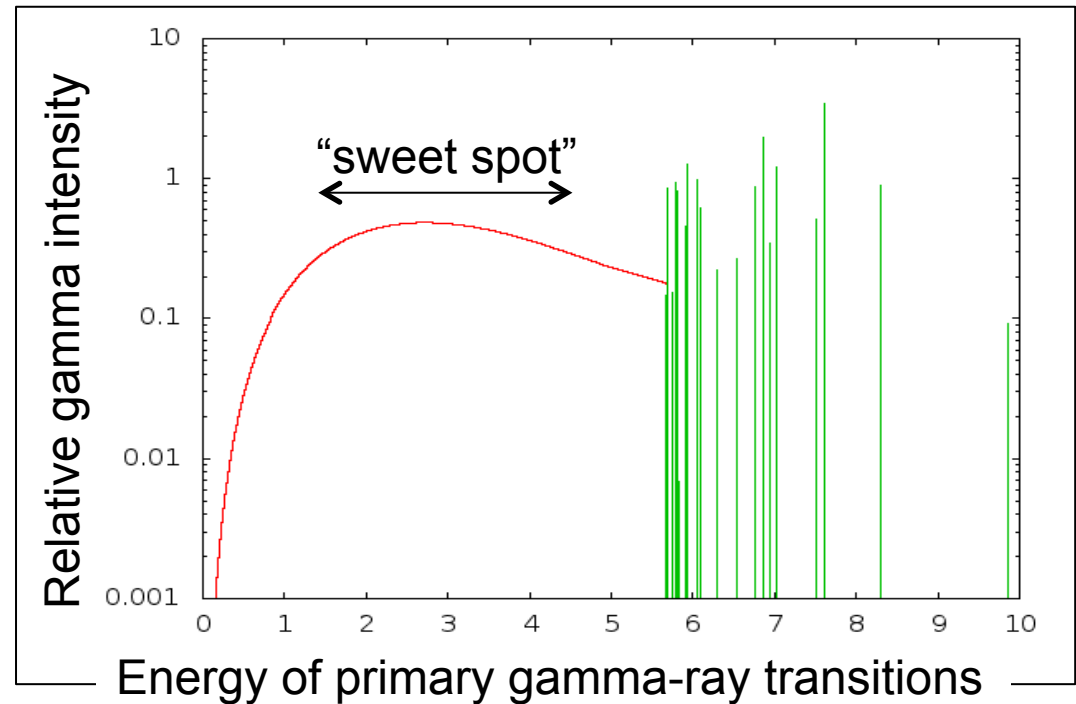
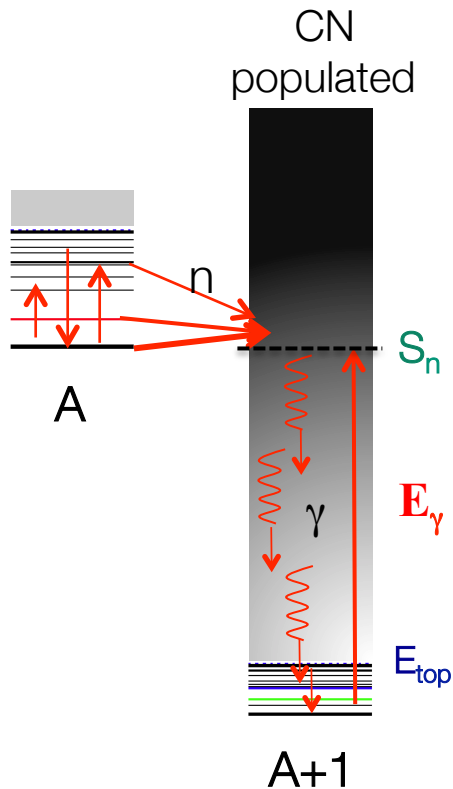
$$T_{E1}(E_\gamma) = 2\pi E_\gamma^3 f_{E1}(E_\gamma) \quad (\text{electric dipole})$$

$\sigma_{abs}(E_\gamma)$ : measured directly with real photons

$$\sigma_{abs}(E_\gamma) = 3(\pi \hbar c)^2 E_\gamma f_{E1}^\uparrow(E_\gamma)$$



# Direct measurement of the $\gamma$ strength function via $(\gamma, \gamma')$ measurement

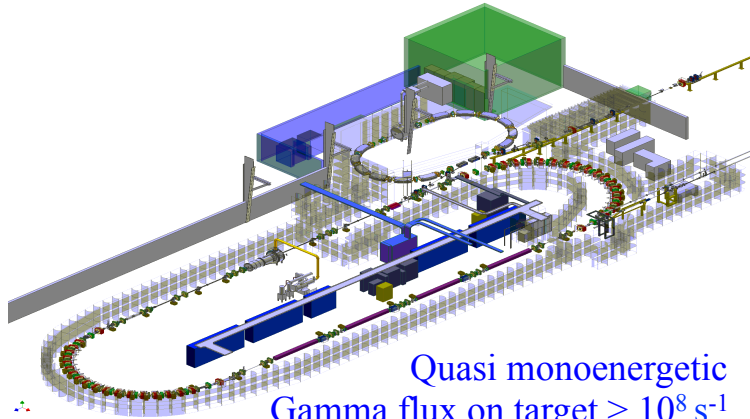


The  $\gamma$  strength function is governed by low-energy transitions !!!

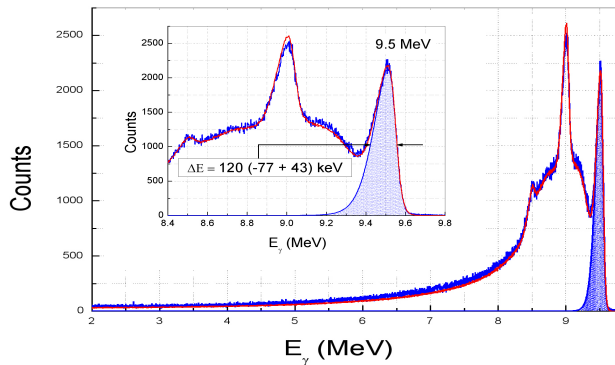


# Measuring the Nuclear Dipole Response using Real Photons

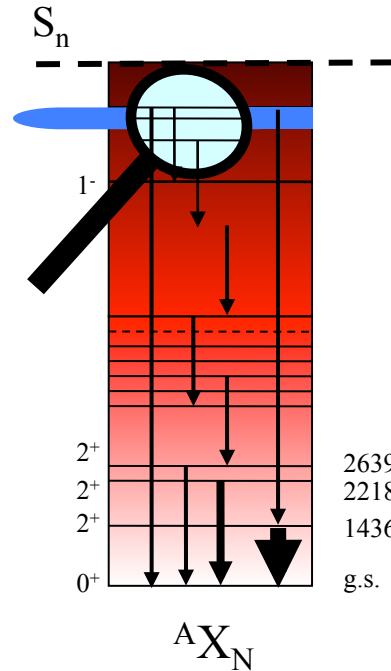
## High-Intensity Gamma-Ray Source



Quasi monoenergetic  
Gamma flux on target  $> 10^8 \text{ s}^{-1}$   
Tunable beam from 1 – 100 MeV  
100% linear or circular polarization



## NRF Technique



## Experimental observables

- Excitation energy  $E_x$
- Spin and parity  $J, \pi$
- Decay width  $\Gamma_0$
- Branching ratio  $\Gamma_i/\Gamma$

In a completely model independent way !

## HIGS Advantages

- $\sigma_{\text{el}} = f(E_\gamma)$  (from primary g.s. transitions)
- $\sigma_{\text{inel}} = f(E_\gamma)$  (from secondary transitions)
- $\sigma_{\text{tot}} = \sigma_{\text{el}} + \sigma_{\text{inel}} = \sigma_{\text{abs}}$

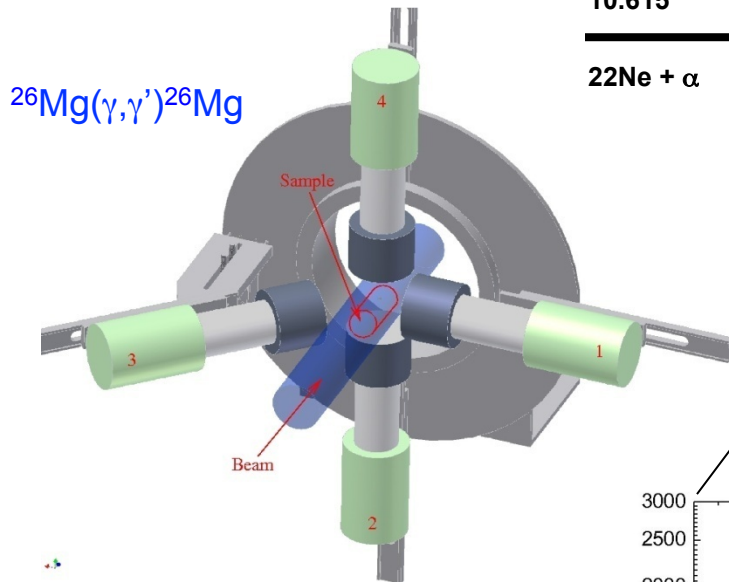
N. Pietralla, A. Tonchev et al. PRL **88** (2001) 012502



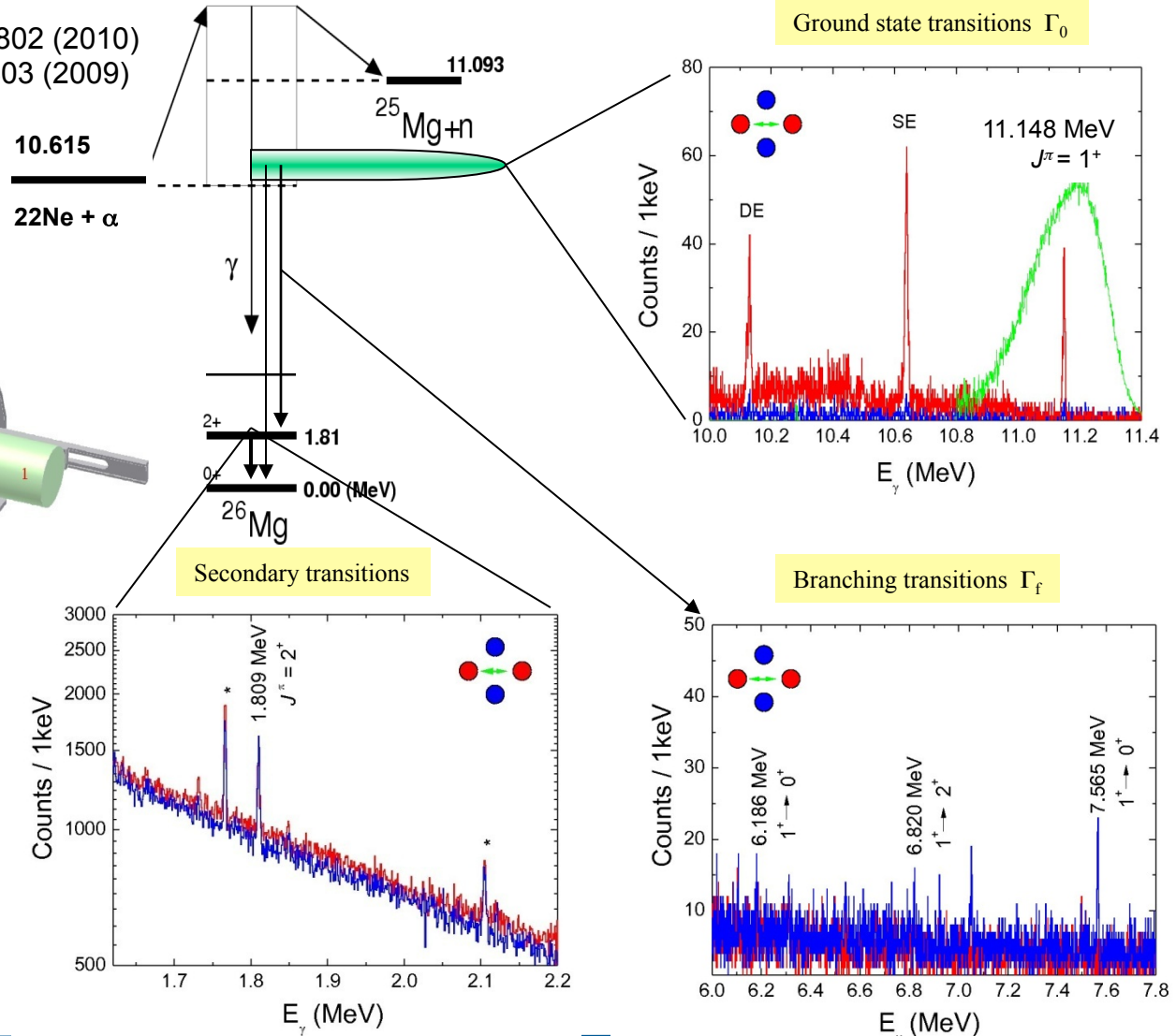


# Photon Scattering from Light-Mass Nuclei

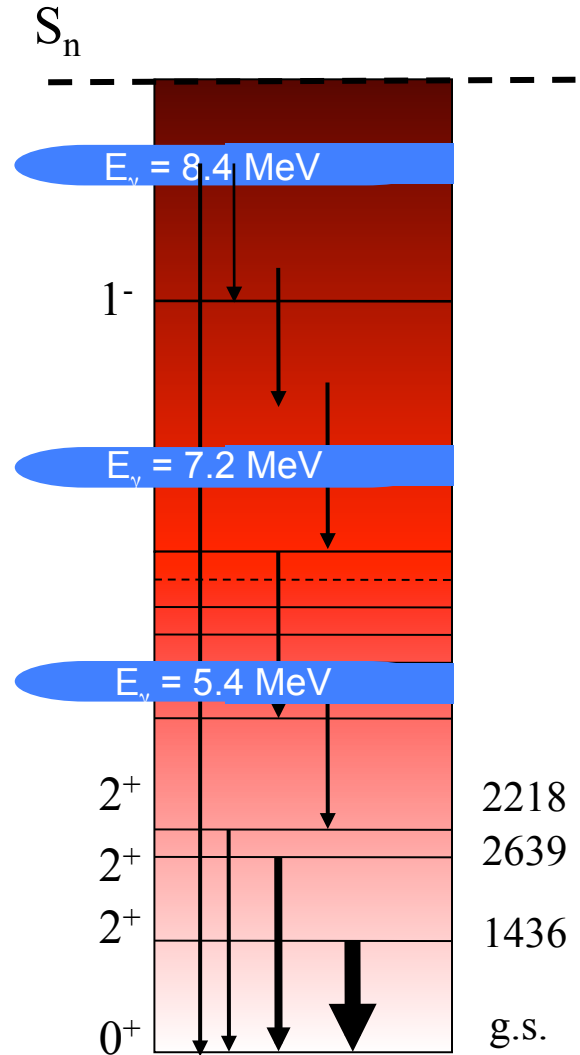
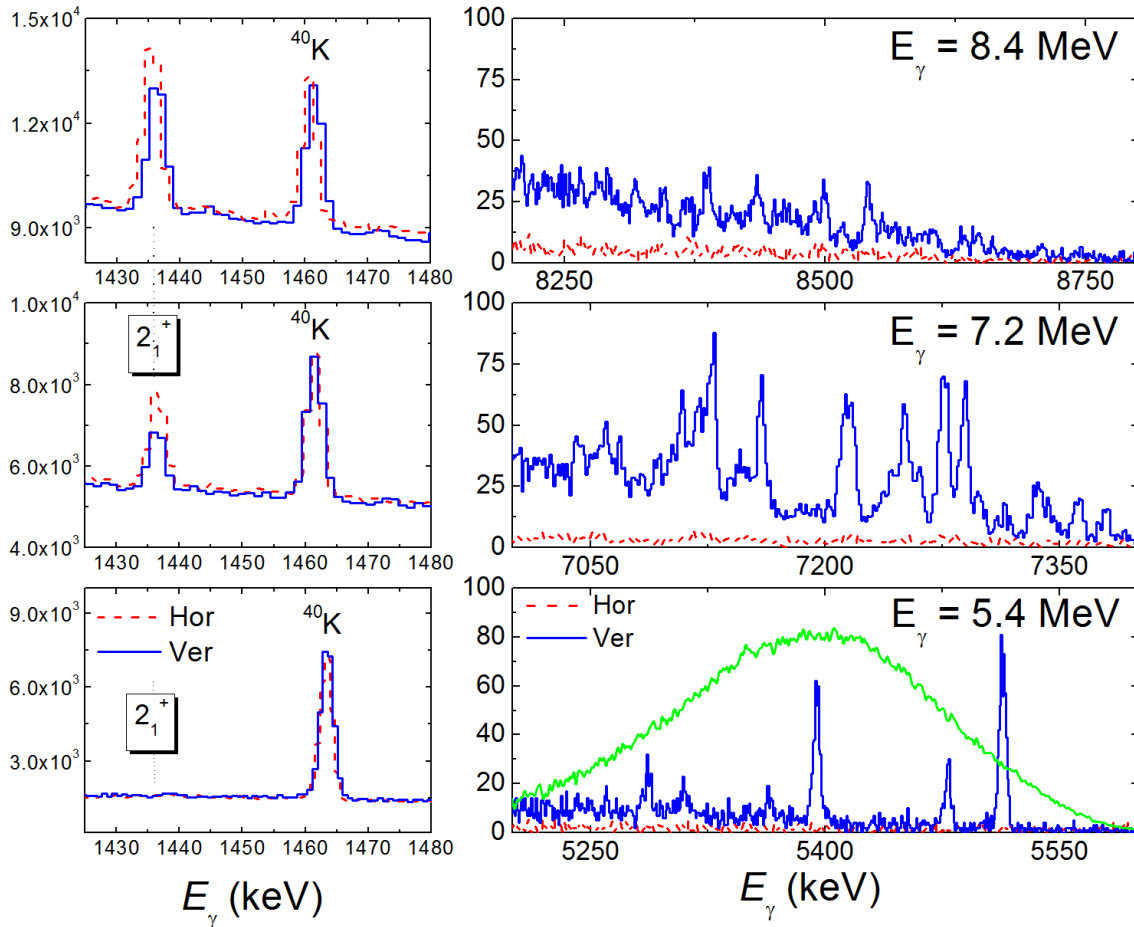
R.J. De Boer et al. Phys. Rev. C **82**, 025802 (2010)  
 R. Longland et al. Phys. Rev. C **80**, 055803 (2009)



**HIGS detection sensitivities:**  
 resonance states with  
 $\Gamma_{\text{tot}} \geq 1\text{meV}$

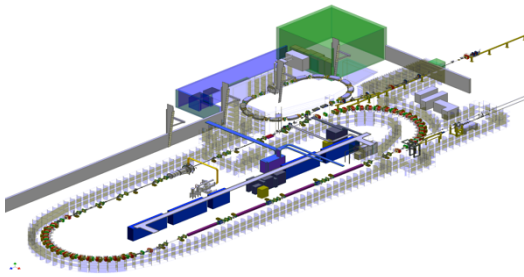


# Photon Scattering from Mid- and Heavy-Mass Nuclei

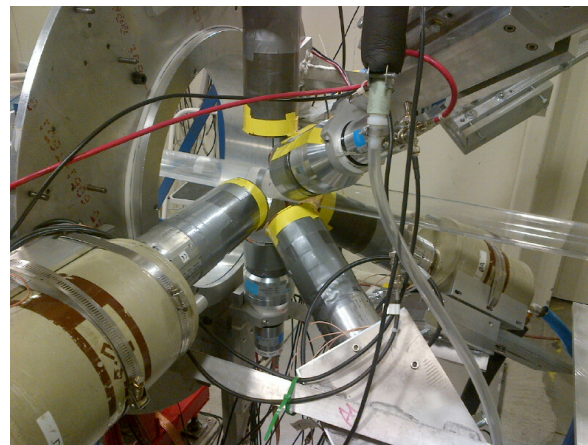
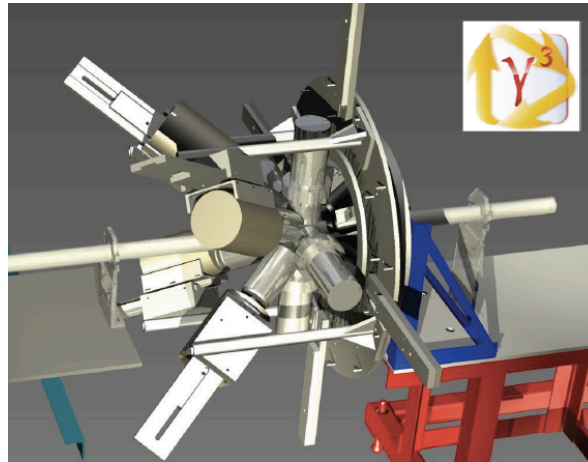


# Constructed the New $\gamma^3$ Setup at the High Intensity Gamma-Ray Source Facility (HIGS) at Duke University

HIGS facility



Schematic drawing of the  $\gamma^3$  experimental setup

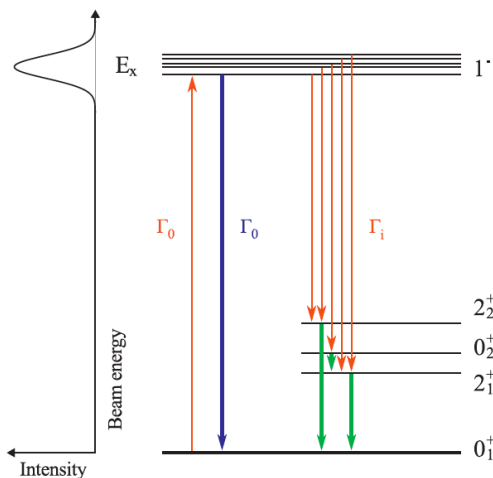


$\gamma^3$  Setup:

Array of four LaBr, four HPGe, and four neutron detectors.

Provides:

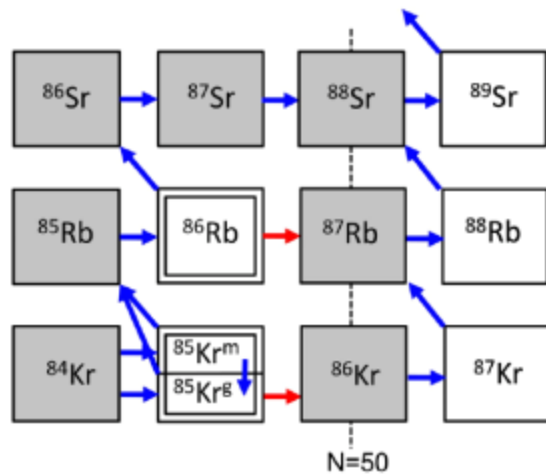
Absolute and model independent cross-section measurements below or above the neutron separation energy, hence direct measurement of the gamma-ray strength function.



B. Loher, V. Derya et al. NIMA 723, 136 (2013)



# Cross-Section Measurements of the $^{86}\text{Kr}(\gamma, n)$ Reaction to Probe the s-Process Branching at $^{85}\text{Kr}$



- The branching at  $^{85}\text{Kr}$  is significant for s-process modeling.
- This branching is independent of temperature and depends only on the neutron density.
- Helps to constrain the neutron density parameter in models of AGB stars.
- $^{85}\text{Kr}$  ( $T_{1/2} = 10.76$  y) neutron capture not measured.

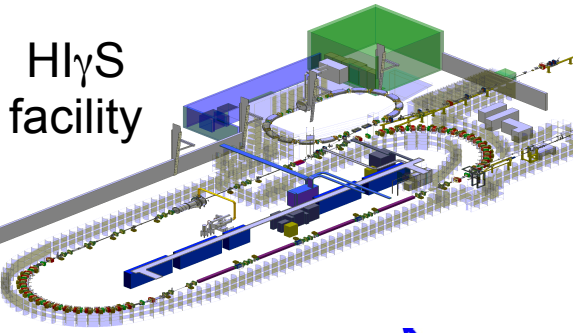
## Solution

Measure  $^{86}\text{Kr}(\gamma, \gamma')^{86}\text{Kr}$  and inverse  $^{86}\text{Kr}(\gamma, n)^{85}\text{Kr}$  reaction

Reproduce experimental results with the statistical model calculations

Apply same parameter set for  $(n, \gamma)$  cross section





# Proof-of-Principle Experiment on $^{85}\text{Kr}$ radioactive target

## Cross-Section Measurements of the $^{86}\text{Kr}(\gamma, n)$ Reaction to Probe the $s$ -Process Branching at $^{85}\text{Kr}$

R. Raut,<sup>1,2,\*</sup> A. P. Tonchev,<sup>1,2,†</sup> G. Rusev,<sup>1,2,‡</sup> W. Tornow,<sup>1,2</sup> C. Iliadis,<sup>3,2</sup> M. Lugaro,<sup>4</sup> J. Buntain,<sup>4</sup> S. Goriely,<sup>5</sup>  
J. H. Kelley,<sup>2,6</sup> R. Schwengner,<sup>7</sup> A. Banu,<sup>8</sup> and N. Tsoneva<sup>9,10</sup>

<sup>1</sup>Department of Physics, Duke University, Durham, North Carolina 27708, USA

<sup>2</sup>Triangle Universities Nuclear Laboratory, Durham, North Carolina 27708, USA

<sup>3</sup>Department of Physics and Astronomy, University of North Carolina at Chapel Hill, Chapel Hill, North Carolina 27695-8202, USA

<sup>4</sup>Monash Centre for Astrophysics (MoCA), Monash University, Victoria 3800, Australia

<sup>5</sup>Institut d'Astronomie et d'Astrophysique, Université Libre de Bruxelles, CP 226, 1050 Brussels, Belgium

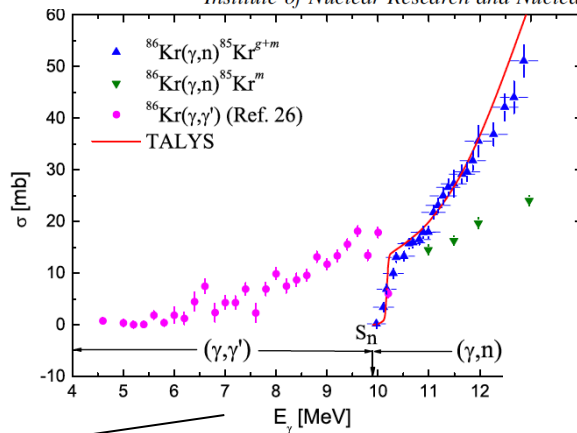
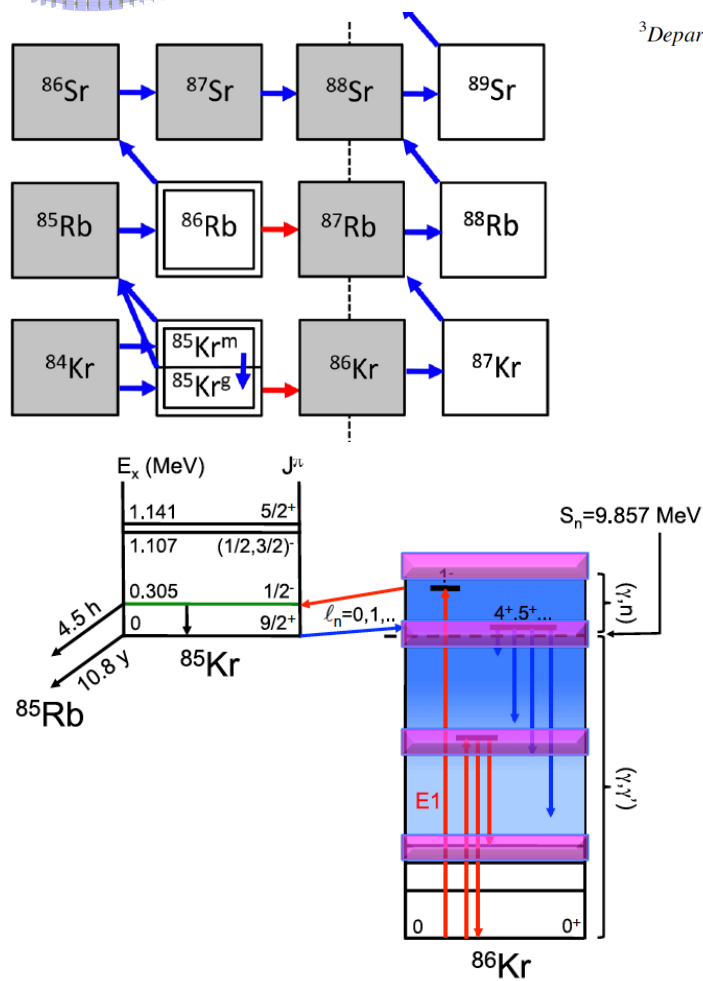
<sup>6</sup>Department of Physics, North Carolina State University, Raleigh, North Carolina 27695, USA

<sup>7</sup>Institut für Strahlenphysik, Helmholtz-Zentrum Dresden-Rossendorf, 01314 Dresden, Germany

<sup>8</sup>Department of Physics and Astronomy, James Madison University, Harrisonburg, Virginia 22807, USA

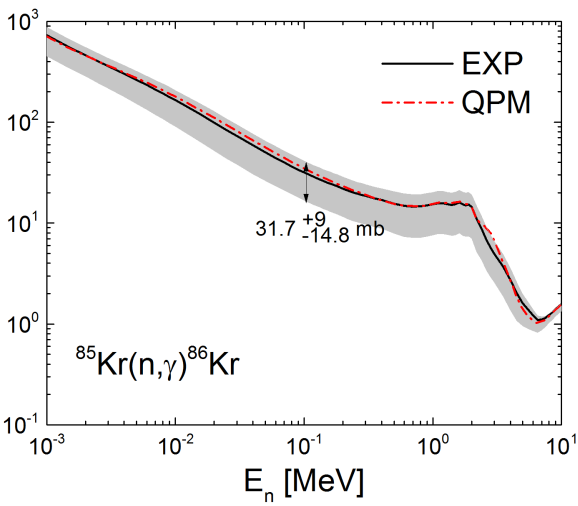
<sup>9</sup>Institut für Theoretische Physik, Universität Gießen, Gießen D-35392, Germany

<sup>10</sup>Institute of Nuclear Research and Nuclear Energy, 1784 Sofia, Bulgaria

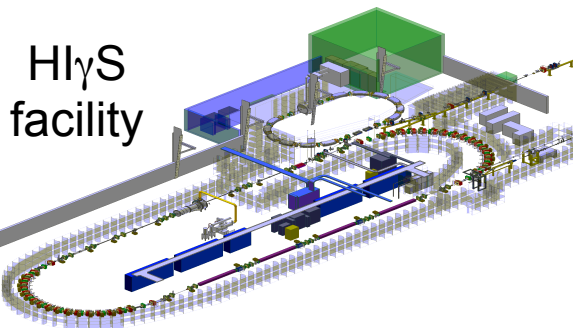


Obtained by statistical (HF) model

Measured directly with monoenergetic photon beams at HIGS



R. Schwengner *et al.*, PRC 87, 024306 (2013)



HIγS  
facility

# Proof-of-Principle Experiment on $^{85}\text{Kr}$ radioactive target

PRL 111, 112501 (2013)

PHYSICAL REVIEW LETTERS

week ending  
13 SEPTEMBER 2013

## Cross-Section Measurements of the $^{86}\text{Kr}(\gamma, n)$ Reaction to Probe the $s$ -Process Branching at $^{85}\text{Kr}$

R. Raut,<sup>1,2,\*</sup> A. P. Tonchev,<sup>1,2,†</sup> G. Rusev,<sup>1,2,‡</sup> W. Tornow,<sup>1,2</sup> C. Iliadis,<sup>3,2</sup> M. Lugaro,<sup>4</sup> J. Buntain,<sup>4</sup> S. Goriely,<sup>5</sup>  
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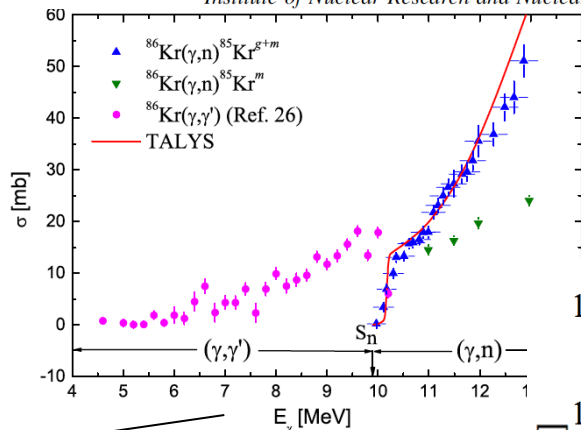
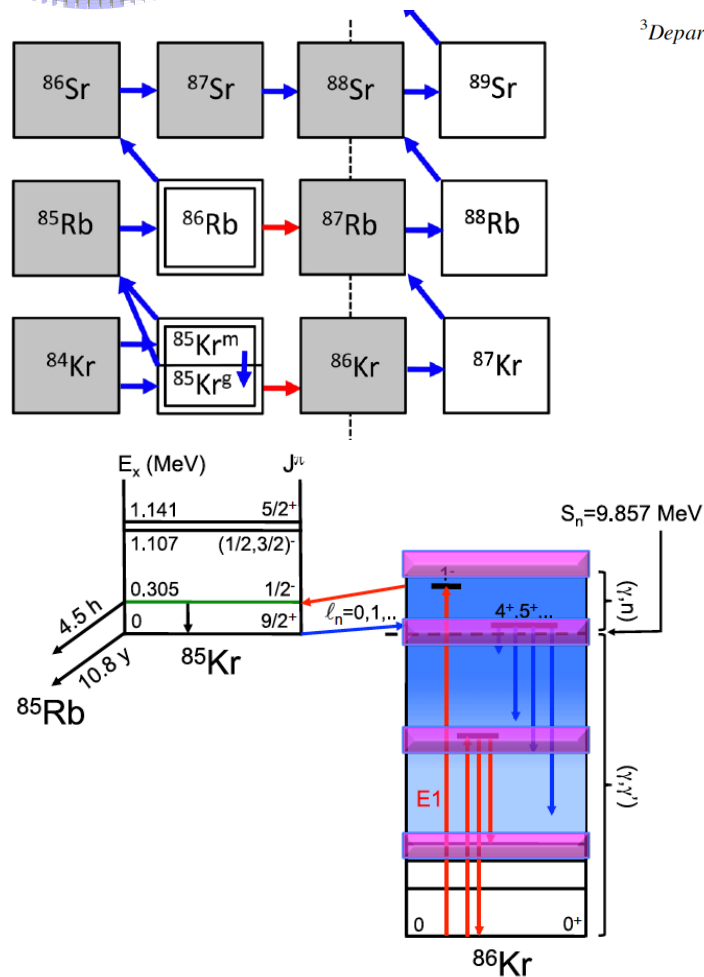
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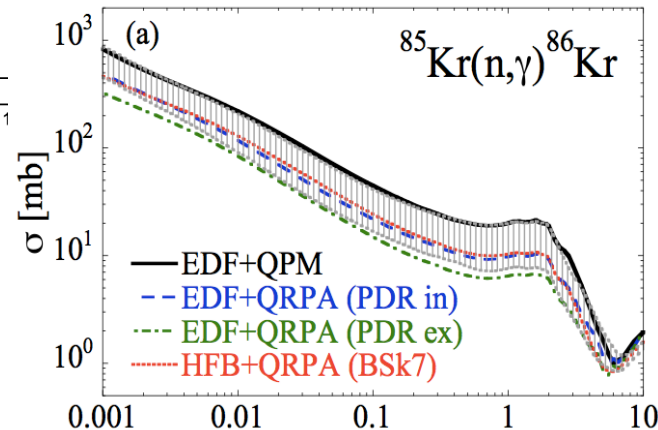
<sup>10</sup>Institute of Nuclear Research and Nuclear Energy, 1784 Sofia, Bulgaria



Obtained by  
statistical (HF)  
model

Measured directly  
with monoenergetic  
photon beams at  
HIGS

R. Schwengner *et al.*, PRC 87,  
024306 (2013)



N. Tsoneva *et al.*, PRC 91,  
044318 (2015)

# $\beta$ -Delayed $\gamma$ Spectra from Total Absorption Spectrometer (TAS)

PRL 113, 232502 (2014)

PHYSICAL REVIEW LETTERS

week ending  
5 DECEMBER 2014

## Novel technique for Constraining $r$ -Process ( $n, \gamma$ ) Reaction Rates

A. Spyrou,<sup>1,2,3,\*</sup> S. N. Liddick,<sup>1,4,†</sup> A. C. Larsen,<sup>5,‡</sup> M. Guttormsen,<sup>5</sup> K. Cooper,<sup>1,4</sup> A. C. Dombos,<sup>1,2,3</sup>  
D. J. Morrissey,<sup>1,4</sup> F. Naqvi,<sup>1</sup> G. Perdikakis,<sup>6,1,3</sup> S. J. Quinn,<sup>1,7,3</sup> T. Renstrøm,<sup>5</sup> J. A. Rodriguez,<sup>1</sup>  
A. Simon,<sup>1,8</sup> C. S. Sumithrarachchi,<sup>1</sup> and R. G. T. Zegers<sup>1,7,3</sup>

### TECHNIQUE USED:

Combine  $\beta$ -decay and  $\gamma$ -emission measurements using  $\gamma$ -ray total absorption spectrometer (TAS). Analyse  $\gamma$ -ray spectra using the 'Oslo'-method.

### OBJECTIVE:

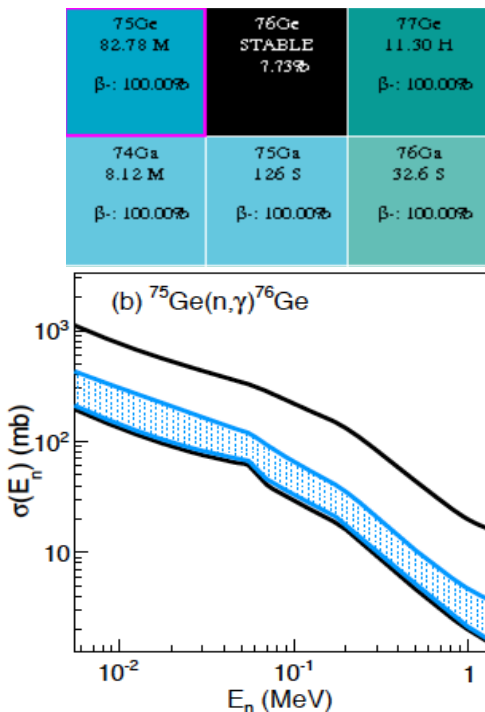
"Experimental" determination of the nuclear level density (NLD) and the  $\gamma$ -ray strength function ( $\gamma$ SF). These two quantities (and the nucleon-nucleus optical model potential) are essential for calculating the neutron-capture cross section.

### RESULT:

Proof-of-principle application to  $^{75}\text{Ge}(n,\gamma)$  shows good agreement with known result.

### OPEN QUESTIONS:

1. Approach made heavy use of available nuclear structure information. How much independent information does this type of measurement yield?
2. Only the functional form of the NLD and  $\gamma$ SF are obtained from the primary  $\gamma$ -ray spectra from 1.2 – 5.8 MeV.
3. The slope and absolute value must be determined by other means.

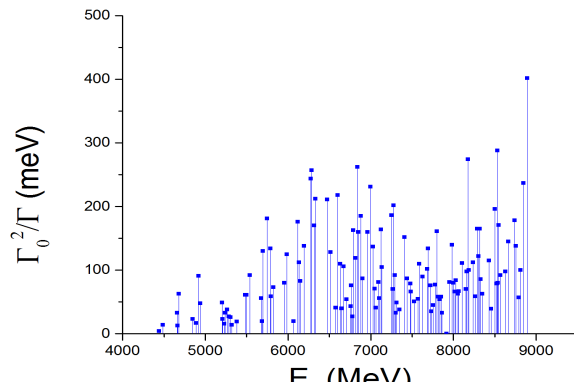


# Nuclear Resonance Fluorescence technique (NRF) to measure the nuclear dipole response

## Dipole response of $^{76}\text{Ge}$ between 4.5 and 9 MeV

R. S. Ilieva,<sup>1,2</sup> P. Humby,<sup>1,2,\*</sup> N. Cooper,<sup>1</sup> V. Werner,<sup>1,3</sup> G. Rusev,<sup>4,5,†</sup> C. Bernards,<sup>1</sup> J. Beller,<sup>3</sup> J. Bliss,<sup>3</sup> B. P. Crider,<sup>6</sup> P. M. Goddard,<sup>2</sup> J. Isaak,<sup>3</sup> J.H. Kelly,<sup>5</sup> B. Löher,<sup>7</sup> E. E. Peters,<sup>8</sup> N. Pietralla,<sup>3</sup> C. Romig,<sup>3</sup> D. Savran,<sup>7</sup> M. Scheck,<sup>3,‡</sup> A. P. Tonchev,<sup>4,§</sup> W. Tornow,<sup>4,5</sup> S. W. Yates,<sup>6,8</sup> and M. Zweidinger<sup>3</sup>

(Submitted to PRC)



TECHNIQUE USED:

Nuclear resonance fluorescence technique.

OBJECTIVE:

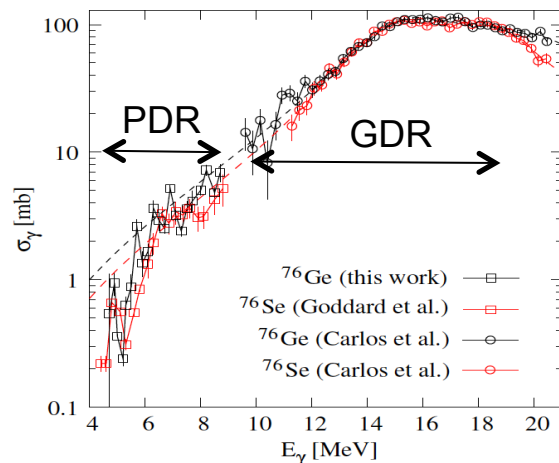
Direct measurements of the photoabsorption cross section below the neutron separation energy, hence direct measure of the  $\gamma$ -ray strength function ( $\gamma$ SF) needed for inverted ( $n,\gamma$ ) calculations.

RESULT:

Photoabsorption cross section of  $^{76}\text{Ge}(\gamma,\gamma')$  from 4 to Sn (9.4 MeV).

OPEN QUESTIONS:

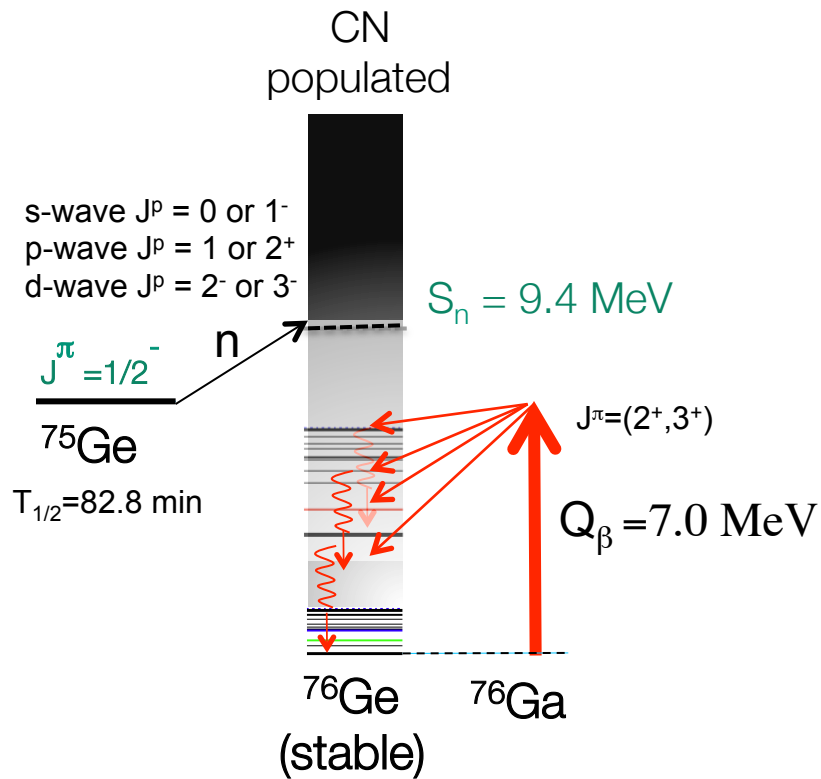
This approach is applicable for systems where  $A$  is stable and  $A-1$  is a radioactive nucleus.



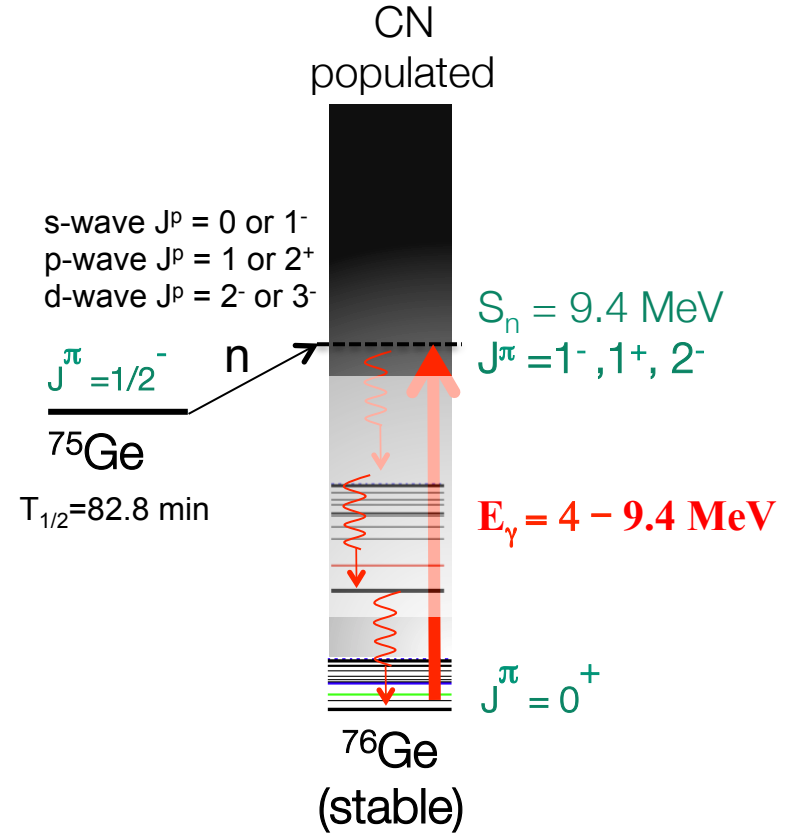


# Reaction Mechanism

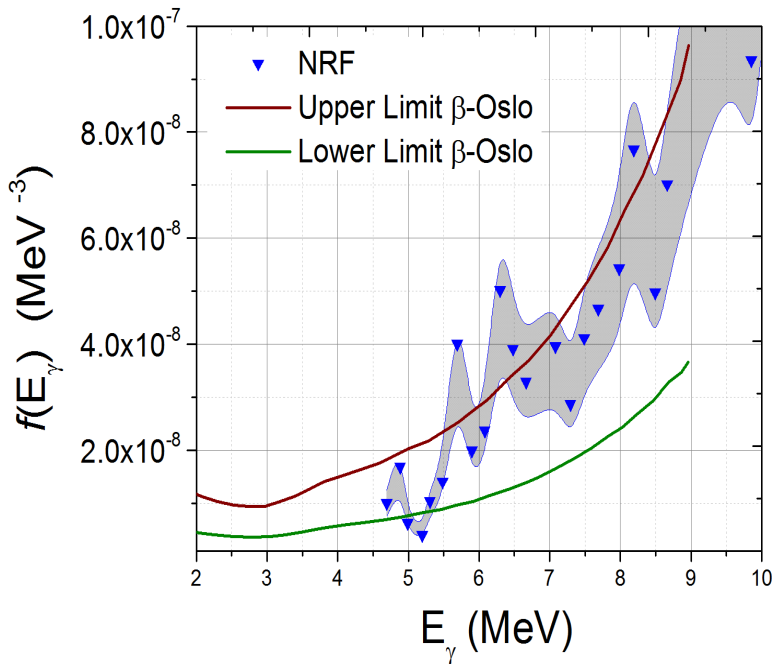
## $\beta$ -delayed neutron emission ( $\beta$ -Oslo)



## Photoabsorption (NRF)



## First Results



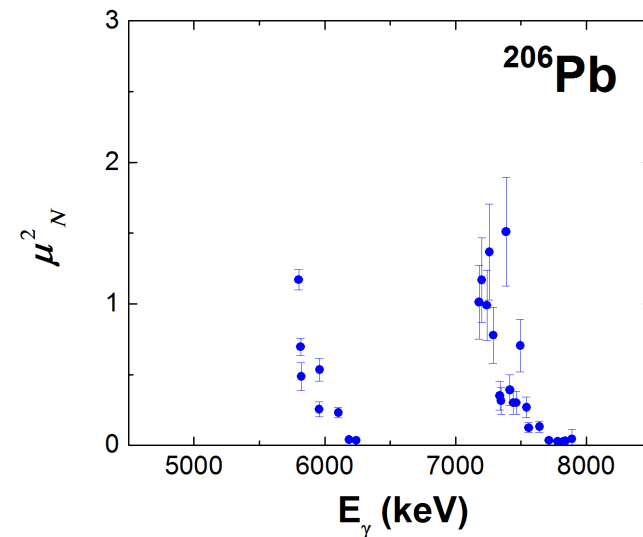
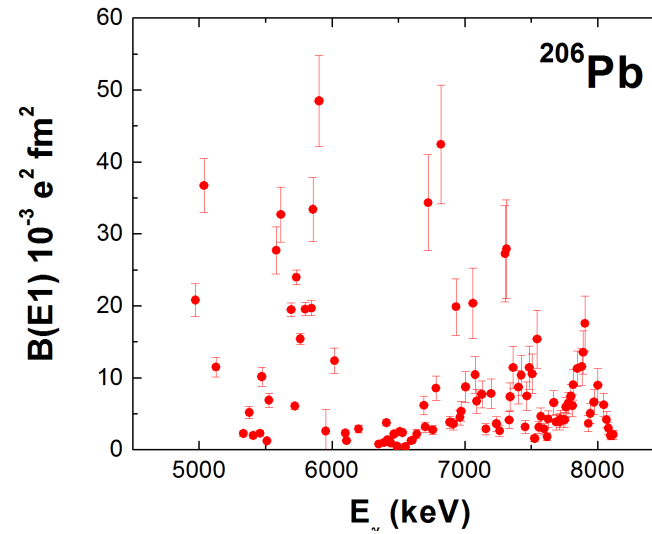
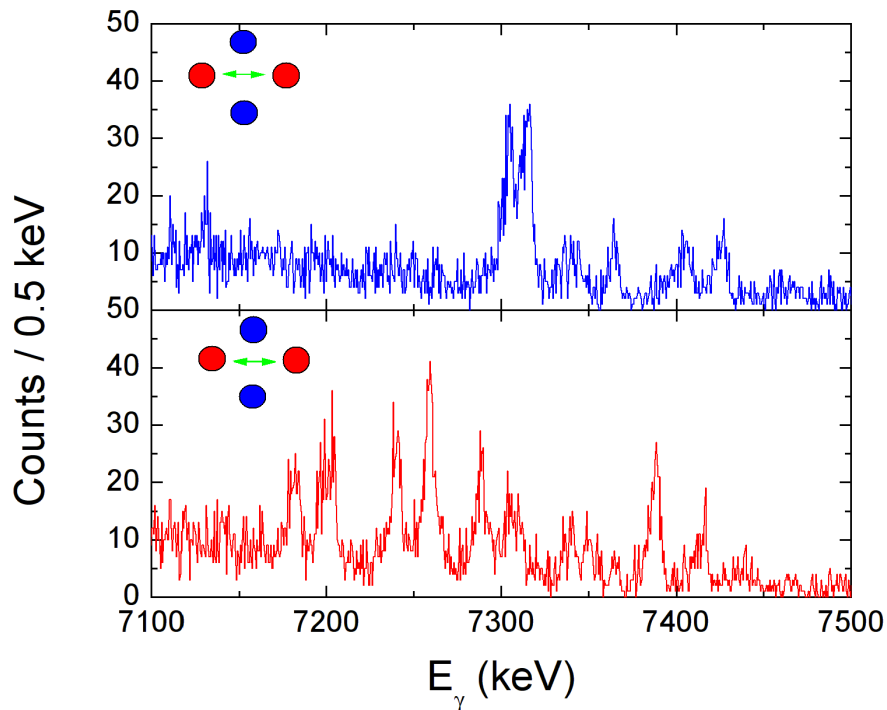
1.  $\gamma$ -ray strength function of  $^{76}\text{Ge}$  measured by the  $\beta$ -Oslo methods is slightly lower compared to the photonscattering or NRF experiment.
2.  $\gamma$ -strength function obtained by the photonscattering method show more structural effect compared to the one obtained via the  $\beta$ -Oslo methods.



# Fine Structure of the Electromagnetic Dipole Strength in $^{206}\text{Pb}$

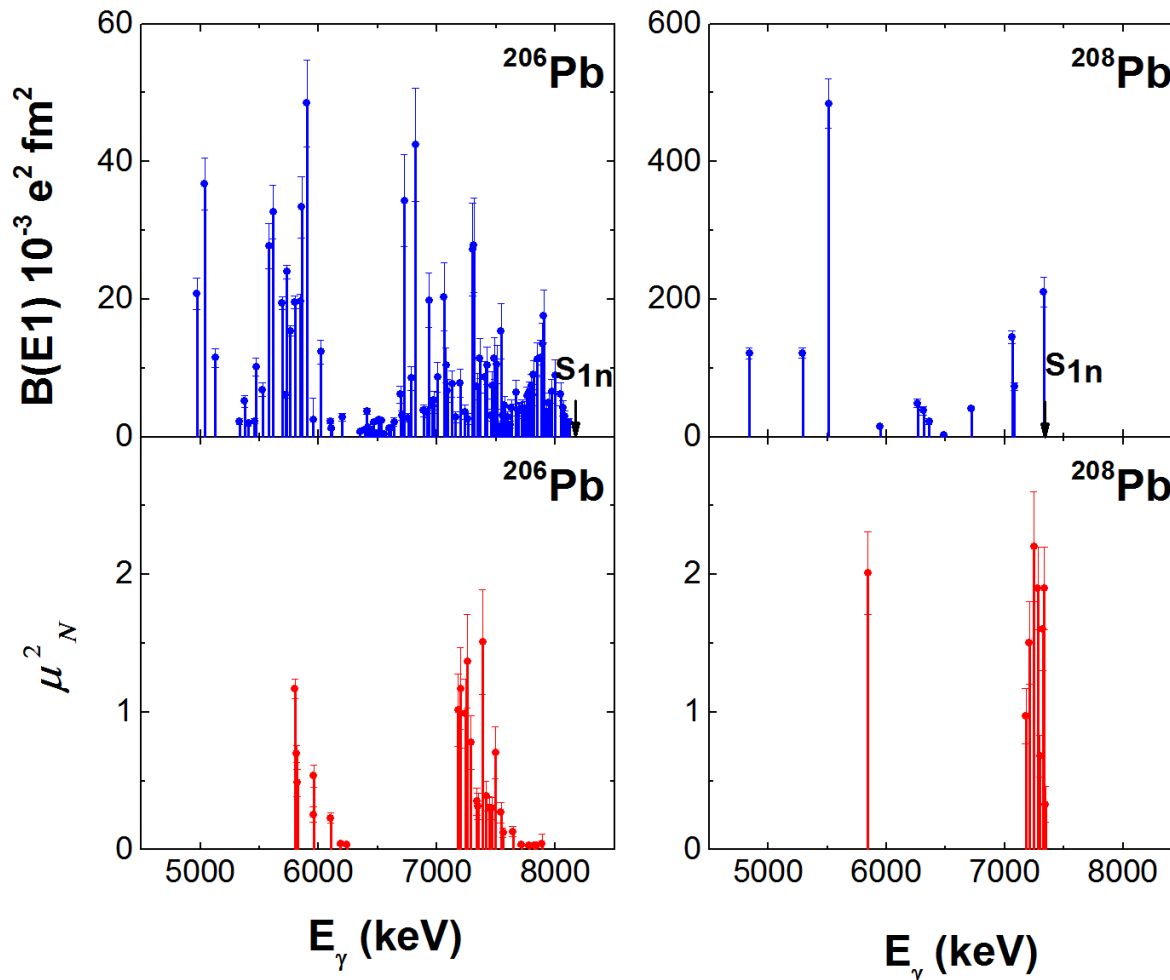
Measured:

- 99 E1 levels (40 new)
- 28 M1 levels (23 new)



# Fine Structure of the Electromagnetic Dipole Strength in $^{206}\text{Pb}$

C. Bhatia et al. In preparation



$^{208}\text{Pb}$

Number of E1 states = 12  
 $\Sigma B(E1)\uparrow = 1.23 \pm 0.11 \text{ e}^2 \text{ fm}^2$   
 $\Sigma B(E1)_{\text{QRPA}} = 1.13 \text{ e}^2 \text{ fm}^2$

$^{206}\text{Pb}$

Number of E1 states = 99  
 $\Sigma B(E1)\uparrow = 0.91 \pm 0.17 \text{ e}^2 \text{ fm}^2$   
 $\Sigma B(E1)_{\text{QPM}} = 1.03 \text{ e}^2 \text{ fm}^2$

$^{208}\text{Pb}$

Number of M1 states = 9  
 $\Sigma B(M1)\uparrow = 13.09 \pm 0.30 \text{ m}_N^2$

$^{206}\text{Pb}$

Number of M1 states = 28  
 $\Sigma B(M1)\uparrow = 13.3 \pm 3.0 \text{ m}_N^2$



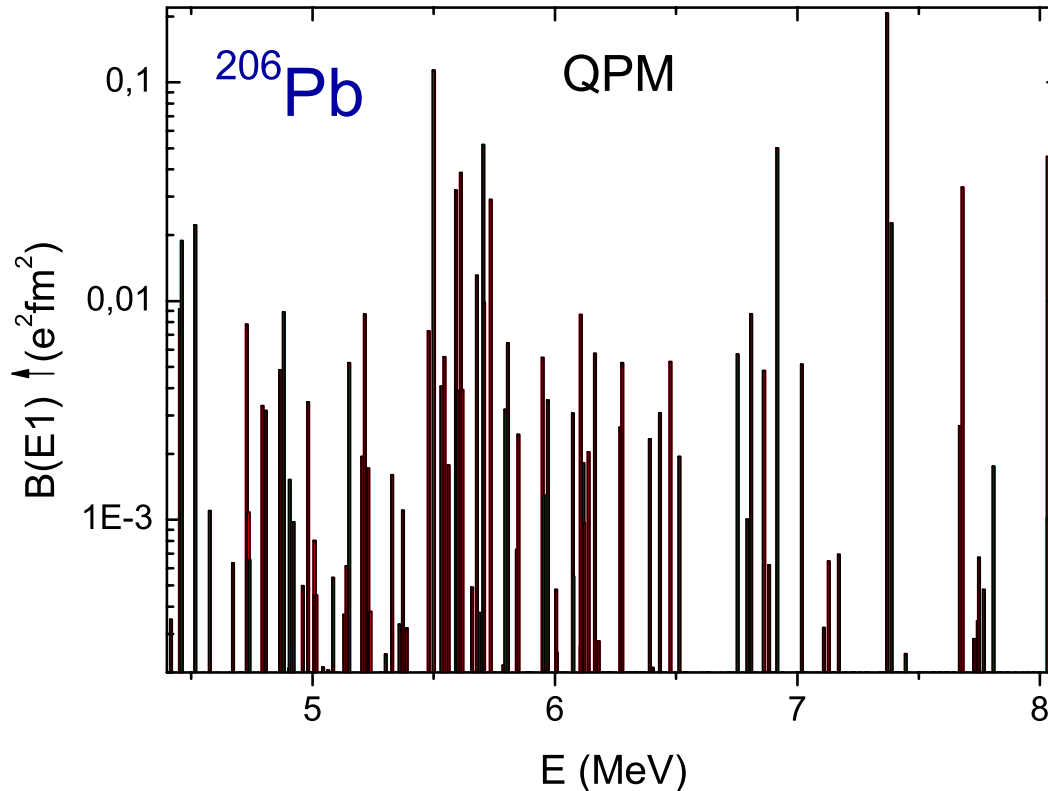
## Summary and Future Work

1. Good agreement is obtained for capture cross-section measurement on  $^{75}\text{Ge}$  from the NRF and  $\beta$ -Oslo methods.
2. Constraining the level density of  $^{76}\text{Ge}$  by using the discrete levels at low excitation energy and information from neutron-resonance experiments at the neutron separation energy  $S_n$  on neighboring germanium isotopes. Use the same normalization data as in the  $\beta$ -Oslo approach.
3. Plug the  $\gamma$ -ray strength function of  $^{76}\text{Ge}$  directly into the Hauser-Feshbach calculations to compare with the capture cross section on  $^{75}\text{Ge}$  nucleus obtained by the  $\beta$ -Oslo methods. Define the uncertainty of the predicted capture cross section based on both methods.



## Spare Slides

# B(E1) Strength in $^{206}\text{Pb}$ from QPM Calculations



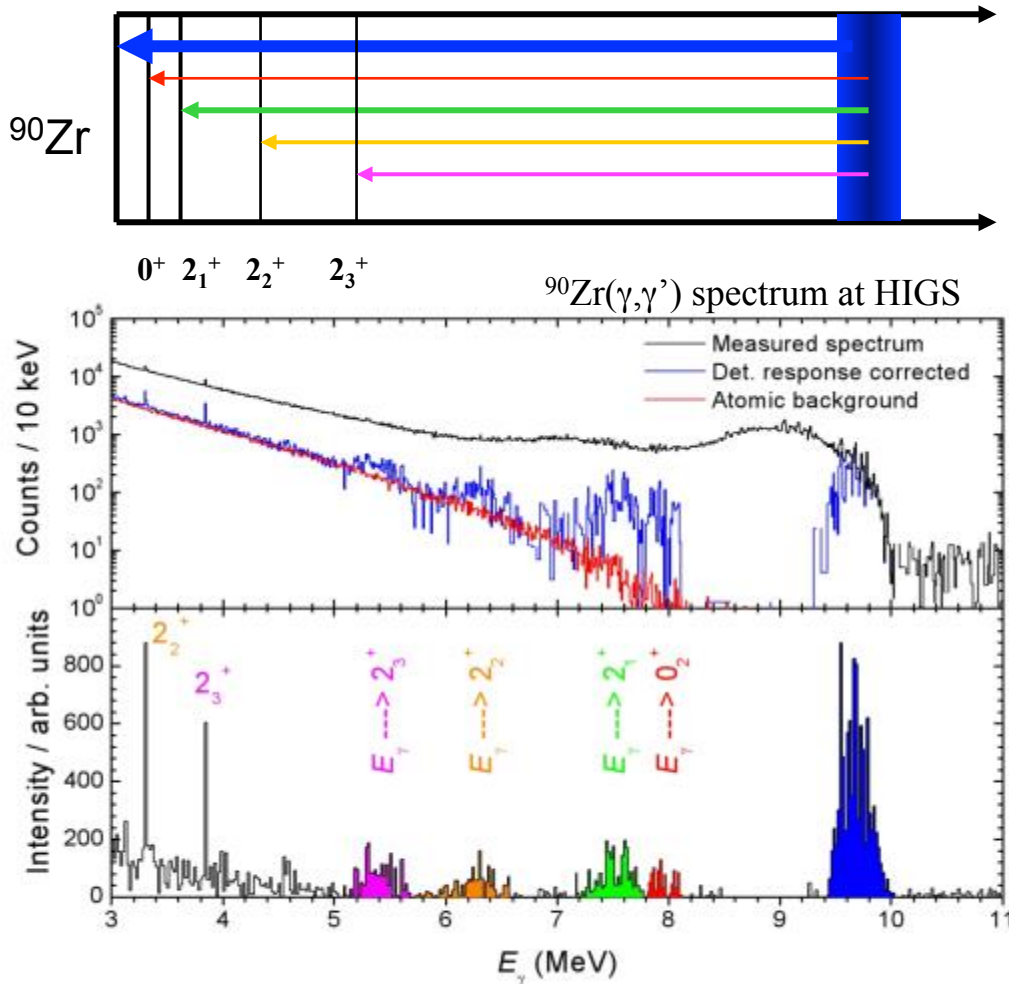
$^{206}\text{Pb}$ : QPM Calculations

Number of E1 states = 430

$\Sigma B(E1)_{\text{QPM}} \uparrow = 2.18 e^2\text{fm}^2$

Includes 2 and 3 phonon states

# Investigation of Photon-Strength Functions: $^{90}\text{Zr}$ case



□ Photon-strength function describes energy distribution of photon emission from high-energy states.

$$f(E_\gamma) = \langle \Gamma_\gamma / D E_\gamma^3 \rangle$$

□ Importance: Astrophysical network calculations; new fast nuclear reactors, statistical models.

## Preliminary results

□ E1 is the dominant multipolarity transition

□ Primary transitions are strongly dictated by the microscopic properties of the low-lying levels.

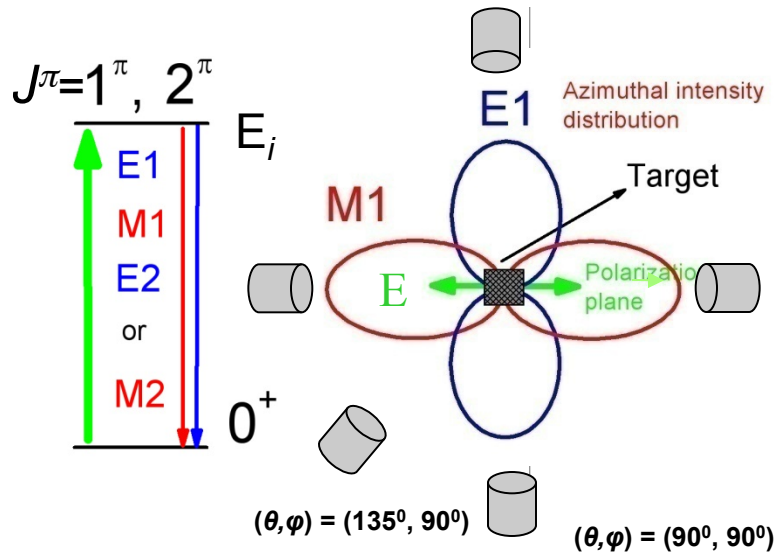
➤ PSF is not smooth curve below the  $B_n$  and  $E_\gamma > 4$  MeV.





# Parity Measurements with a Linear Polarized Photon Beams

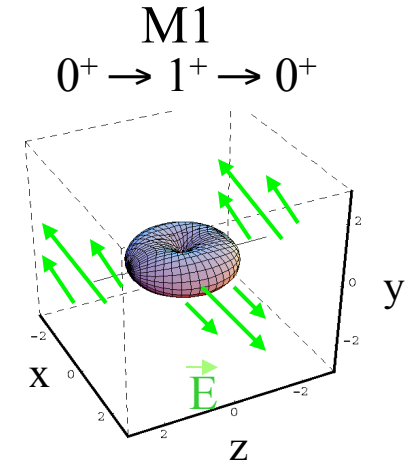
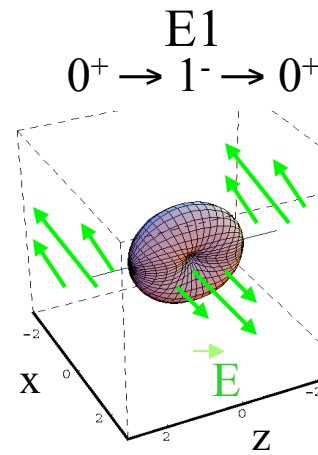
## Azimuthal distribution



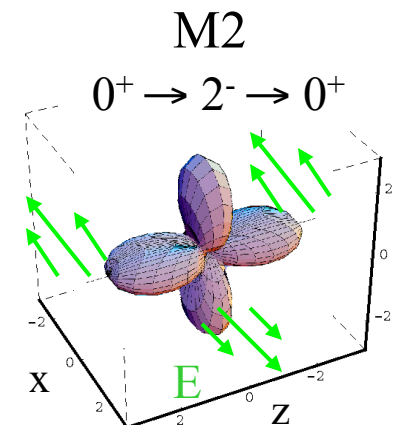
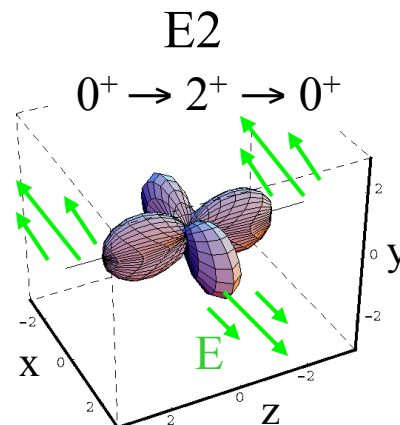
$$\Sigma = \frac{W(90^\circ, 0^\circ) - W(90^\circ, 90^\circ)}{W(90^\circ, 0^\circ) + W(90^\circ, 90^\circ)} = \pi_1 = \begin{cases} +1 & \text{for } J^\pi = 1^+, 2^+ \\ -1 & \text{for } J^\pi = 1^-, 2^- \end{cases}$$

Experimental Asymmetry of 0.96

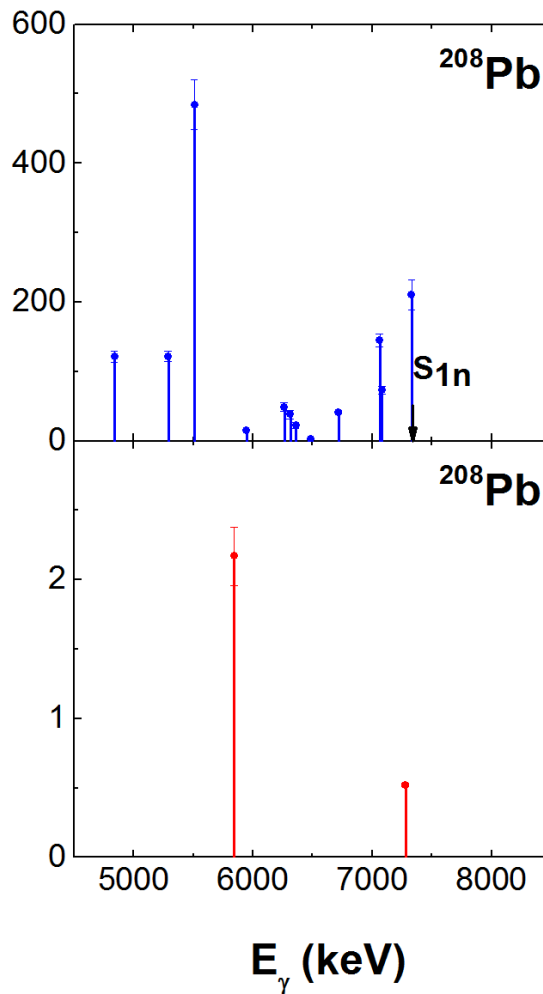
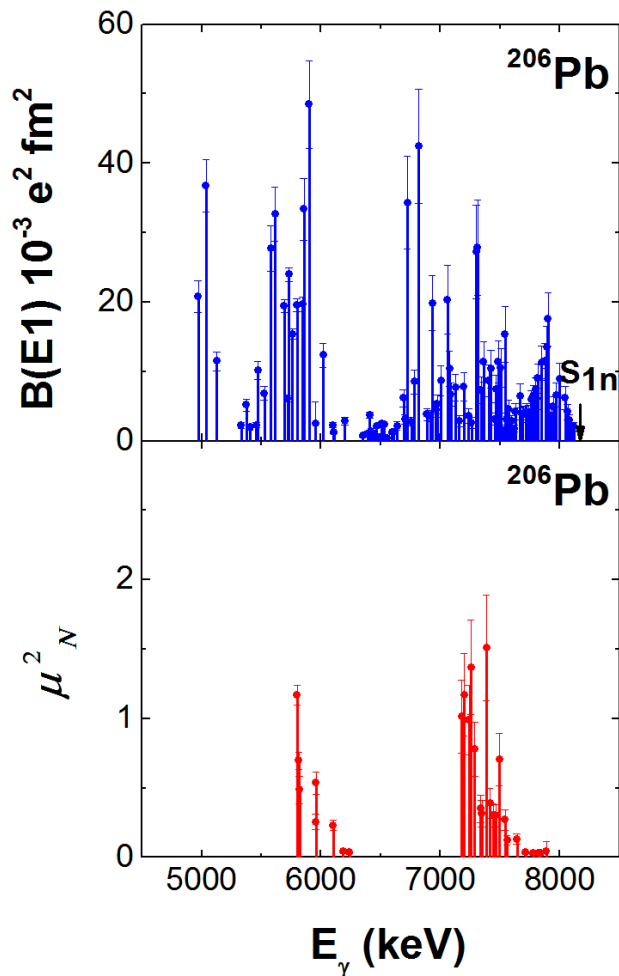
## Dipole



## Quadrupole



# Fine Structure of the Electromagnetic Dipole Strength in $^{206}\text{Pb}$



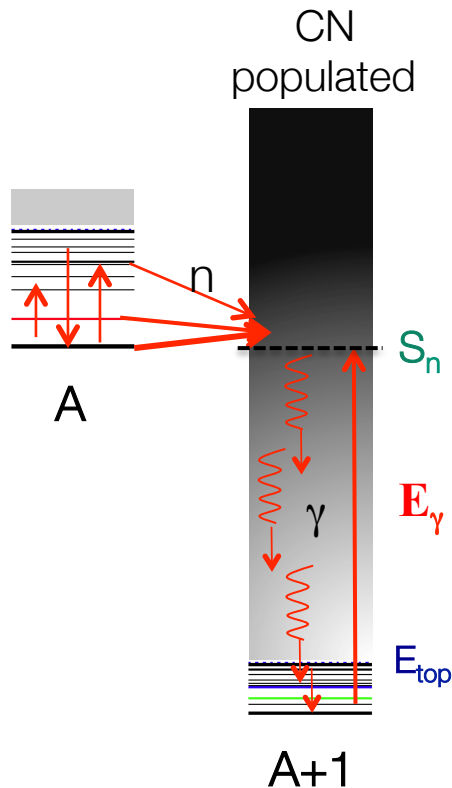
- $^{208}\text{Pb}$
- $\Sigma B(E1) \uparrow = 1.23 \pm 0.11 \text{ e}^2 \text{ fm}^2$
  - $B(E1)_{\min} = 0.003 \text{ e}^2 \text{ fm}^2$
  - $B(E1)_{\max} = 0.484 \text{ e}^2 \text{ fm}^2$
  - Number of states = 12
- $^{206}\text{Pb}$
- $\Sigma B(E1) \uparrow = 0.91 \pm 0.17 \text{ e}^2 \text{ fm}^2$
  - $B(E1)_{\min} = 0.37 \times 10^{-3} \text{ e}^2 \text{ fm}^2$
  - $B(E1)_{\max} = 48.5 \times 10^{-3} \text{ e}^2 \text{ fm}^2$
  - Number of states = 99
- $^{208}\text{Pb}$
- $\Sigma B(M1) \uparrow = 2.7 \pm 0.3 \text{ m}_N^2$
  - Number of states = 2
- $^{206}\text{Pb}$
- $\Sigma B(M1) \uparrow = 13.3 \pm 3.0 \text{ m}_N^2$
  - Number of states = 28



# Calculating (n,g) Cross Sections

## Capture reaction

n+target → population of compound nucleus (CN)  
 Subsequent decay by competition of  $\gamma$  emission and



## Theoretical description

Hauser-Feshbach (HF) formalism:

$$\sigma_{n\gamma} = \frac{\pi}{k_n^2} \sum_{J,\pi} g_J \frac{T_\gamma(E, J, \pi) T_n(E, J, \pi)}{T_{tot}}$$

where

$T_\gamma$  is the gamma transmission coefficient

$$T_\gamma(E, J, \pi) = \sum_{X,\lambda} \int T_{X\lambda}(\epsilon_\gamma) \rho(E - \epsilon_\gamma) d\epsilon_\gamma$$

$T_n$  is the neutron transmission coefficient

$$\text{and } T_{tot} = T_n + T_p + T_d + T_t + T_\alpha + T_\gamma$$

## In the case of (n, $\gamma$ ):

- $T_\gamma \ll T_n + T_p + T_d + T_t + T_\alpha$ , ignore  $T_\gamma$  in  $T_{tot}$
- $T_p, T_d, T_t, T_\alpha$  all small below thresholds

So  $T_n T_\gamma / T_n$  reduces to  $T_\gamma$

$$\sigma_{n\gamma} = \frac{\pi}{k_n^2} \sum_{J,\pi} g_J \sum_{X,\lambda} \int \underbrace{T_{X\lambda}(\epsilon_\gamma) \rho(E - \epsilon_\gamma) d\epsilon_\gamma}_{\text{Product of } \gamma\text{SF and LD}}$$

# Measuring the $^{142}\text{Ce}(g,n)$ reaction to infer the inverse $^{141}\text{Ce}(n, g)$ reaction

## Proof-of-Principle Experiment on $^{141}\text{Ce}$ radioactive target

$^{140}\text{Pr}$ 3.39 M $\epsilon$ : 100.00%	$^{141}\text{Pr}$ STABLE 100%	$^{142}\text{Pr}$ 19.12 H $\beta^-$ : 99.98% $\epsilon$ : 0.02%	$^{143}\text{Pr}$ 13.57 D $\beta^-$ : 100.00%	$^{144}\text{Pr}$ 17.28 M $\beta^-$ : 100.00%
$^{139}\text{Ce}$ 137.641 D $\epsilon$ : 100.00%	$^{140}\text{Ce}$ STABLE 88.450%	$^{141}\text{Ce}$ 32.508 D $\beta^-$ : 100.00%	$^{142}\text{Ce}$ >5E+16 Y 11.114% 2 $\beta^-$	$^{143}\text{Ce}$ 33.039 H $\beta^-$ : 100.00%
$^{138}\text{La}$ 1.02E+11 Y 0.08881% $\epsilon$ : 65.60% $\beta^-$ : 34.40%	$^{139}\text{La}$ STABLE 99.9119%	$^{140}\text{La}$ 1.67855 D $\beta^-$ : 100.00%	$^{141}\text{La}$ 3.92 H $\beta^-$ : 100.00%	$^{142}\text{La}$ 91.1 M $\beta^-$ : 100.00%

Knowing the energy dependence of  $(\gamma,n)$  cross sections is mandatory to predict the abundances of heavy elements using astrophysical models. The data can be applied directly or used to constrain the cross section of the inverse  $(n,\gamma)$  reaction.

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### Determination of the $^{142}\text{Ce}(\gamma,n)$ cross section using quasi-monoenergetic Compton backscattered $\gamma$ rays

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